

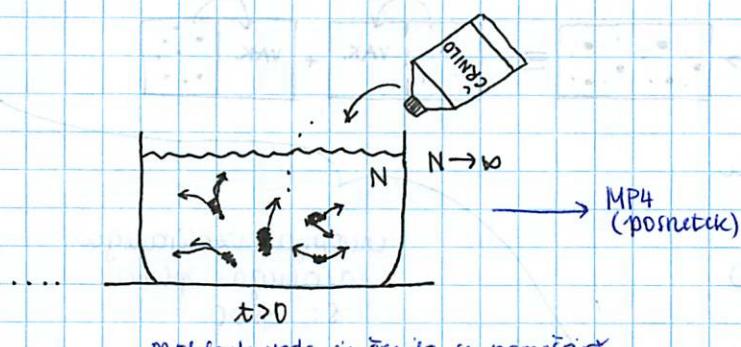
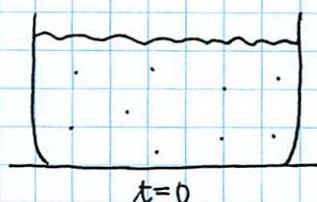
1. zakon termodinamike = energijski zakon
(= ohranjač Joule) $\Delta W_h = Q + A$

Če je posoda izolirana, ne „deluje“ močno okolje \rightarrow ne obstaja PERPETUUM MOBILE I. stopnje, ker energija ni ohranjač (superpredniki, vrtalka ti se vrčo \leftrightarrow ni trčja).

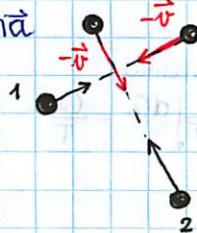
2. zakon termodinamike = entropijski zakon
(= za izoliran sistem) $\Delta S \geq 0$

ENTROPIJA

postkus:



$$\vec{F} = m\vec{a}$$

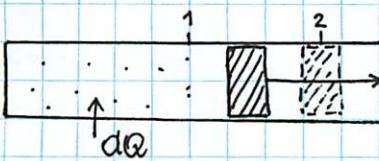


$$\vec{F} = m\ddot{\vec{r}} = m \frac{d^2\vec{r}}{dt^2} = m \frac{d\vec{r}}{dt} \cdot \frac{d\vec{r}}{dt} \Rightarrow ; t \rightarrow \tau \quad \vec{r} = \vec{r}(t) = \frac{d\vec{r}}{dt}$$

$$\frac{d\vec{r}}{dt} = \vec{v} = -\vec{v}$$

(= prav tako so sile enake, ne glede ali čas teče naprej ali nazaj)

Newtonovi zakoni omogočajo, da ne glede na smrek česar, "ologodški" potekajo enako.



premaknemo fat!

$$dW = dA + dQ$$

$$mc_v dT = -pdV + dQ \quad / \frac{1}{T}$$

$$\frac{dQ}{T} = \frac{pdV}{T} + mc_v \frac{dT}{T}$$

$$pV = \frac{m}{M} RT$$

$$\frac{p}{T} = \frac{m}{M} R \frac{1}{V}$$

$$c_v = \frac{R}{M} \frac{1}{\kappa-1}$$

$$ds = \frac{dQ}{T} = m \frac{R}{M} \cdot \frac{1}{\kappa-1} \frac{dT}{T} + \frac{m}{M} R \frac{dV}{V}$$

$$\Delta S = \int_{S_1}^{S_2} ds = m \frac{R}{M} \frac{1}{\kappa-1} \ln \frac{T_2}{T_1} + \frac{m}{M} R \ln \frac{V_2}{V_1} = S_2 - S_1$$

Definicija:

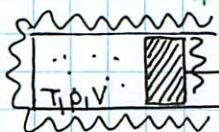
$$S = S_0 + \frac{mR}{M} \frac{1}{\kappa-1} \ln T + \frac{m}{M} R \ln V$$

Entropija idealnega plina:

$$W_n = \frac{m}{M} R \frac{1}{\kappa-1}$$

ΔS nam pove, koliko joulov toplote je bilo dovedeno pri meki T (= kelvini).

