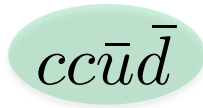


Evidence for a doubly charm tetraquark pole with lattice QCD



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F1, IJS, 24th March 2022

based on **2202.101101**

done in collaboration
with **M. Padmanath**

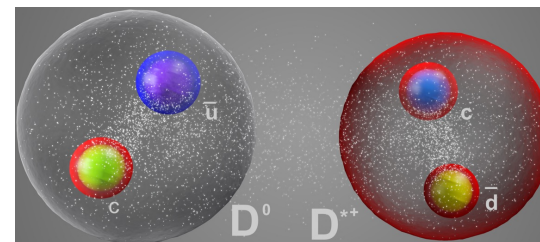


'Doubly charming' tetraquark is the longest-lived exotic-matter particle ever found
By Ben Turner published August 05, 2021

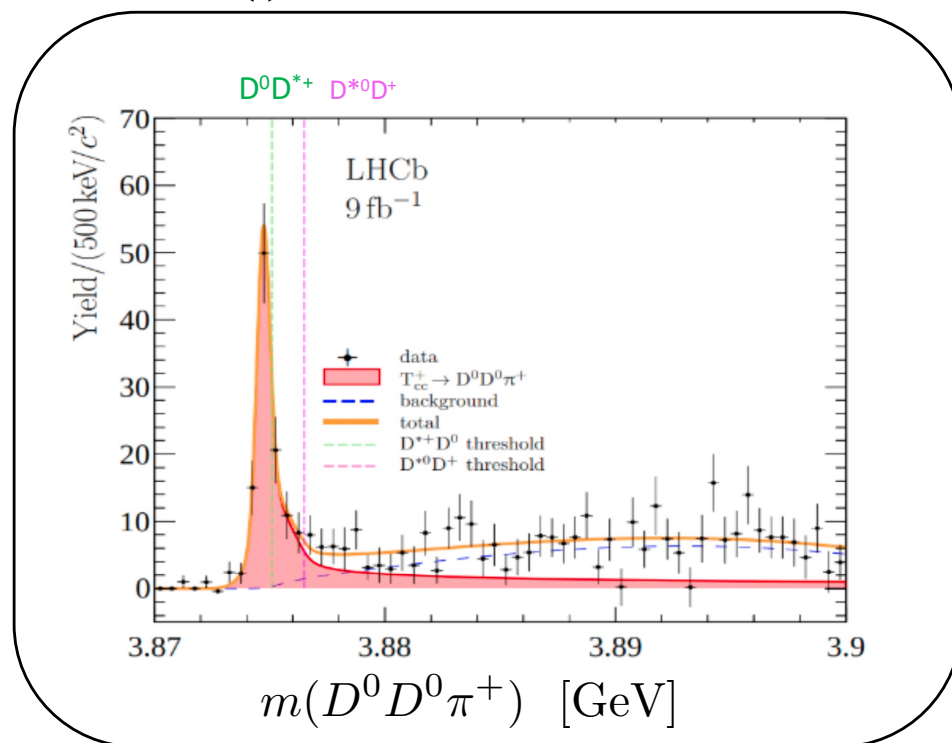
Twice the charm: long-lived exotic particle discovered
CERN Accelerating science

LHCb discovery of T_{cc}

$cc\bar{u}\bar{d}$



$$\delta m \equiv m_{T_{cc}^+} - (m_{D^{*+}} + m_{D^0}) \quad I=0, J^P=1^+$$



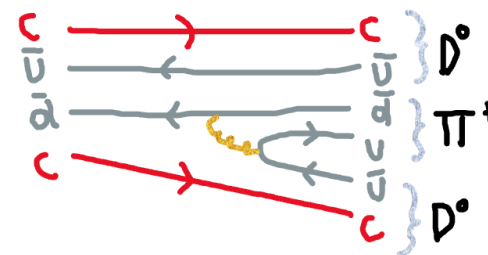
The longest lived discovered hadron with explicitly exotic quark content

LHCb July 2021, 2109.01038, 2109.01056

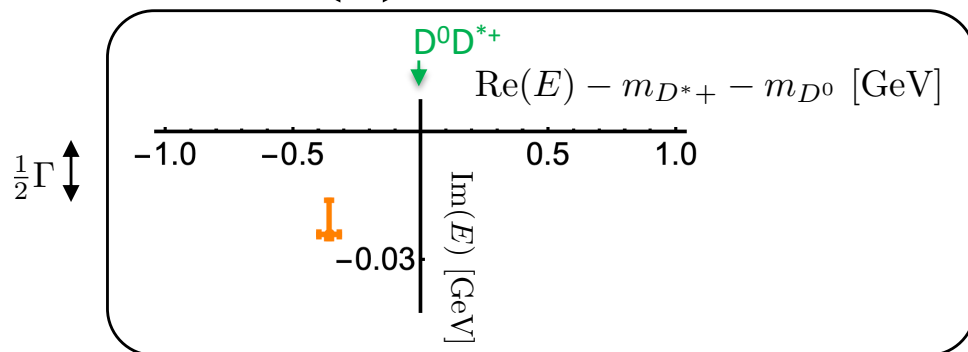
The doubly charmed tetraquark T_{cc}^+ , $I = 0$ and favours $J^P = 1^+$.

No states observed in $D^0 D^+ \pi^+$: eliminates possibility of $I = 1$.

Near-threshold state: Demands pole identification to confirm existence.



Pole in $T(E)$ $\delta m = -0.36(4)$ GeV



$$\begin{aligned} \delta m_{\text{pole}} &= -360 \pm 40_{-0}^{+4} \text{ keV}/c^2, \\ \Gamma_{\text{pole}} &= 48 \pm 2_{-14}^{+0} \text{ keV}, \end{aligned}$$

Theoretical predictions

❖ Phenomenological approaches →

* Janc & Rosina , *Few Body Syst.* 35, 175 (2004), [hep-ph/0405208](#)

one of the most sophisticated quark model predictions:

V_{ij} between all pairs of quarks, ground state energy of four-body problem

$$\delta m = -1.6 \pm 1.0 \text{ MeV}$$

❖ Lattice QCD

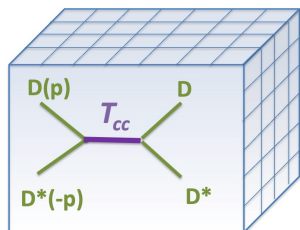
only finite-volume eigen-energy $E_n(L)$ was extracted:

this does not suffice to establish a near-threshold state

[Junnarkar, Mathur, Padmanath, PRD 99, 034507 \(2019\), 1810.12285](#)

[Hadron Spectrum, JHEP 11, 033 \(2017\), 1709.01417](#)

To establish a near-threshold state: pole in $T(E)$ needs to be found:



scattering amplitude $T(E)$

$$T(E) \propto \frac{1}{s - m^2} = \frac{1}{E^2 - m^2}$$

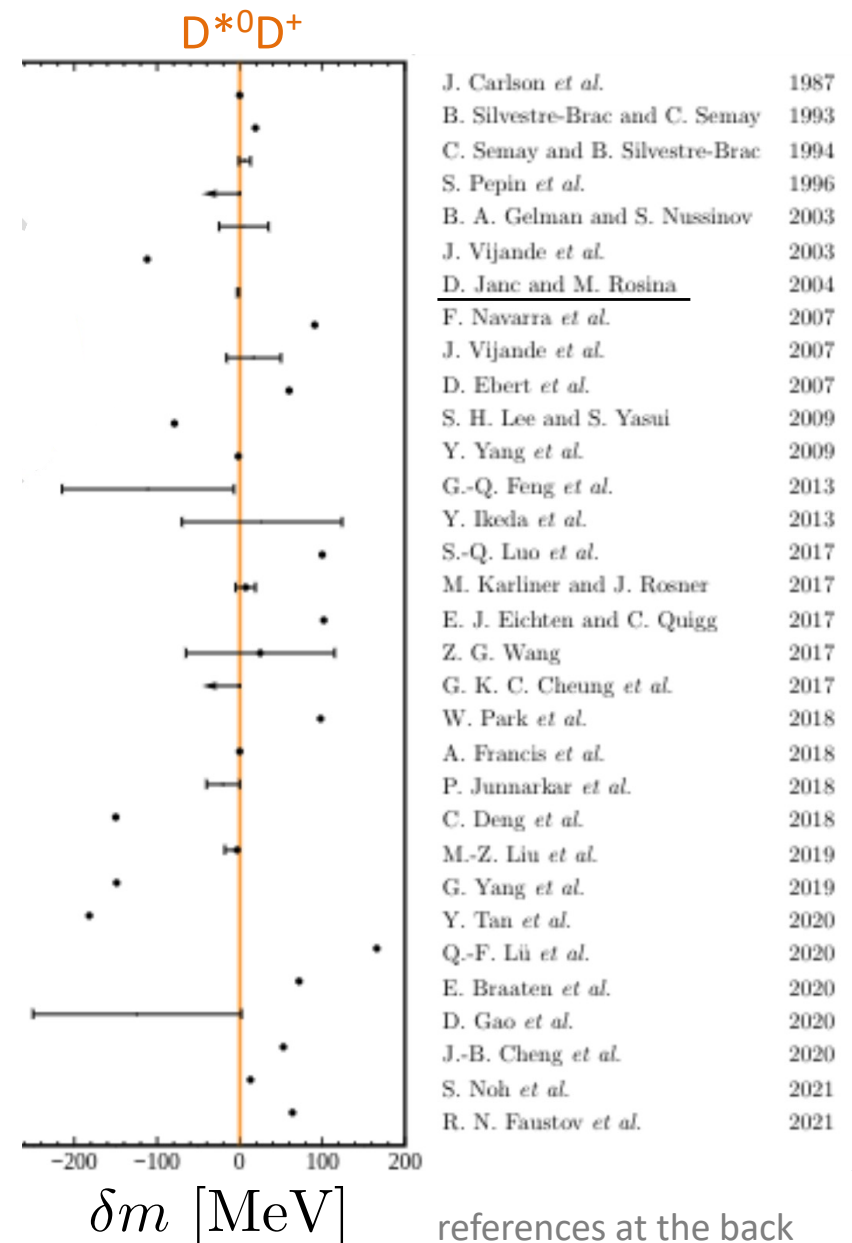
$$E = E_{cm}$$

$T(E)$ has not been extracted by lattice QCD before our study

Our study [2202.101101](#)

first and still the only extraction of $T(E)$ with lattice QCD

pole related to T_{cc} established for the first time with lattice QCD



references at the back

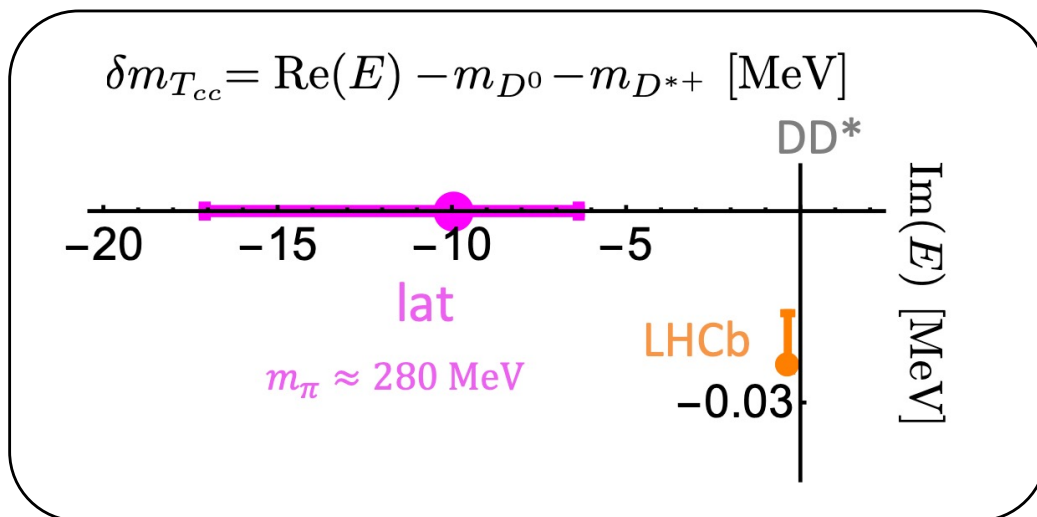
Theoretical PRedictions

courtesy: Ivan Polyakov, EPS-HEP 2021

(references at the back)

Summary of our lattice results

Pole of $T(E)$



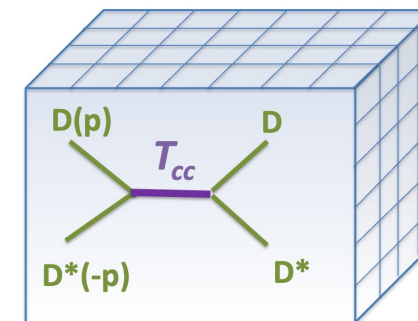
	$\delta m_{T_{cc}} \text{ [MeV]}$	T_{cc}
lat	$-9.9^{+3.6}_{-7.2}$	virtual bound st.
LHCb	$-0.36(4)$	bound st.

omitting $D^0 D^0 \pi^+$

$$m_{u,d} > m_{u,d}^{\text{phy}}$$

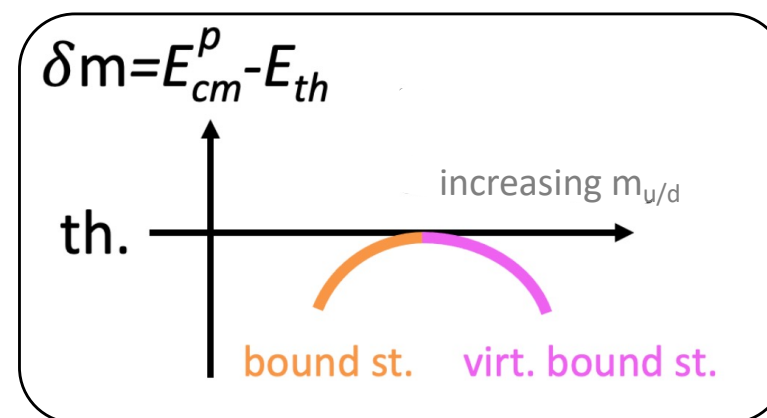
$$m_\pi \approx 280 \text{ MeV}$$

$$m_D \approx 1927 \text{ MeV}$$



- ❖ $T(E)$ extracted via the Luscher's method
- ❖ Evidence for pole related to T_{cc}
- ❖ For $m_{u,d} > m_{u,d}^{\text{phy}}$ one expects decreased attraction
 T_{cc} : bound state become virtual bound state
indeed this is what we find