(a) Adiabatska spremembra (izentropna spremembra)



konstantna izentropija

$$\Delta S = 0$$

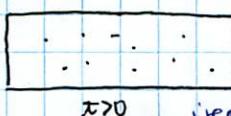
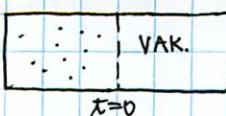
$pV^\kappa = \text{konstanta}$

$TV^{\kappa-1} = \text{konstanta}$

Entropija idealnega plina \approx

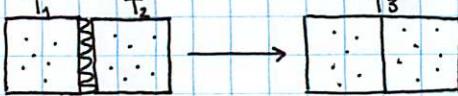
$$S = S_0 + \ln V + \ln T$$

(b) Hirnov postkus



reversibilna razširovanja

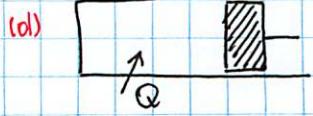
$$\Delta S = \frac{m}{M} R \ln \frac{2V}{V} = \frac{m}{M} R \ln 2 > 0$$



$$\Delta S = 2b \ln T_3 - (b \ln T_1 - b \ln T_2) = b \ln \frac{T_3^2}{T_1 T_2} = b \ln \frac{T_1^2 + T_2^2 + 2T_1 T_2}{4T_1 T_2} = b \ln \frac{(T_1 - T_2)^2 + 4T_1 T_2}{4T_1 T_2} \geq 0$$

$T_3 = \frac{T_1 + T_2}{2}$

(c) $\Delta S \geq 0$ (če sta $T_1 = T_2$)



$T = \text{konstanta}$
 $\Delta S \geq 0$



Nisi pogoji so enaki,
ki molekulama maza
je različna

$$\Delta S = \alpha_1 \ln 2 + \alpha_2 \ln 2 \geq 0$$

ustrezni konstanti
za 1. in 2. posodo

entropija izoliranega
idealnega plina
 $S: \Delta S \geq 0$

Sklep - 2. zakon (aksiom) termodinamike (S)

$$\frac{dW}{T} = \frac{dA}{T} + \frac{dQ}{T}$$

$$ds = \frac{dQ}{T} = -\frac{pdV}{T} + \frac{mc_v dT}{T}$$

$$\Delta S \geq \int \frac{dQ}{T} = \int \sum_{i=1}^n \frac{\Delta Q_i}{T_i}$$

↪ pogledaš limito $N \rightarrow \infty$

$$W_n = mc_v T = \frac{m}{M} R \frac{1}{N-1} T$$

$W_n(T)$!

$N = 1,6$ za enatomne molekule

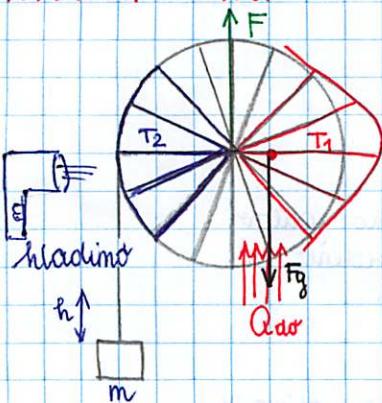
$N = 1,4$ za dvoatomne molekule

$N = 1,3$ za troatomne molekule

Če je T konstanta:

$$\frac{1}{T} \int dQ = \frac{Q}{T}$$

TOPLOTNI STROJI



- × težišči se premakne
- × $T_1 > T_2$
- × na kolo je navita vrv
- × $mgh = A$
- × izkoristek toplotnega stroja (= vedno manjši od 1)

$$\eta = \frac{A_{kr}}{Q_{do}} < 1$$

- × ciklične spremembe

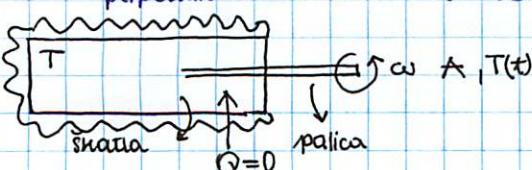


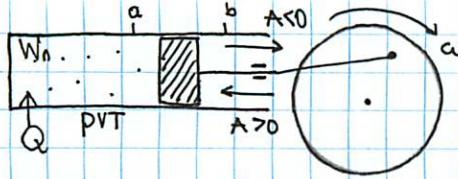
$$\gamma = \frac{A_{kr}}{Q_{do}}$$

koliko dela
iz čemer ciklu
dovedeni toplotni
v enem ciklu

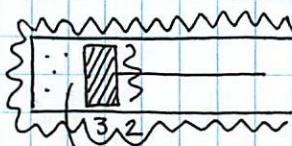
Poznamo: perpetuum mobile 1. vrste (stroj se bi večno vrtil)

perpetuum mobile 2. vrste (ne obstaja - stroj bi sam opravljaj neko delo?)