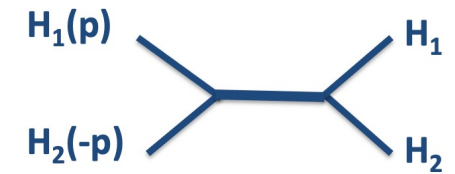
Sketch of expected binding energy



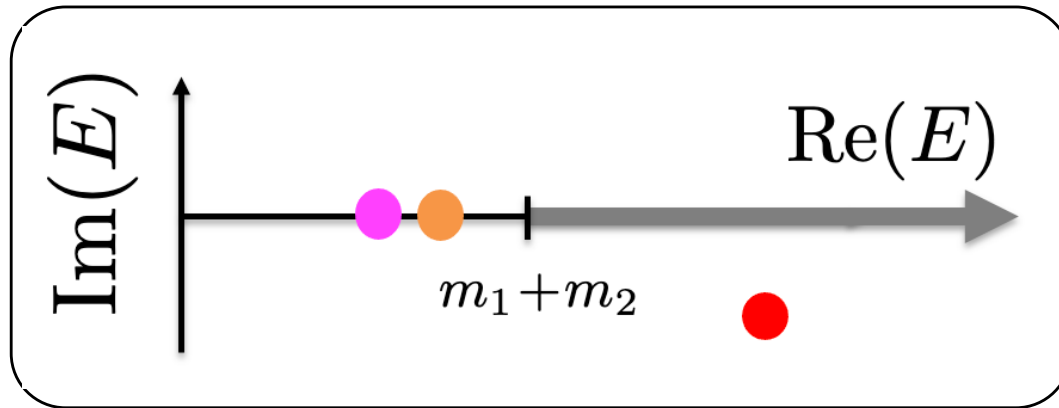
Definitions: bound state, virtual bound state & resonance

$$T(E) \propto \frac{1}{E^2 - m^2}$$

$$T(E) \propto \frac{1}{E^2 - m^2 + iE\Gamma}$$



Poles of $T(E)$, $E=E_{cm}$



Virtual bound st. Bound st. Resonance

$p = -i|p|$ $p = i|p|$

$$E = \sqrt{m_1^2 + p^2} + \sqrt{m_2^2 + (-p)^2} < m_1 + m_2$$

How did we arrive at these lattice QCD result ?

eigen-energies on the lattice $\rightarrow T(E)$

Lattice QCD ensembles employed

CLS Consortium

with dynamical quarks: u,d,s

$m_u = m_d > m_{u,d}^{\text{phy}}$, $m_\pi \approx 280 \text{ MeV}$

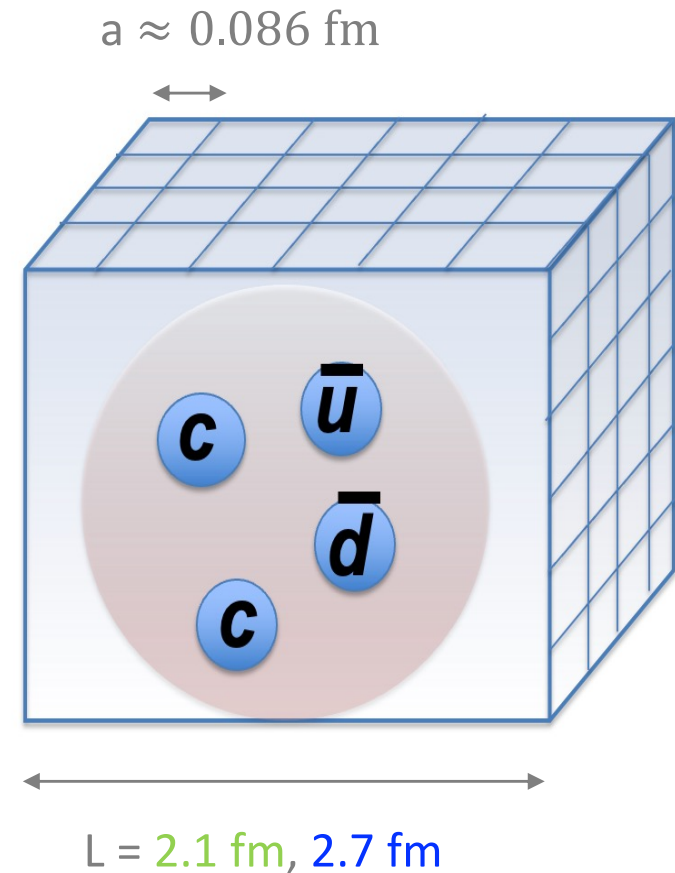
Clover Wilson fermions

$$\langle C \rangle = \int DG Dq D\bar{q} C e^{-S_{QCD}/\hbar}$$

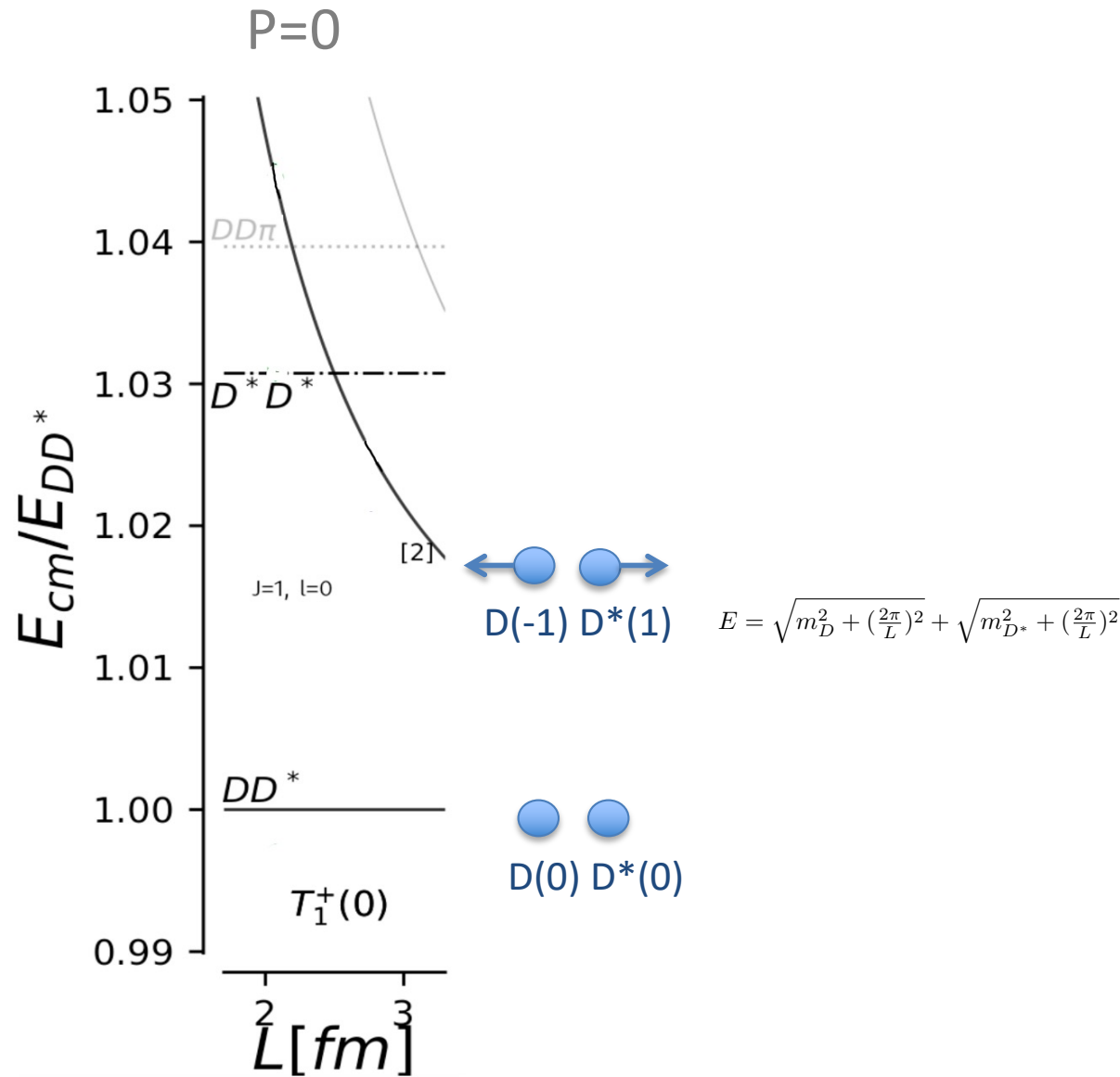
Euclidian space-time

$$S_{QCD}^E = \int d^4x_E \mathcal{L}_{QCD}^E(m_q, g_s)$$

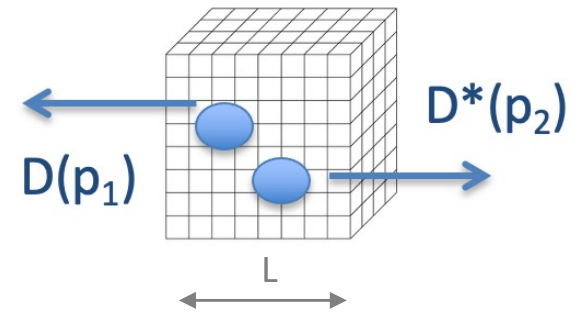
strategy: $C \rightarrow E \rightarrow T(E)$



Energies of DD* in non-interacting limit



$$E_{DD^*} \equiv m_D + m_{D^*}$$

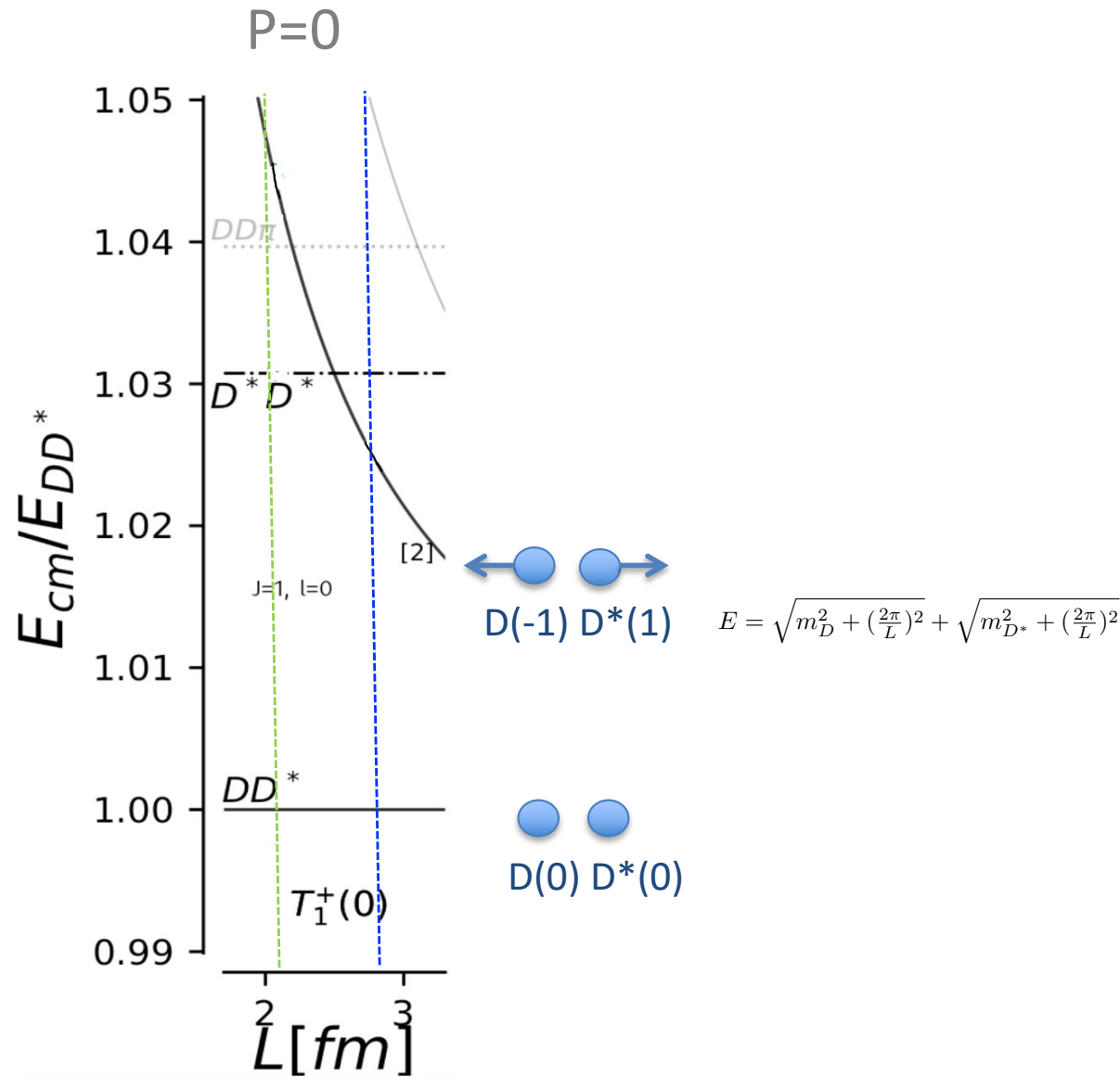


periodic bc in space

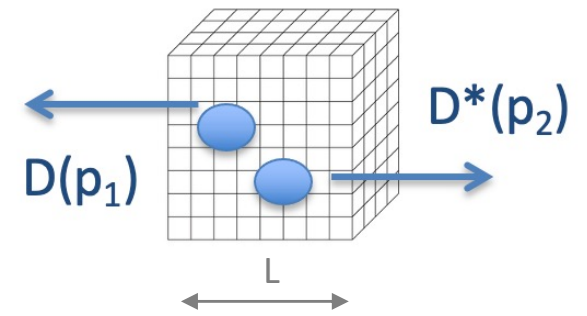
$$\vec{p}_{1,2} = \vec{n}_{1,2} \frac{2\pi}{L}$$