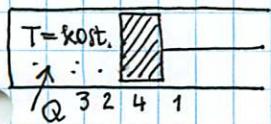




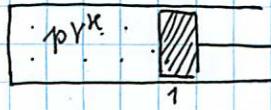
Ko grc fat nazaj ($A > 0$), se olvajaja hladeni zrak, ki ga segreje in ta zrak ponine bat naprej, navorzgor ($A < 0$) ... to je EN cikel.



pV^k (adiabatska spremembra)

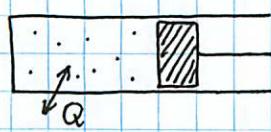


Odstoranim izolacijo in dajamo Q (kopal z vročo vodo).



Bat na kiro raztegnemo na začetno stanje (1).

Skup:



$$\Delta W_n = A + Q = 0$$

$$0 = -A_{kr} + Q_{do} - Q_{dv}$$

$$\eta = \frac{A_{kr}}{Q_{do}} = \frac{Q_{do} - Q_{dv}}{Q_{do}} = 1 - \frac{Q_{dv}}{Q_{do}} < 1$$

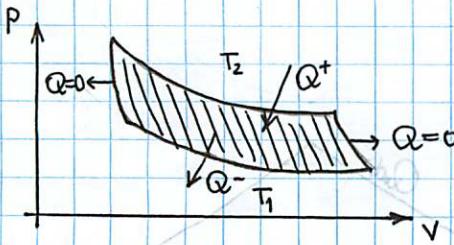


Pomembno: Kolaj dolimo največje delo?
 $\Delta W_n = A + Q = 0$

Qdo ... izgubljena toplota pri hlađenju npr. ☺

Qdo ... tenac, ki ga plačamo, ali pa ogenj ☺

Carnot



$$\Delta S \geq \int \frac{dQ}{T}$$

$$\Delta S = \frac{Q}{T} \quad \text{tch. Oftih nikoli ne dobro}$$

$Q=0$... "bat" je izoliran (na batu premaknemo fat)

$$\Delta S_{\text{Carnot}} = 0 = -\frac{Q_{dv}}{T_1} + 0 + \frac{Q_{do}}{T_2} + 0$$

olvajanje toplote

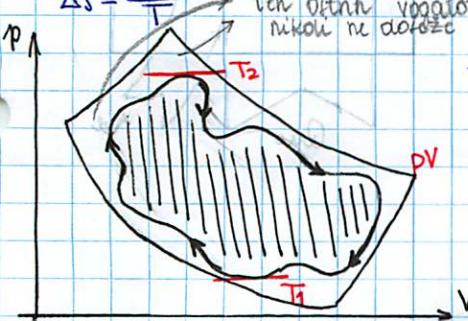
nikoli ne more biti enak 0

$$\frac{Q_{do}}{T_2} = \frac{Q_{dv}}{T_1}$$

$$\frac{Q_{dv}}{Q_{do}} = \frac{T_1}{T_2}$$

$$\eta \leq \eta_c = 1 - \frac{T_1}{T_2} < 1$$

druga oblika
izkoristka idealnega toplotnega stroja
(Carnot)



$T_2 > T_1$

Q_{do} konstantna

Cikel je maksimalen.

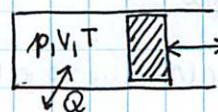
$$\eta \leq \eta_c = 1 - \frac{T_1}{T_2} < 1$$

$$\eta = 1 - \frac{Q_{dv}}{Q_{do}}$$

$$1. \Delta W_n = A + Q$$

$$2. \Delta S \geq \int \frac{dQ}{T}$$

→ Ni izoliran = $\frac{Q}{T}$



A_{kr} je plotečina .
 $\eta = \frac{A_{kr}}{Q_{dov}} \dots$ izkorišček

$$\Delta W_n = A + Q \dots$$

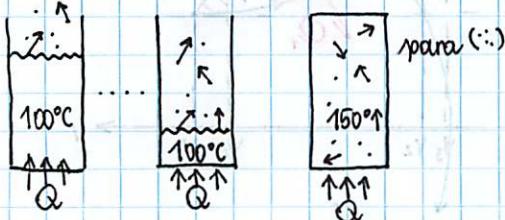
$$Q_v = m c_v \Delta T$$

$$Q_p = m c_p \Delta T$$

$$Q_{takna} = m g t$$

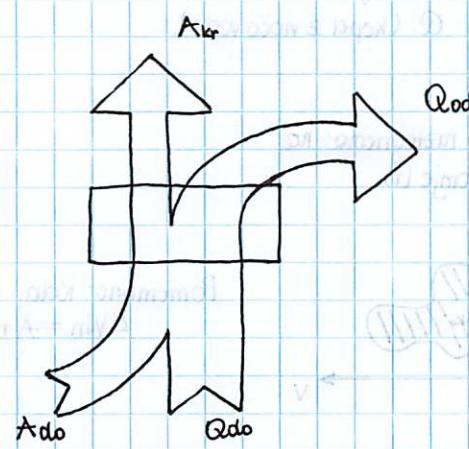
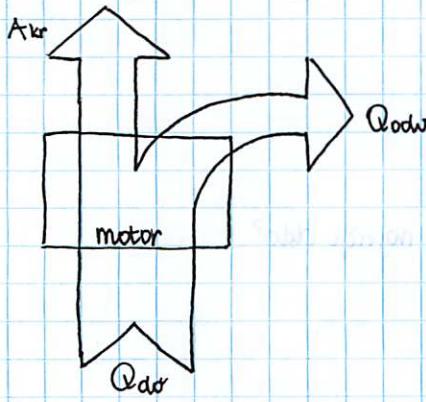
$$Q_{izpanha} = m g_i$$

$$Q_{rezigna} = m g_s$$



$$\Delta S = 0 + \frac{Q_{dov}}{T_{dov}} + 0 - \frac{Q_{odv}}{T_{odv}} = 0$$

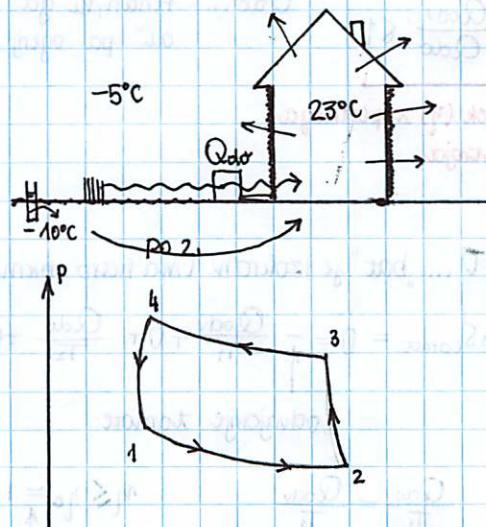
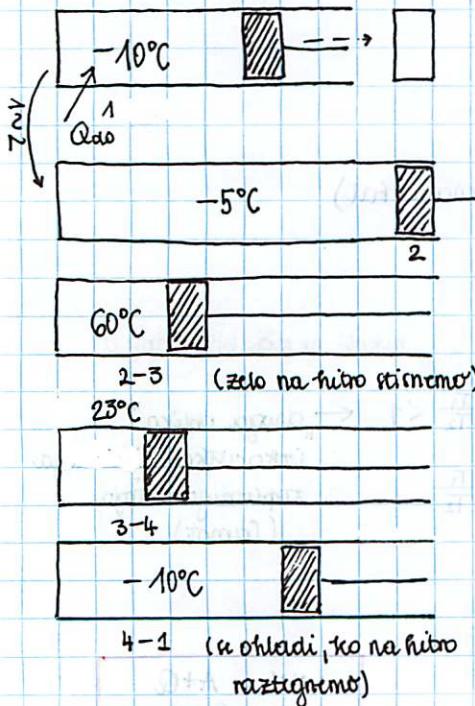
$$\eta = 1 - \frac{Q_{odv}}{Q_{dov}} = 1 - \frac{T_{odv}}{T_{dov}} < 1$$



Specifična topota snovi - tekoča
topota je tako dovedena, da se je
temperatura spremenila zol takto,
tekoč se je spremenila.

HLADILNI STROJ (KLIMATSKA NAPRAVA), HLADILNIK, TOPLOTNA ČRPALKA

(A) Toplotna črpalka



$$\eta = \frac{Q_{dov}}{A_{kr}} > 1$$

definicija (izkoriščka)
za hladilni stroj

$$\eta = \frac{Q_{dov}}{Q_{odv} - Q_{dov}} = \frac{T_{dov}}{T_{odv} - T_{dov}}$$

ker se stroj v enem ciklu vrne
nазaj v isto stanje,
lahko namreco A priremo T.