$$E = \sqrt{m_D^2 + \vec{p}_1^2} + \sqrt{m_{D^*}^2 + \vec{p}_2^2}$$

Energies of DD* in non-interacting limit



$$E_{DD^*} \equiv m_D + m_{D^*}$$



$$L = 2.1 \text{ fm}, 2.7 \text{ fm}$$

periodic bc in space

$$\vec{p}_{1,2} = \vec{n}_{1,2} \frac{2\pi}{L}$$

$$E = \sqrt{m_D^2 + \vec{p}_1^2} + \sqrt{m_{D^*}^2 + \vec{p}_2^2}$$

Extracting eigen-energies from correlation functions

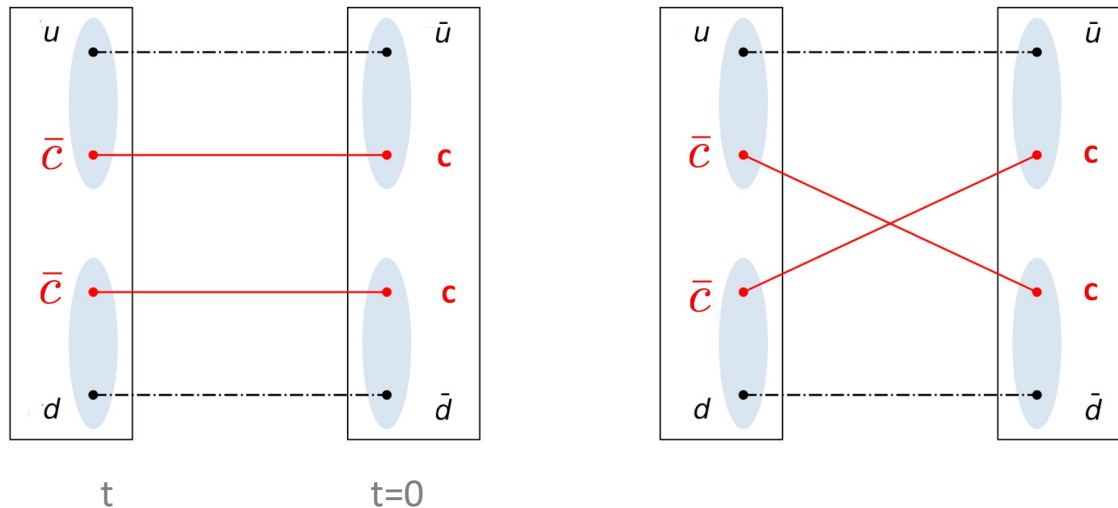
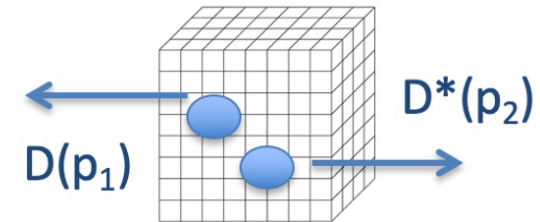
$cc\bar{u}\bar{d}$

$J^P=1^+, I=0$

$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^\dagger(0) | 0 \rangle = \sum_n \langle 0 | \mathcal{O}_i | n \rangle e^{-E_n t} \langle n | \mathcal{O}_j^\dagger | 0 \rangle$$

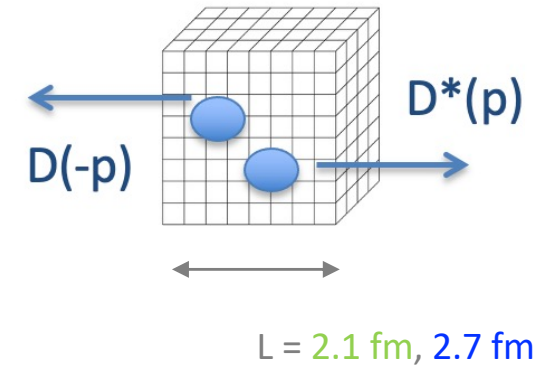
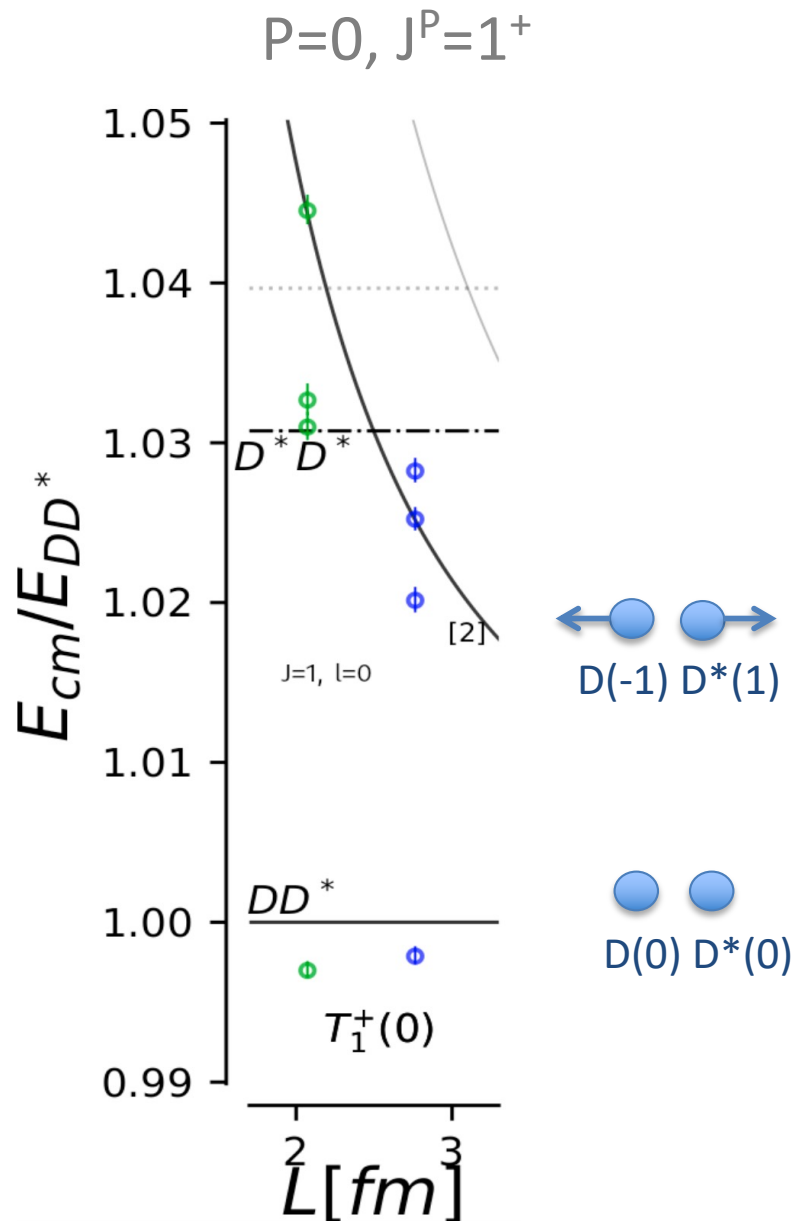
$$\mathcal{O} = \overset{\mathbf{D}}{(\bar{u}\gamma_5 c)_{\vec{p}_1}} \overset{\mathbf{D}^*}{(\bar{d}\gamma_i c)_{\vec{p}_2}} - (\vec{p}_1 \leftrightarrow \vec{p}_2) \quad \vec{p}_{1,2} = \vec{n}_{1,2} \frac{2\pi}{L}$$

$$(\bar{u}\gamma_5 \gamma_t c)_{\vec{p}_1} (\bar{d}\gamma_i \gamma_t c)_{\vec{p}_2}$$



Diquark antidiquark operators $[cc][\bar{u}\bar{d}]$ not incorporated: Emmanuel, Luka

Energies of DD^* from lattice QCD



- ❖ energies shifted from non-interacting energies:
renders info on $T(E)$
- ❖ focus on energy region near DD^* threshold
- ❖ scattering in partial wave $l=2$ negligible

[2]: DD^* with $J^P=1^+$ in $l=0,2$:

both E degenerate in noninteracting limit

Luscher's relation $E \rightarrow T(E)$

Relation between E and $\delta(E)$, $T(E)$: 1D quantum mechanics

$V=0$: outside the region of potential

$$\psi(x) = A \cos(p|x| + \delta) = \begin{cases} A \cos(px + \delta) & x > R \\ A \cos(-px + \delta) & x < -\frac{R}{2} \end{cases}$$

• this form already ensures

$$\psi(L/2) = \psi(-L/2)$$

• the other BC:

$$\psi'(L/2) = \psi'(-L/2)$$

this requires

$$A p \sin(p(\frac{L}{2}) + \delta) = -A p \sin(-p(-\frac{L}{2}) + \delta)$$

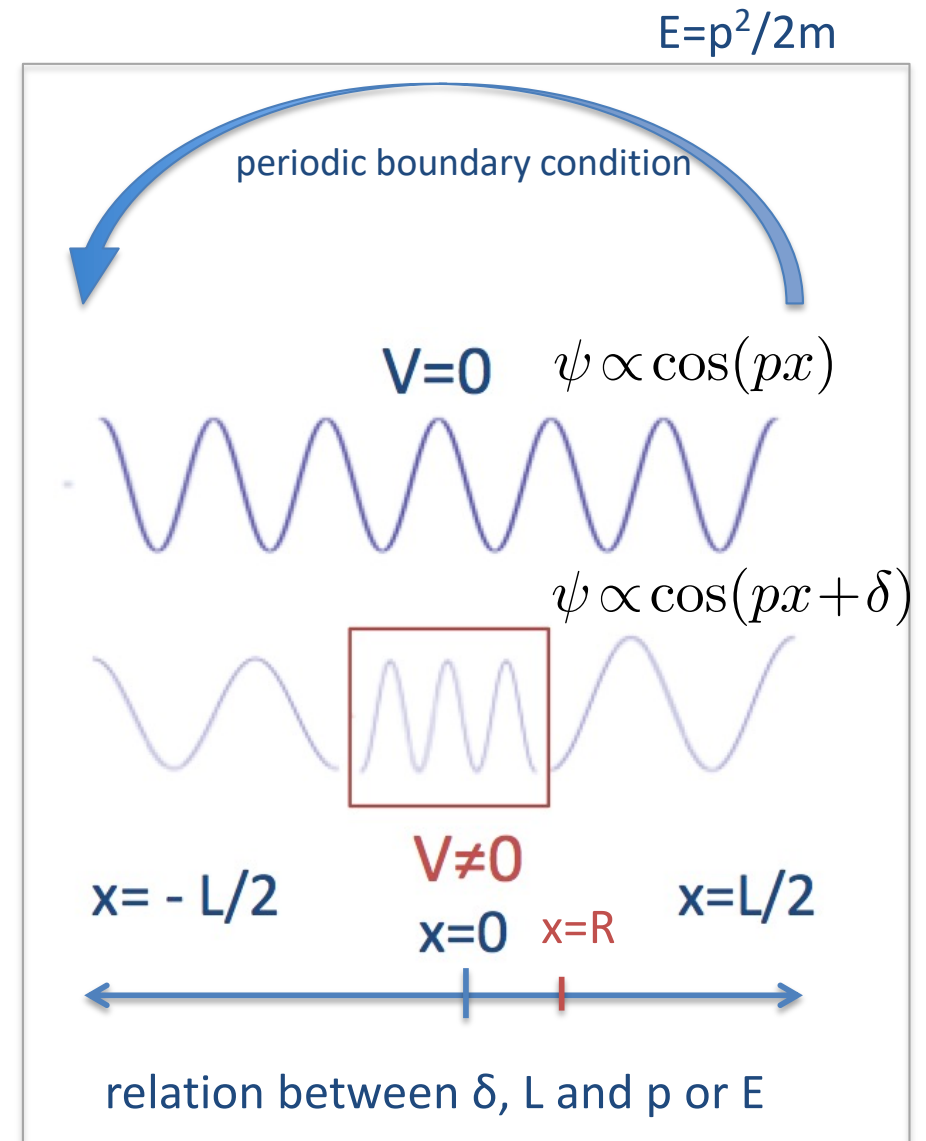
$$\rightarrow \psi'(L/2) = 0, \sin(p\frac{L}{2} + \delta) = 0$$

$$p\frac{L}{2} + \delta = n\pi \quad \boxed{p = n\frac{2\pi}{L} - \frac{2}{L}\delta}$$

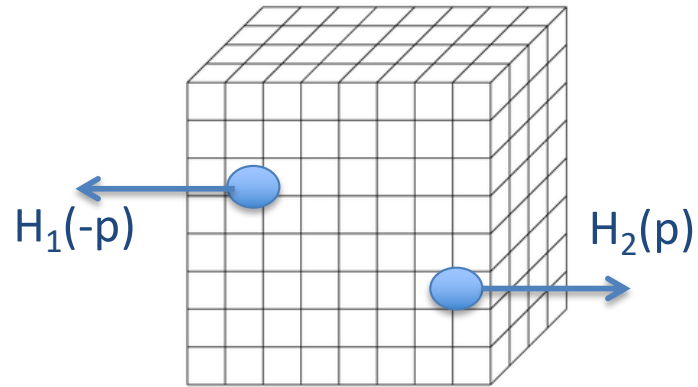
relation between δ, L

$$p = \frac{2\pi}{L}n$$

$$p = \frac{2\pi}{L}n - \frac{2}{L}\delta$$



Relation between E and $\delta(E), T(E)$



E = eigen-energy lattice from lattice in cmf

$$E = \sqrt{m_1^2 + p^2} + \sqrt{m_2^2 + p^2}$$

$$S = 1 + i \frac{4p}{E} T = e^{2i\delta}$$

$$T = \frac{E}{2} \frac{1}{p \cot \delta - ip}$$

Luscher's relation :

$$E \rightarrow T(E), \delta(E)$$

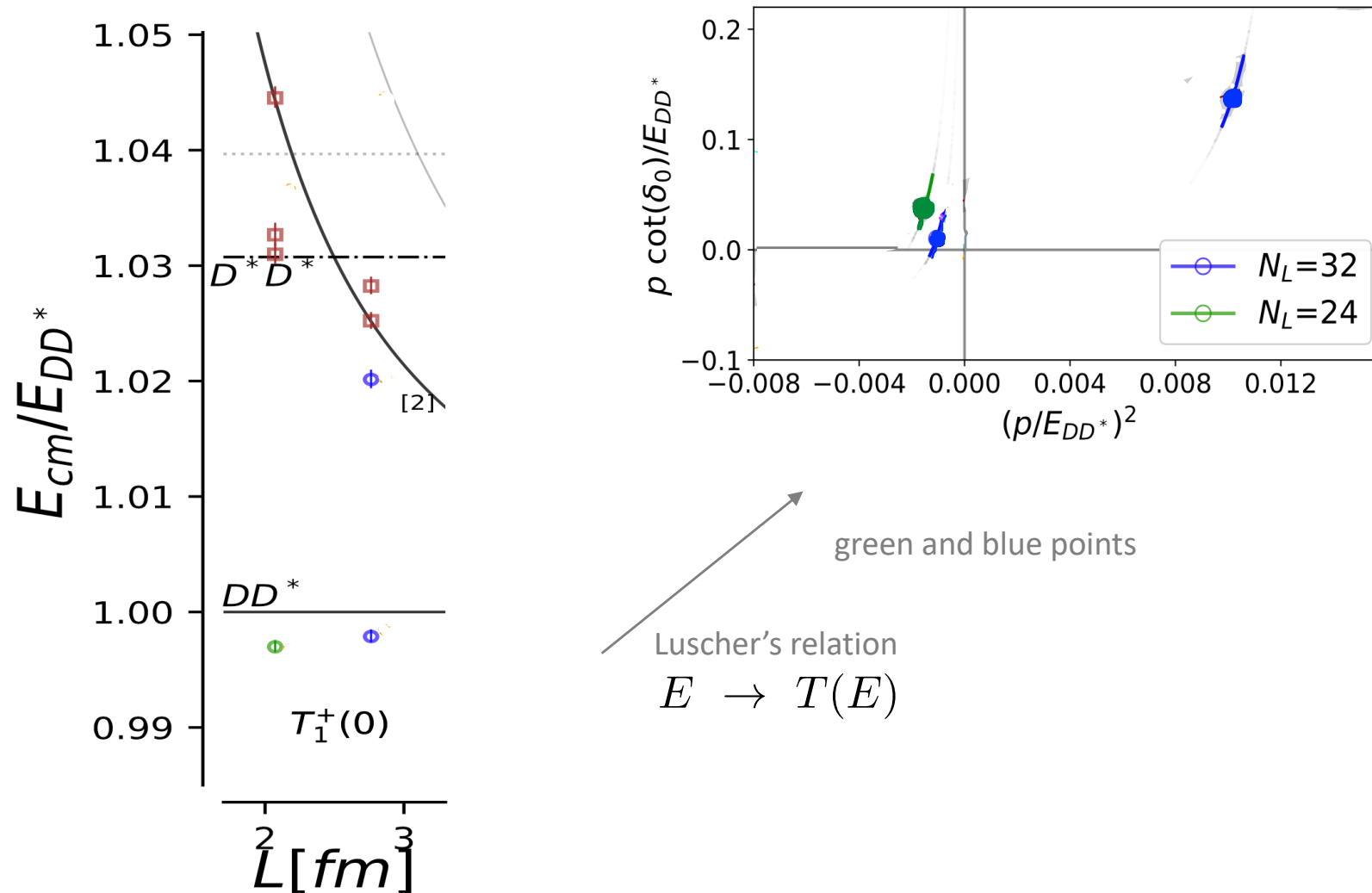
known function
↓

$$l=0: \quad p \cot \delta(p) = \frac{2 Z_{00}(1, (\frac{pL}{2\pi})^2)}{\sqrt{\pi} L}$$

DD* scattering amplitude with $l=0$

fit using just $P=0$

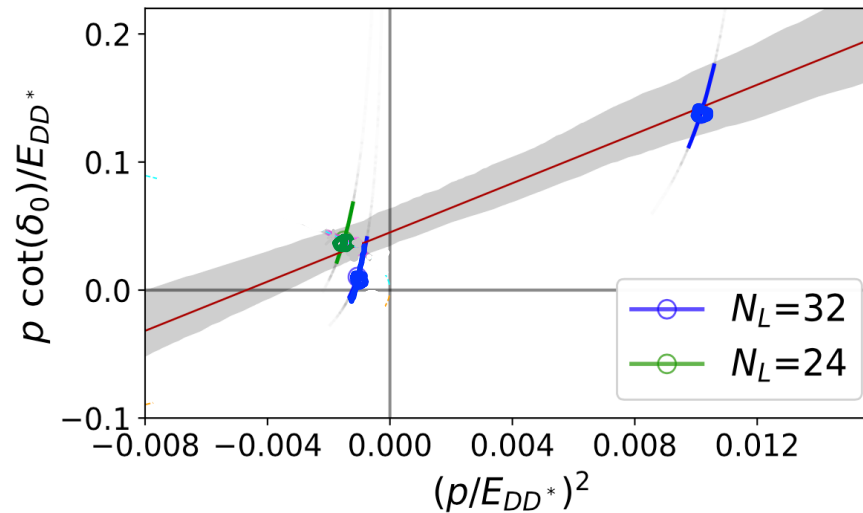
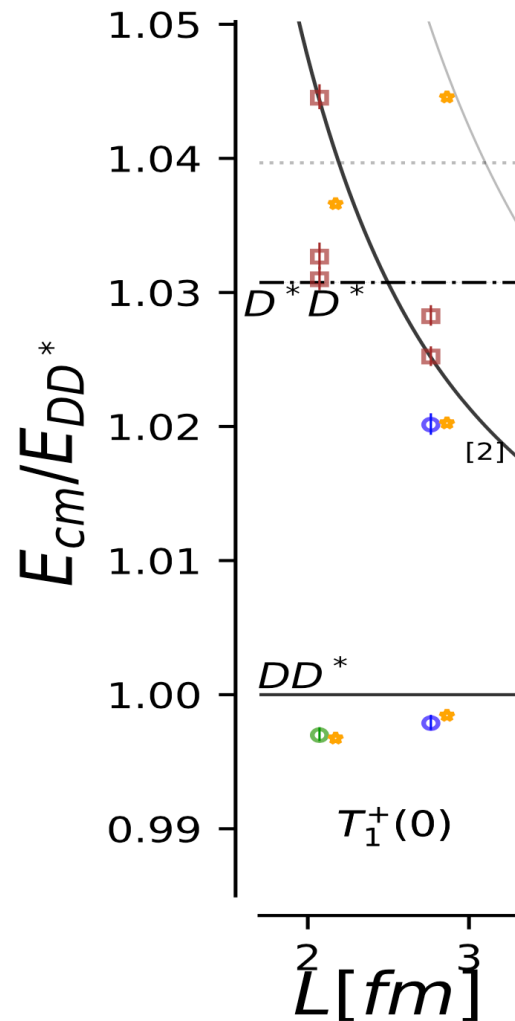
$P=0, J^P=1^+$



DD* scattering amplitude with $l=0$

fit using just $P=0$

$P=0, J^P=1^+$



$$T_{l=0}^{(J=1)} : p \cot \delta_0 = \frac{1}{a_0^{(1)}} + \frac{r_0^{(1)} p^2}{2}$$

Fit parameters:

$$a_0^{(1)} = 1.10^{(+0.34}_{-0.31)} \text{ fm} \quad r_0^{(1)} = 0.95^{(+0.24}_{-0.22)} \text{ fm}$$

• E_{analytic} (fit params)

red line

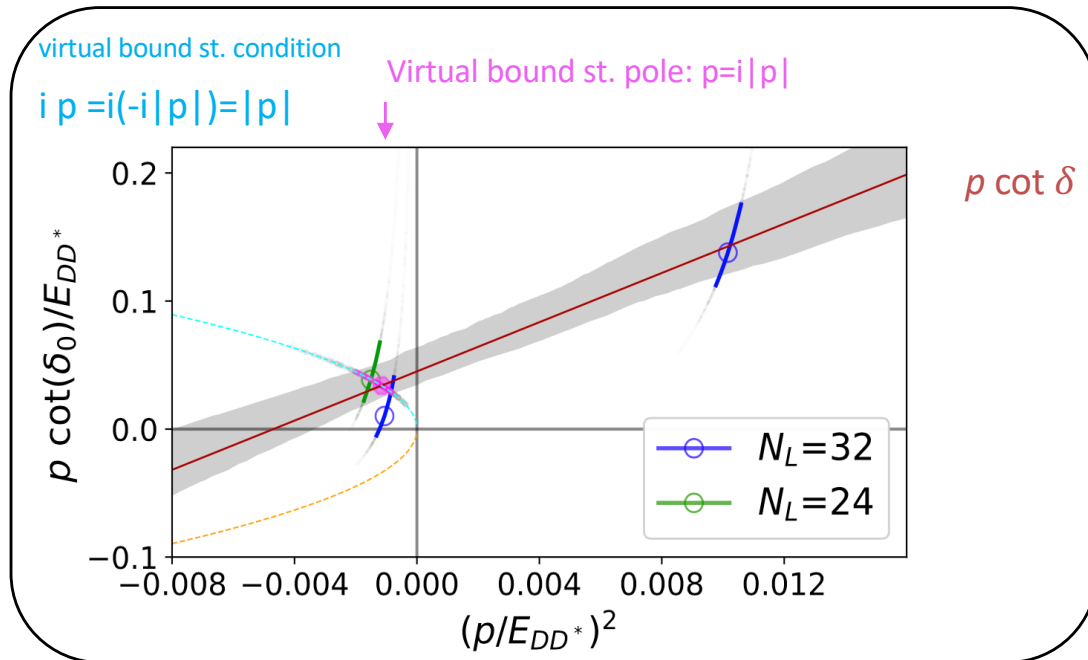
effective range

expansion near threshold

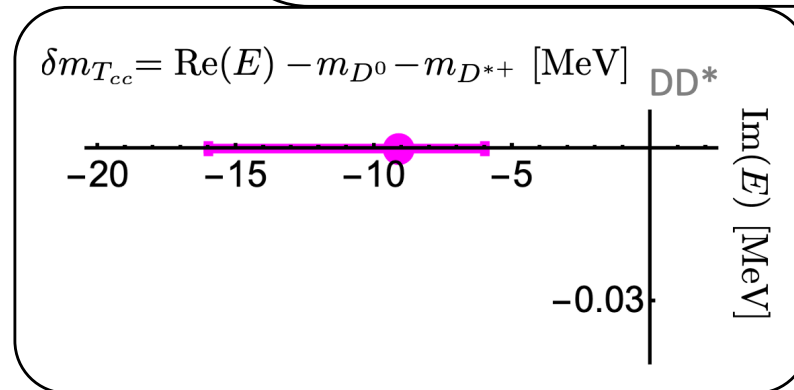
DD* scattering amplitude with $l=0$

fit using just $P=0$

$$T = \frac{E}{2} \frac{1}{p \cot \delta - ip}$$



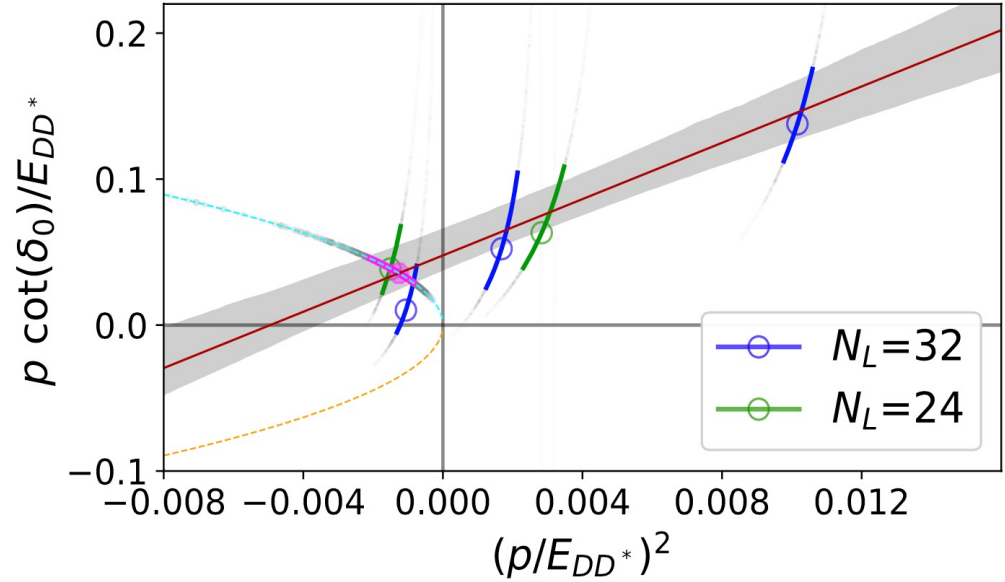
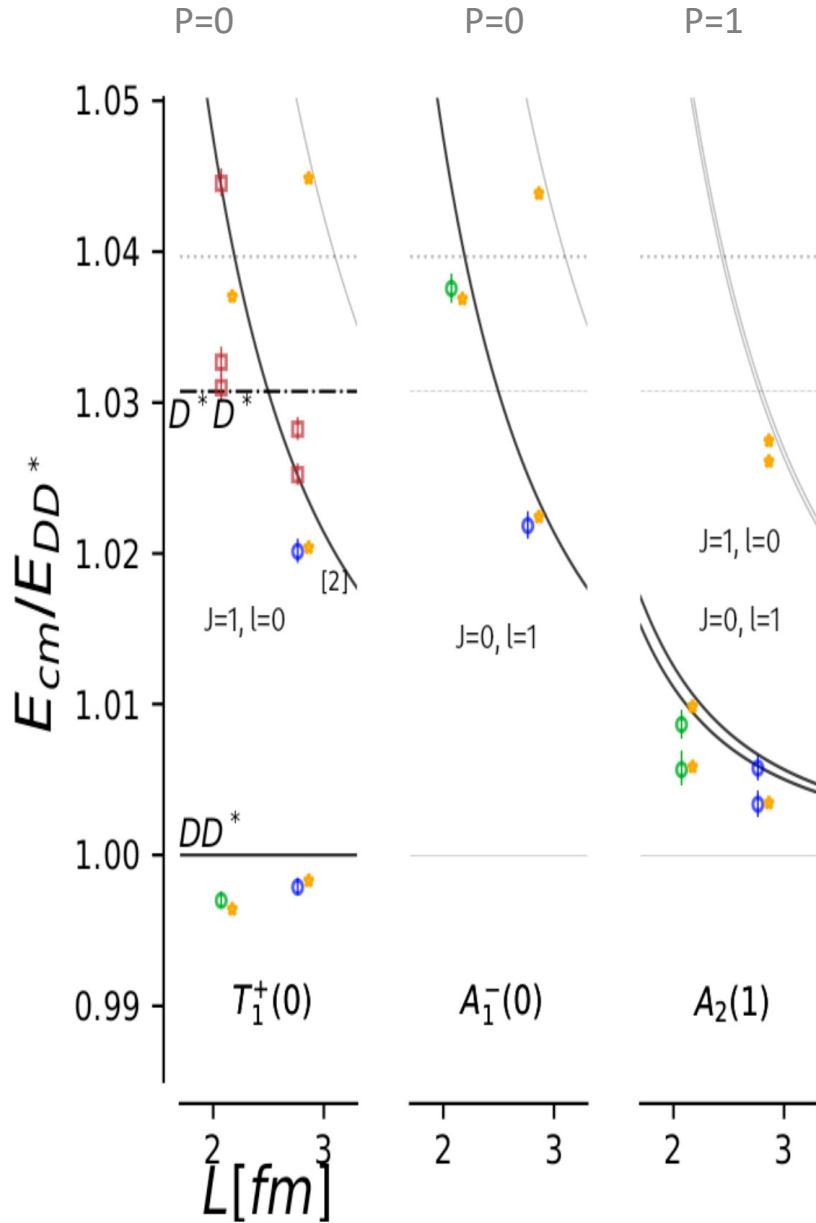
virtual bound st.
 pole of $T(E)$



✿ Binding energy:

$$\delta m_{T_{cc}} = -9.1^{(+3.1)}_{(-6.9)} \text{ MeV.}$$

DD* scattering amplitude with $l=0,1$



$$p^{2l+1} \cot \delta_l^{(J)} = \frac{1}{a_l^{(J)}} + \frac{r_l^{(J)}}{2} p^2$$

Fit parameters:

$$a_0^{(1)} = 1.04(0.29) \text{ fm} \ \& \ r_0^{(1)} = 0.96({}^{+0.18}_{-0.20}) \text{ fm}$$

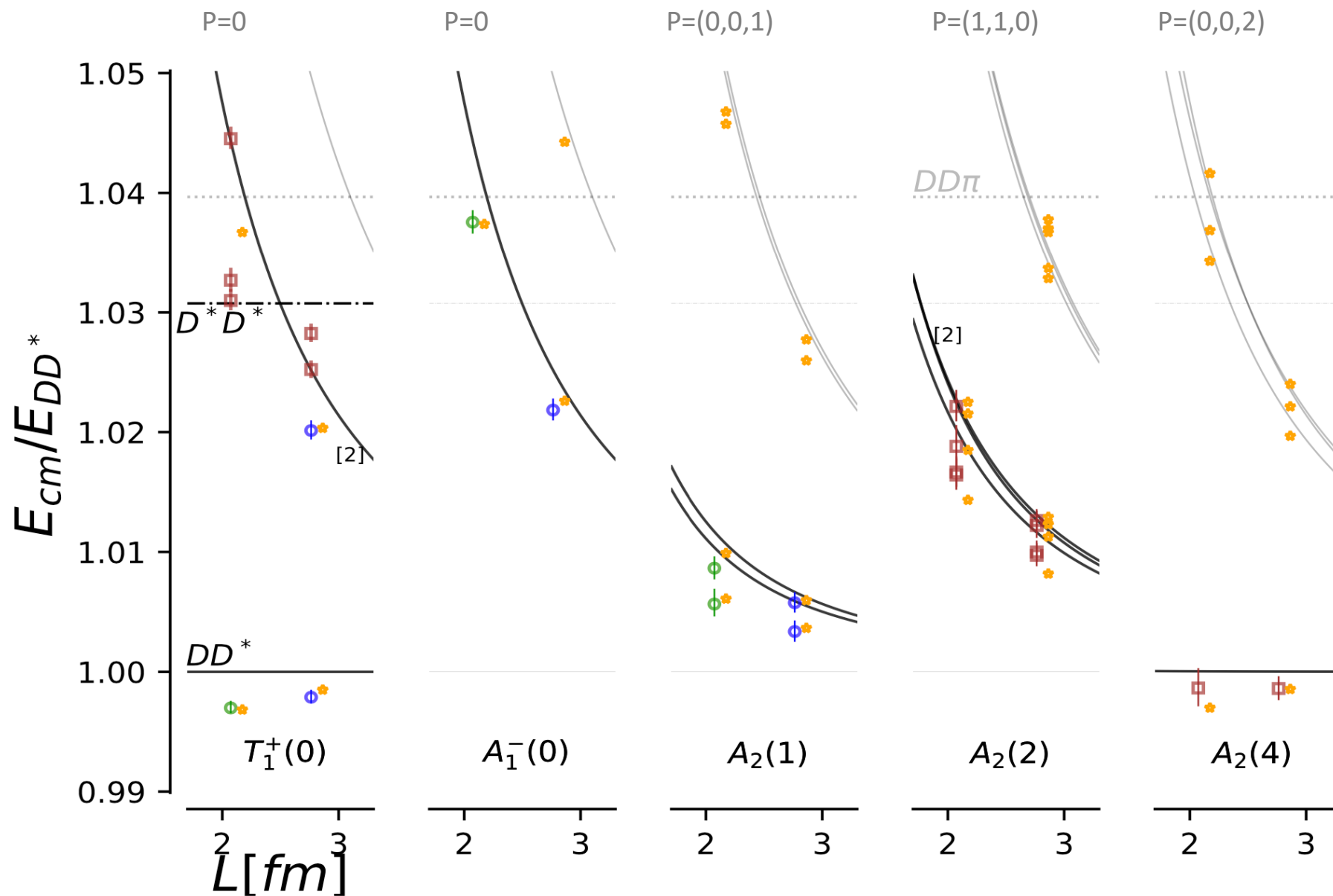
$$a_1^{(0)} = 0.076({}^{+0.008}_{-0.009}) \text{ fm}^3 \ \& \ r_1^{(0)} = 6.9(2.1) \text{ fm}^{-1}$$

Binding energy:

$$\delta m_{T_{cc}} = -9.9({}^{+3.6}_{-7.2}) \text{ MeV}.$$

Predicting eigen-energies

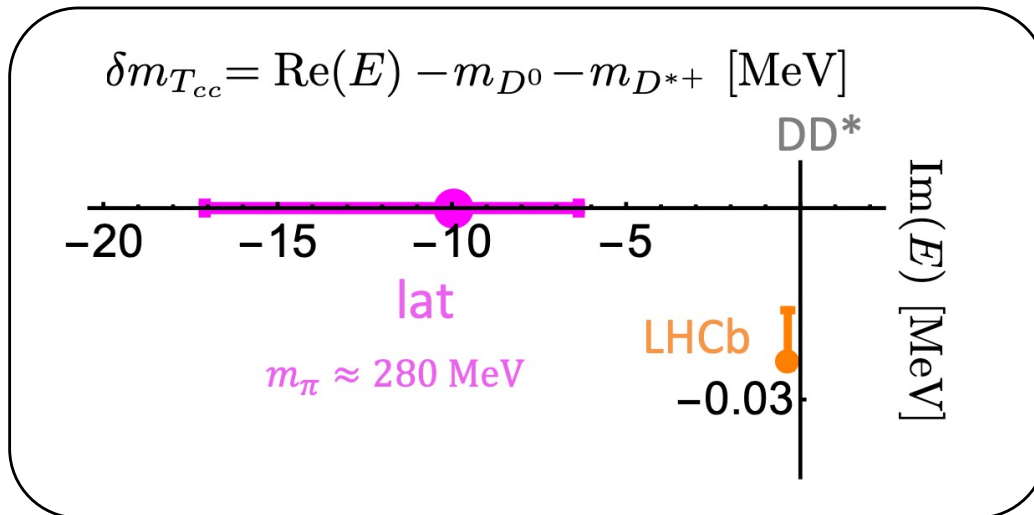
■ E^{analytic} (fit params)



Location of T_{cc} pole

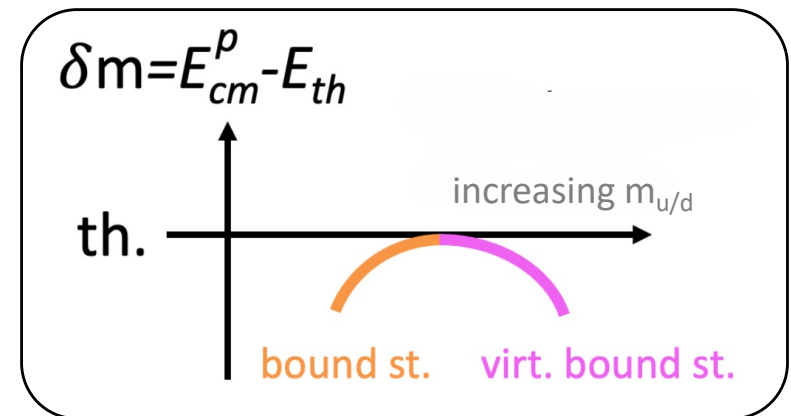
$$\begin{aligned} m_\pi &\approx 280 \text{ MeV} \\ m_D &\approx 1927 \text{ MeV} \\ m_{D^*} &\approx 2049 \text{ MeV} \end{aligned}$$

Pole of $T(E)$



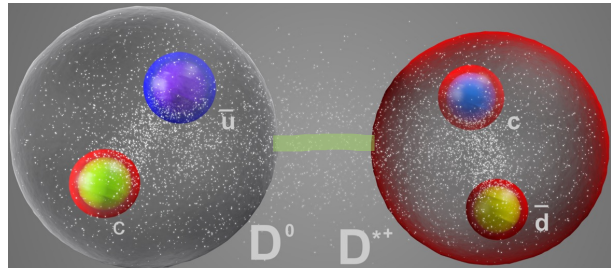
- ❖ lat: evidence for virtual bound state pole
- ❖ we expect this pole is related to T_{cc} bound state pole found by LHCb : arguments in Supplement of [2202.101101](#) and in the following slides

Sketch of expected binding energy



	$\delta m_{T_{cc}} \text{ [MeV]}$	T_{cc}
lat	$-9.9^{+3.6}_{-7.2}$	virtual bound st.
LHCb	$-0.36(4)$	bound st.

Expected dependence of T_{cc} on $m_{u/d}$: simple QM arguments



exchanged particles:
light mesons $\pi, \rho, \dots$

increasing $m_{u/d}$

increasing m_{ex}

decreasing attraction $|V|$

Yukawa-like potential

$$V(r) \propto -\frac{e^{-m_{ex}r}}{r}$$

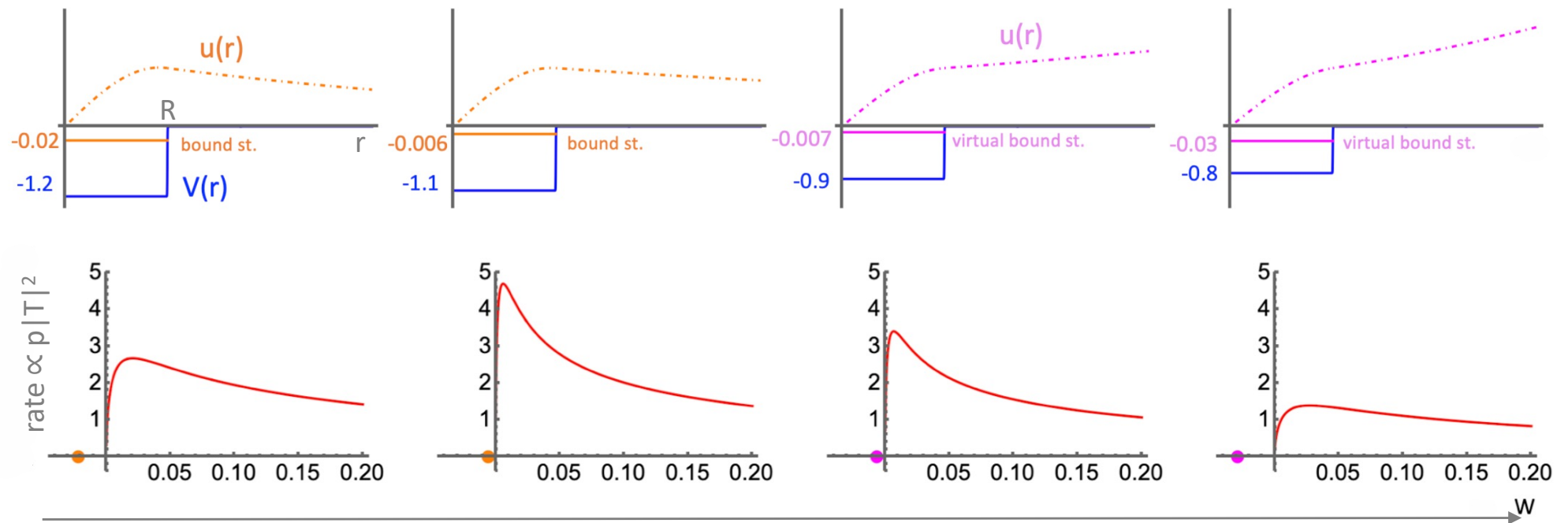
Simplest Example: scattering in square-well potential in QM

$$\delta_l = \arctan[\tan(qR)\frac{p}{q}] - pR$$

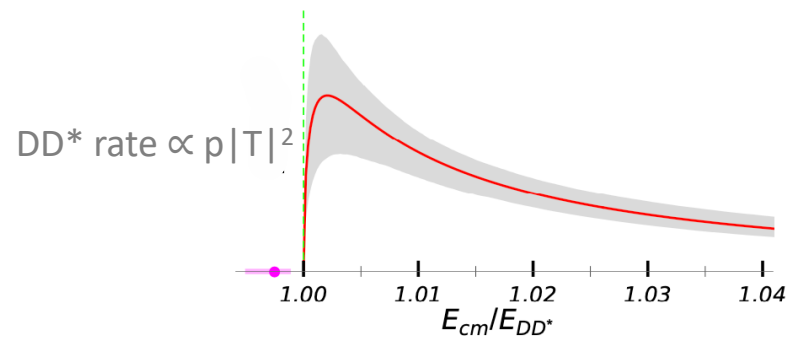
$$u(r) = A \sin(qr) \quad u(r) = B \sin(pr + \delta)$$

$$p=i|p| \quad e^{ipr} = e^{-|p|r}$$

$$p=-i|p| \quad e^{ipr} = e^{|p|r}$$



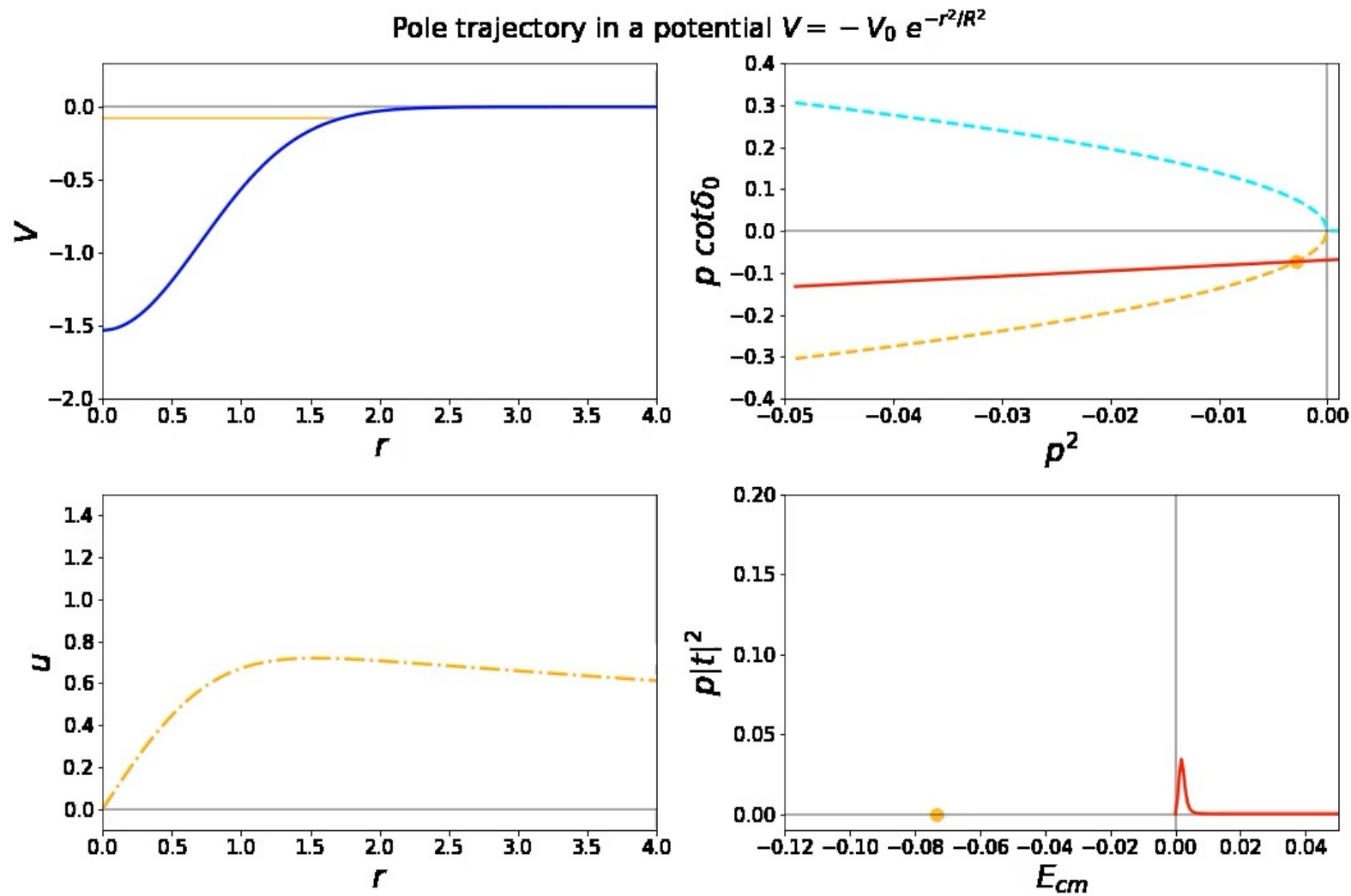
increasing $m_{u/d}$, decreasing attraction



for T_{cc} extracted on lattice

All fully attractive potentials lead to analogous conclusions

video: courtesy M. Padmanath



Dependence on the charm quark mass

simulation at two charm-quark masses

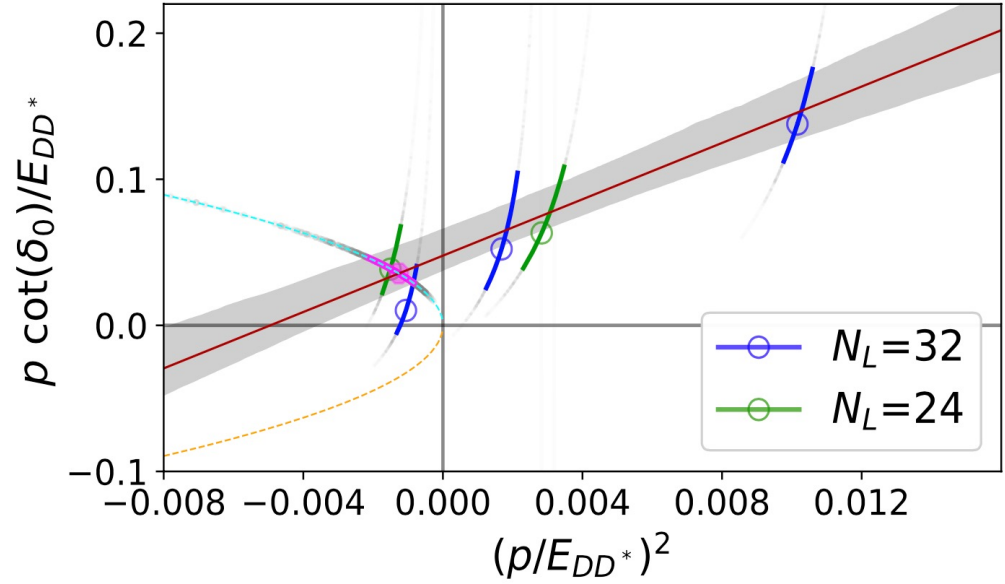
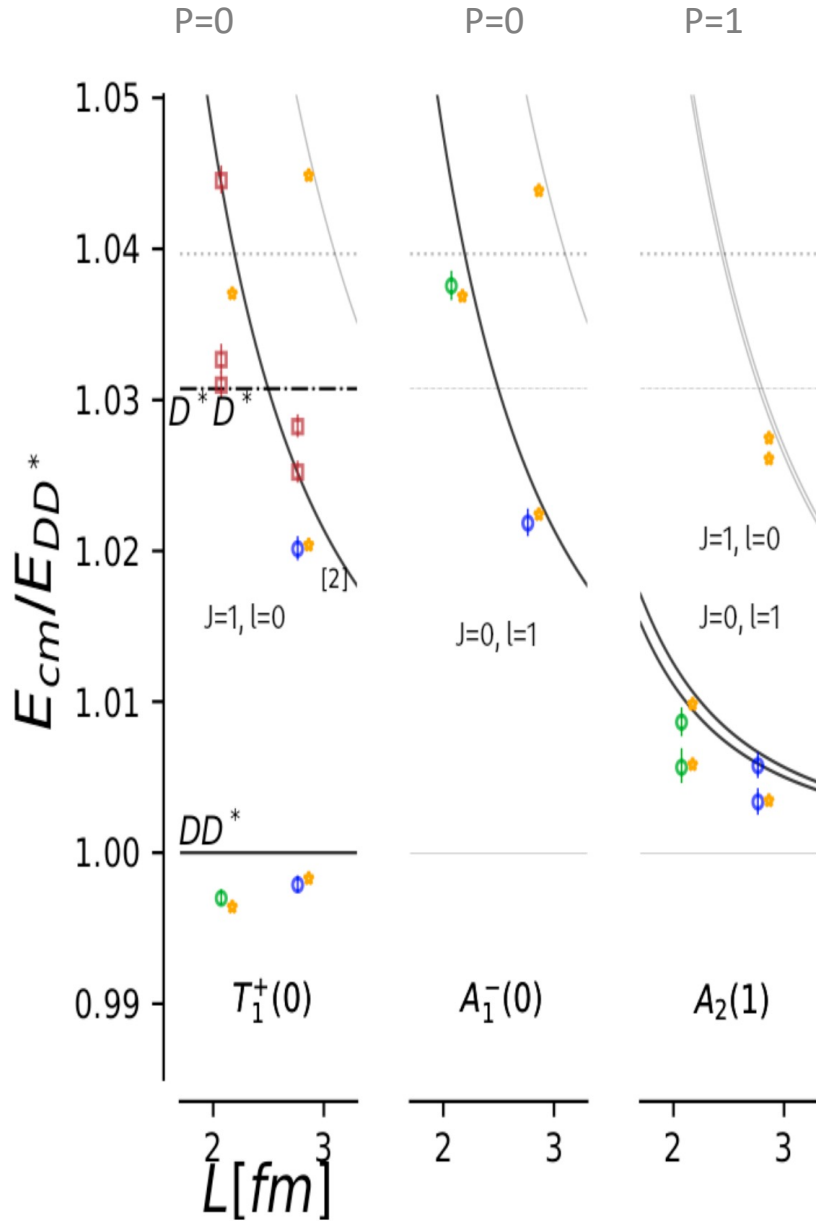
$$M_{av} \equiv \frac{1}{4}(m_{\eta_c} + 3m_{J/\psi})$$

closer to physical
(presented till now)

	m_D [MeV]	m_{D^*} [MeV]	M_{av} [MeV]
lat. ($m_\pi \simeq 280$ MeV, $m_c^{(h)}$)	1927(1)	2049(2)	3103(3)
lat. ($m_\pi \simeq 280$ MeV, $m_c^{(l)}$)	1762(1)	1898(2)	2820(3)
exp. [2, 37]	1864.85(5)	2010.26(5)	3068.6(1)

DD* scattering amplitude with $l=0,1$

at $m_D \approx 1927 \text{ MeV}$



$$p^{2l+1} \cot \delta_l^{(J)} = \frac{1}{a_l^{(J)}} + \frac{r_l^{(J)}}{2} p^2$$



Fit parameters:

$$a_0^{(1)} = 1.04(0.29) \text{ fm} \ \& \ r_0^{(1)} = 0.96({}^{+0.18}_{-0.20}) \text{ fm}$$

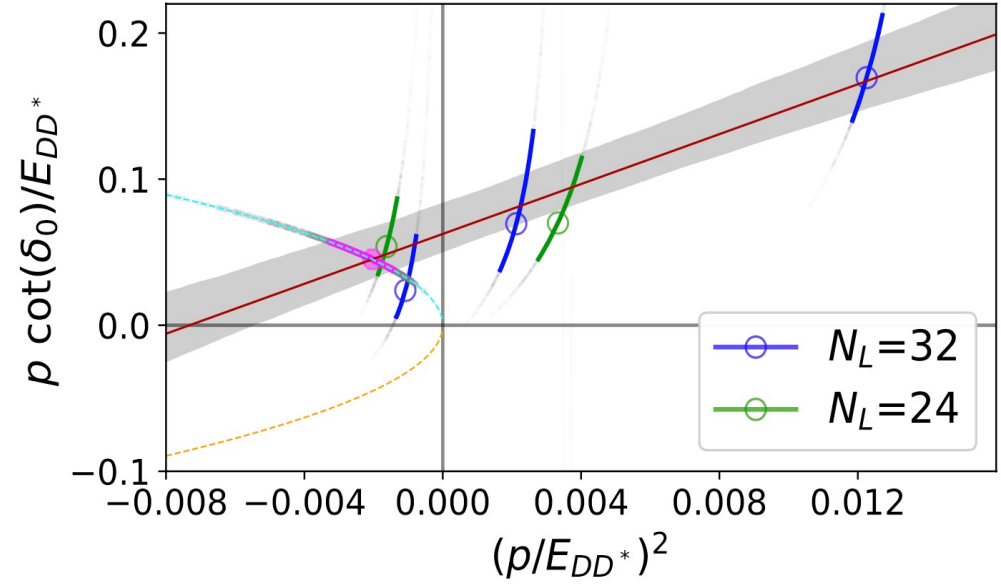
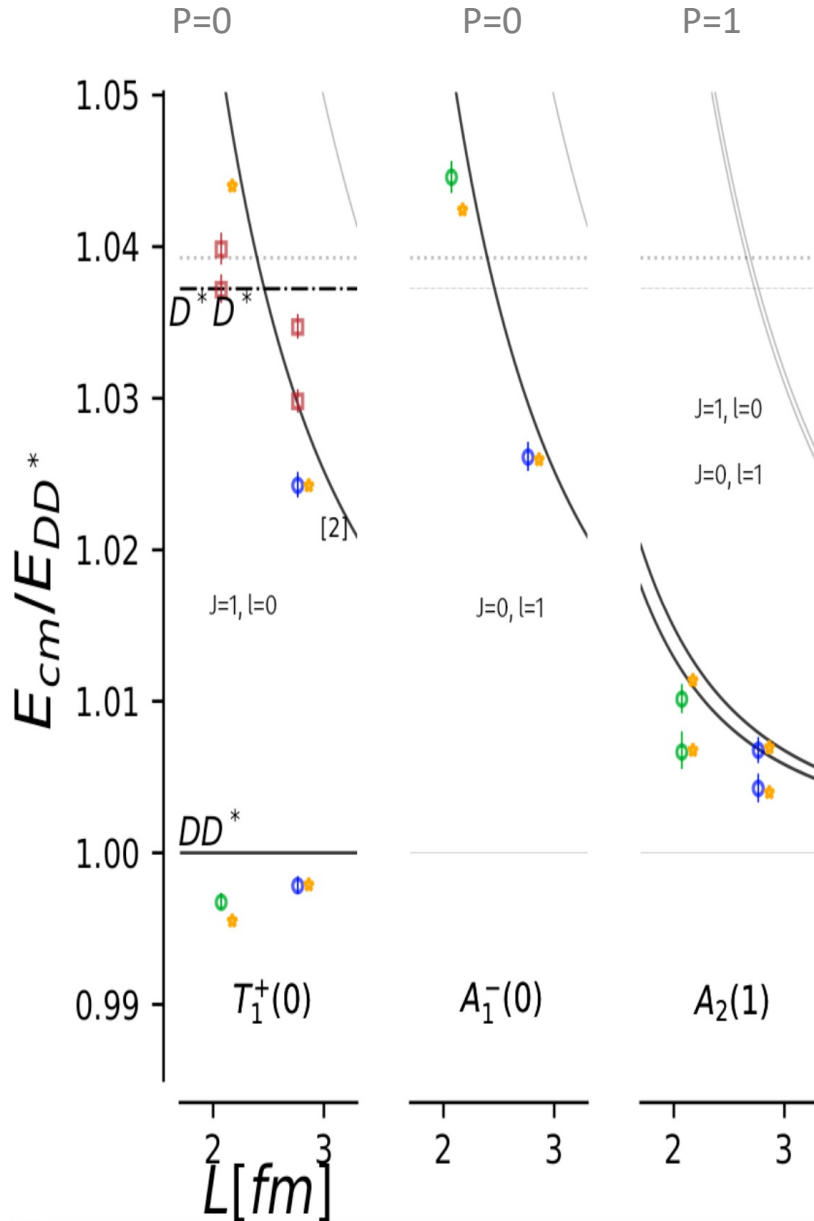
$$a_1^{(0)} = 0.076({}^{+0.008}_{-0.009}) \text{ fm}^3 \ \& \ r_1^{(0)} = 6.9(2.1) \text{ fm}^{-1}$$



Binding energy:

$$\delta m_{T_{cc}} = -9.9({}^{+3.6}_{-7.2}) \text{ MeV}.$$

DD* scattering amplitude with $l=0,1$ at $m_D \approx 1762$ MeV (lighter charm quark mass)



$$p^{2l+1} \cot \delta_l^{(J)} = \frac{1}{a_l^{(J)}} + \frac{r_l^{(J)}}{2} p^2$$

Fit parameters:

$$a_0^{(1)} = 0.86(0.22) \text{ fm} \text{ \& } r_0^{(1)} = 0.92({}^{+0.17}_{-0.19}) \text{ fm}$$

$$a_1^{(0)} = 0.117({}^{+0.013}_{-0.014}) \text{ fm}^3 \text{ \& } r_1^{(0)} = 8.6({}^{+1.5}_{-1.1}) \text{ fm}^{-1}$$

Binding energy:

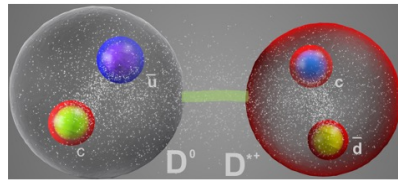
$$\delta m_{T_{cc}} = -15.0({}^{+4.6}_{-9.3}) \text{ MeV.}$$

Lattice results at two m_c

	m_D [MeV]	m_{D^*} [MeV]	M_{av} [MeV]	$a_{l=0}^{(J=1)}$ [fm]	$r_{l=0}^{(J=1)}$ [fm]	$\delta m_{T_{cc}}$ [MeV]	T_{cc}
lat. ($m_\pi \simeq 280$ MeV, $m_c^{(h)}$)	1927(1)	2049(2)	3103(3)	1.04(29)	$0.96^{(+0.18)}_{(-0.20)}$	$-9.9^{+3.6}_{-7.2}$	virtual bound st.
lat. ($m_\pi \simeq 280$ MeV, $m_c^{(l)}$)	1762(1)	1898(2)	2820(3)	0.86(0.22)	$0.92^{(+0.17)}_{(-0.19)}$	$-15.0^{(+4.6)}_{(-9.3)}$	virtual bound st.
exp. [2, 37]	1864.85(5)	2010.26(5)	3068.6(1)	-7.15(51)	$[-11.9(16.9), 0]$	-0.36(4)	bound st.

Observed m_c dependence in agreement with QM arguments for fully attractive potential

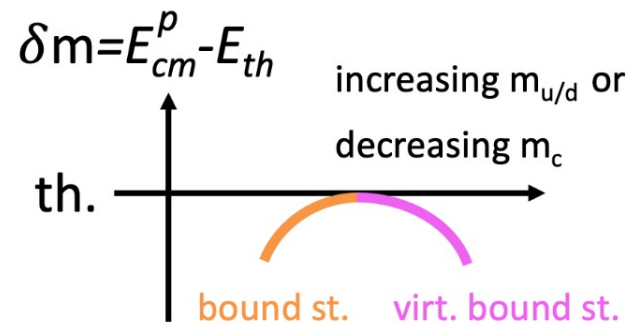
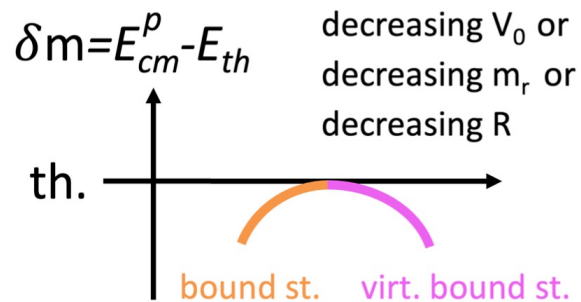
$$V(r) = -V_0 f(r/R)$$



exchanged particles:
light mesons $\pi, \rho, \dots$

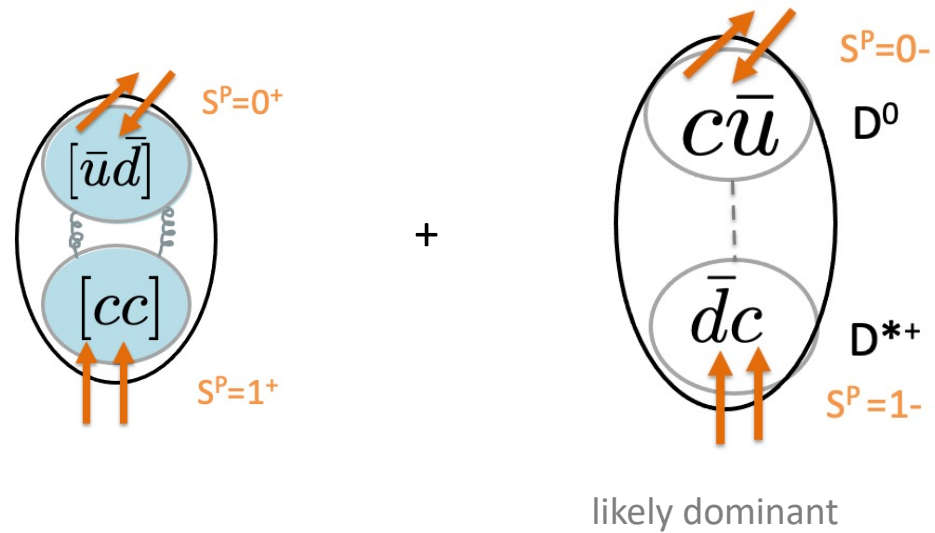
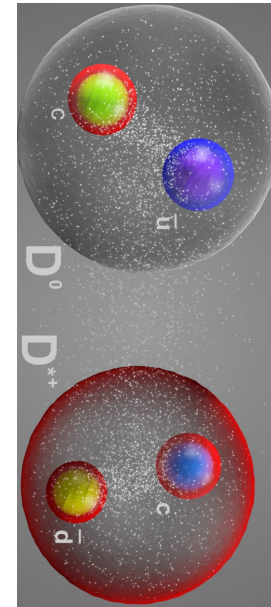
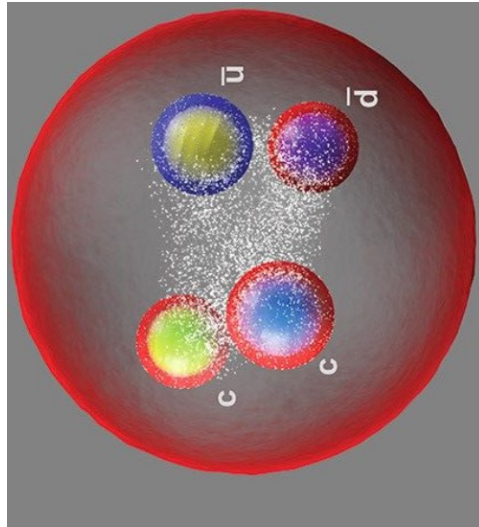
$V(r)$ independent on m_c ,

reduced mass m_r of D, D^* system increases with m_c

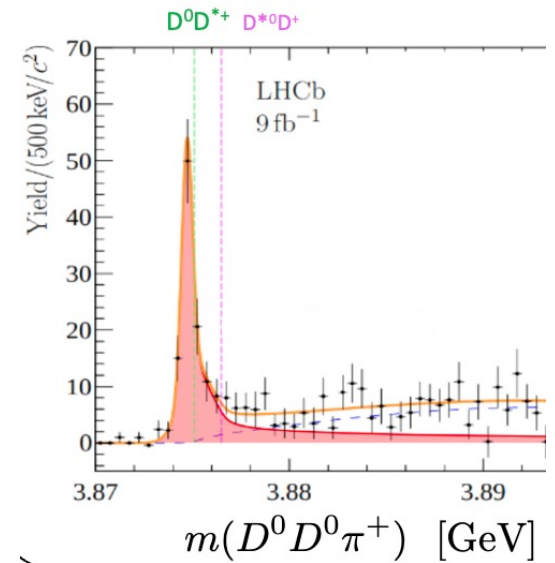
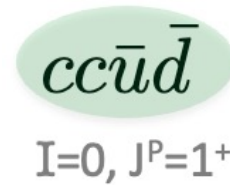


$T_{cc} (I=0, J^P=1^+)$

$cc\bar{d}\bar{u}$



Conclusions on doubly charm tetraquark



❖ The longest lived exotic hadron ever found

❖ It lies very close to DD^* threshold

❖ Lattice QCD:

to establish a state near threshold, scattering amplitude has to be extracted and pole identified

Our study [2202.101101](#) :

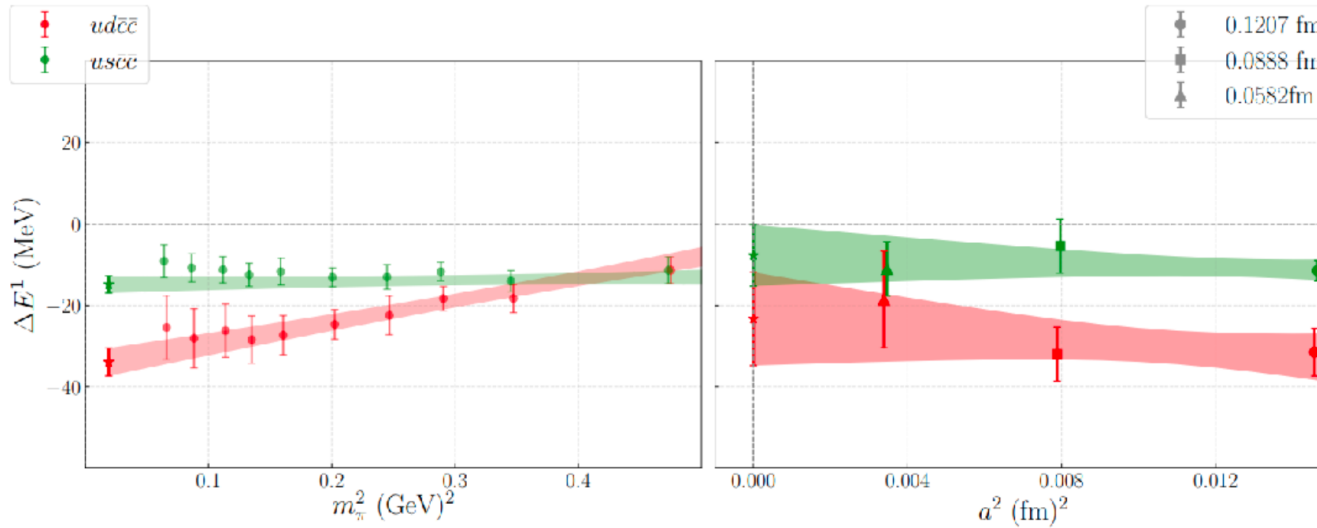
- the only extraction of DD^* scattering amplitude
- virtual bound state pole found at $m_\pi \approx 280$ MeV
- likely related to T_{cc} found by LHCb

Many interesting questions and quantities still to be explored ...

Backup

Previous lattice QCD study of T_{cc} channel

Junnarkar, Mathur, Padmanath, PRD 99, 034507 (2019), 1810.12285



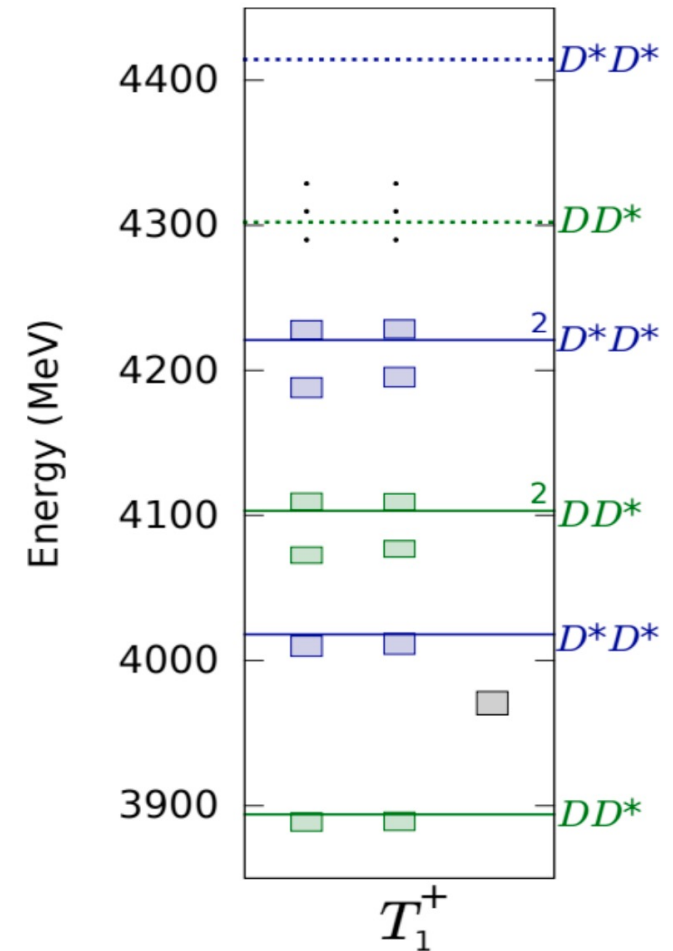
lowest finite-volume
eigen-energy for
 $P=0, J^P=1^+, I=0$

- ✿ Study performed on LQCD ensembles with different lattice spacings. Single volume and only rest frame finite-volume irreps considered.
- ✿ Including a meson-meson and diquark-antidiquark interpolator. Diquark-antidiquark interpolators do not influence the low energy spectrum.
- ✿ The ground state energy subjected to chiral and continuum extrapolations.
- ✿ A finite-volume energy level 23(11) MeV below DD^* threshold. No rigorous scattering analysis and no pole structure determined.

Previous lattice QCD study of T_{cc} channel

Hadron Spectrum, JHEP 11, 033 (2017), 1709.01417

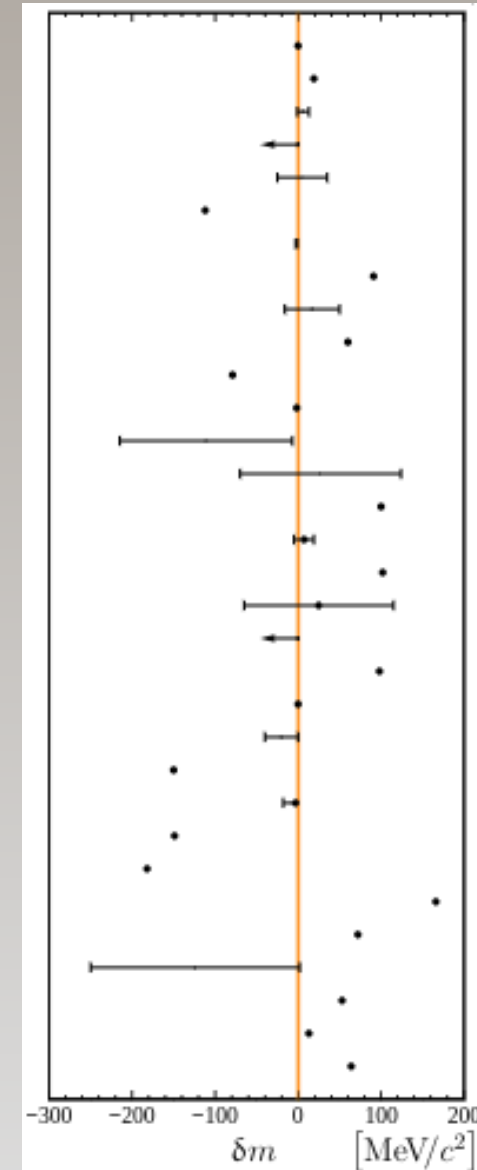
finite-volume
eigen-energies for
 $P=0, J^P=1^+, I=0$



- ❁ Single volume rest frame study on a relatively coarse lattice ($a_s \sim 0.12$ fm).
- ❁ Large basis of meson-meson and diquark-antidiquark interpolators.
- ❁ Diquark-antidiquark interpolators do not influence the low energy spectrum.
- ❁ No statistically significant energy shifts observed near DD^* threshold.
 $\Rightarrow$ No scattering amplitude extraction.

Theory predictions

Reference		Year	$\delta'm$ [MeV/c ²]
J. Carlson, L. Heller and J. A. Tjon	36	1987	~ 0
B. Silvestre-Brac and C. Semay	37	1993	+19
C. Semay and B. Silvestre-Brac	38	1994	[-1, +13]
S. Pepin, F. Stancu, M. Genovese and J. M. Richard	39	1996	< 0
B. A. Gelman and S. Nussinov	40	2002	[-25, +35]
J. Vijande, F. Fernandez, A. Valcarce, A. and B. Silvestre-Brac	41	2003	-112
D. Janc and M. Rosina	42	2004	[-3, -1]
F. Navarra, M. Nielsen and S. H. Lee	43	2007	+91
J. Vijande, E. Weissman, A. Valcarce	44	2007	[-16, +50]
D. Ebert, R. N. Faustov, V. O. Galkin and W. Lucha	45	2007	+60
S. H. Lee and S. Yasui	46	2009	-79
Y. Yang, C. Deng, J. Ping and T. Goldman	47	2009	-1.8
G.-Q. Feng, X.-H. Guo and B.-S. Zou	48	2013	-215
Y. Ikeda, B. Charron, S. Aoki, T. Doi, T. Hatsuda, T. Inoue, N. Ishii, K. Murano, H. Nemura and K. Sasaki	49	2013	[-70, +124]
S.-Q. Luo, K. Chen, X. Liu, Y.-R. Liu and S.-L. Zhu	50	2017	+100
M. Karliner and J. Rosner	51	2017	$7 \pm 12 \rightarrow 1$
E. J. Eichten and C. Quigg	52	2017	+102
Z. G. Wang	53	2017	$+25 \pm 90$
G. K. C. Cheung, C. E. Thomas, J. J. Dudek and R. G. Edwards	54	2017	$\lesssim 0$
W. Park, S. Noh and S. H. Lee	55	2018	+98
A. Francis, R. J. Hudspith, R. Lewis and K. Maltman	56	2018	~ 0
P. Jannan, N. Mathur and M. Padmanath	57	2018	[-40, 0]
C. Deng, H. Chen and J. Ping	58	2018	-150
M.-Z. Liu, T.-W. Wu, V. Pavon Valderrama, J.-J. Xie and L.-S. Geng	59	2019	-3^{+4}_{-15}
G. Yang, J. Ping and J. Segovia	60	2019	-149
Y. Tan, W. Lu and J. Ping	61	2020	-182
Q.-F. Lü, D.-Y. Chen and Y.-B. Dong	62	2020	+166
E. Braaten, L.-P. He and A. Mohapatra	63	2020	+72
D. Gao, D. Jia, Y.-J. Sun, Z. Zhang, W.-N. Liu and Q. Mei	64	2020	[-250, +2]
J.-B. Cheng, S.-Y. Li, Y.-R. Liu, Z.-G. Si, T. Yao	65	2020	+53
S. Noh, W. Park and S. H. Lee	66	2021	+13
R. N. Faustov, V. O. Galkin and E. M. Savchenko	67	2021	+64



Refs. for theory predictions

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Interpolators

Example: $P=0$

$J^P=1^+$ $\rightarrow$ cubic irrep T_1^+

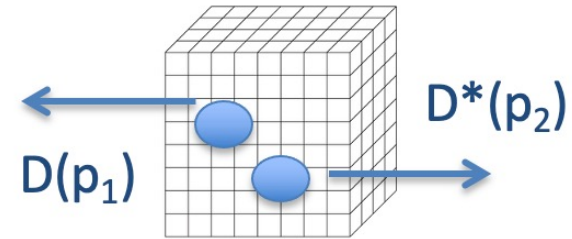
$$O^{l=0} = P(\{0, 0, 0\})V_z(\{0, 0, 0\})$$

$$O^{l=0} = P(\{1, 0, 0\})V_z(\{-1, 0, 0\}) + P(\{-1, 0, 0\})V_z(\{1, 0, 0\}) \\ + P(\{0, 1, 0\})V_z(\{0, -1, 0\}) + P(\{0, -1, 0\})V_z(\{0, 1, 0\}) \\ + P(\{0, 0, 1\})V_z(\{0, 0, -1\}) + P(\{0, 0, -1\})V_z(\{0, 0, 1\})]$$

$$O^{l=2} = P(\{1, 0, 0\})V_z(\{-1, 0, 0\}) + P(\{-1, 0, 0\})V_z(\{1, 0, 0\}) \\ + P(\{0, 1, 0\})V_z(\{0, -1, 0\}) + P(\{0, -1, 0\})V_z(\{0, 1, 0\}) \\ - 2[P(\{0, 0, 1\})V_z(\{0, 0, -1\}) + P(\{0, 0, -1\})V_z(\{0, 0, 1\})]$$

$$O^{l=0} = V_{1x}[0, 0, 0]V_{2y}[0, 0, 0] - V_{1y}[0, 0, 0]V_{2x}[0, 0, 0]$$

$P=D, V=D^*$



Relation between E and $\delta(E), T(E)$

$$S = 1 + i \frac{4p}{E} T = e^{2i\delta}$$

$$T = \frac{E}{2} \frac{1}{p \cot \delta - ip}$$

E = eigen-energy lattice from lattice in cmf

$$E = \sqrt{m_1^2 + p^2} + \sqrt{m_2^2 + p^2}$$

even and odd l contribute to
given irrep for nonzero mom.

Luscher's relation ($l=0,1$):

$$\det \left[\begin{pmatrix} p \cot \delta_0 & 0 \\ 0 & p^3 \cot \delta_1 \end{pmatrix} - \underset{\substack{\downarrow \\ \text{known 2x2 matrix} \\ \text{of kinematical} \\ \text{functions} \\ \text{(non-diagonal)}}}{B(E, L)} \right] = 0$$

Luscher's relation (only $l=0$):

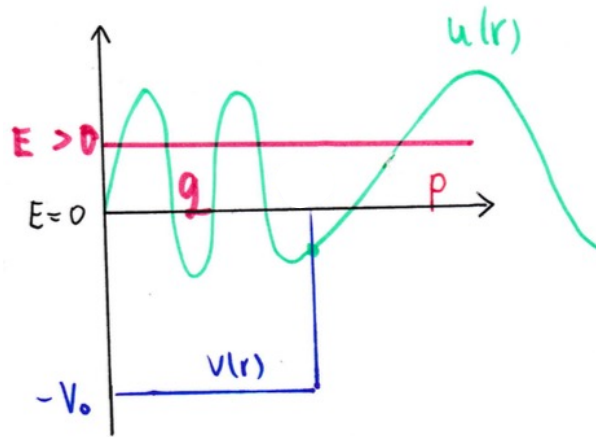
$$p \cot \delta_0 = B(E, L)$$

$$p \cot \delta_0 - B(E, L) = 0$$

$\downarrow$
known
kinematical
function

$\downarrow$
lattice
eigen-energy

s-wave scattering on spherical potential well



$$A \sin qr \quad B \sin(pr + \delta_0)$$

$$\left. \begin{aligned} u(R) &= A \sin qR = B \sin(pR + \delta) \\ u'(R) &= q A \cos qR = p B \cos(pR + \delta) \end{aligned} \right\}$$

$$q = \sqrt{2\mu(V_0 + E)} = \sqrt{2\mu V_0 + p^2}$$

dividing both eqs

$$\frac{1}{q} \tan qR = \frac{1}{p} \tan(pR + \delta)$$

$$\delta_0(p) = \arctan\left(\frac{p}{q} \tan(qR)\right) - pR + n\pi$$