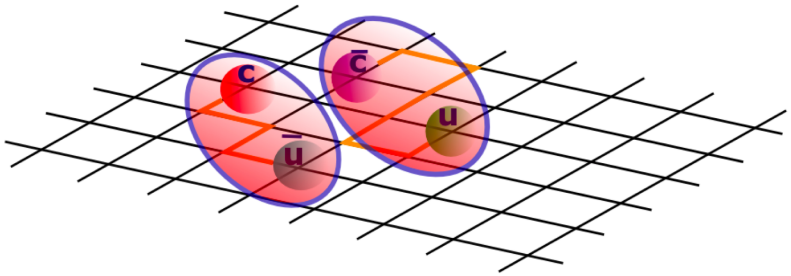
Charmonium(like) and other near-threshold mesons on the lattice



Sasa Prelovsek

University of Ljubljana & Jozef Stefan Institute, Slovenia

Jefferson Lab, Virginia, USA

Fermilab, 31st July 2015

Outline

- Experimental appetizer concerning exotic hadrons:
  - very recent discovery of pentaquarks
  - recent discoveries of tetraquarks (just one example)

- Theoretical approach to study
  conventional and exotic hadrons
  based directly on QCD – lattice QCD

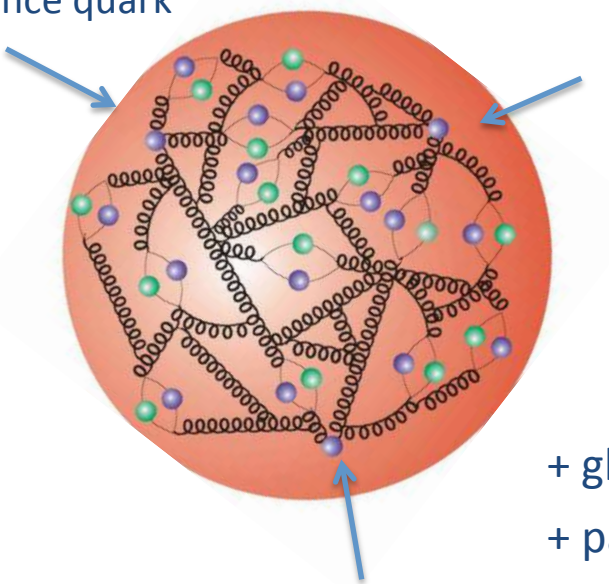
Hadrons

Conventional

Normal baryon

Normal meson

valence quark



+ gluons

+ pairs of sea $\underline{q}q$

Exotic

minimal valence content

$\bar{q}\ \bar{q}\ q\ q$ tetraquarks

$\bar{q}\ q\ q\ q\ q$ pentaquarks

.... hybrids

glueballs

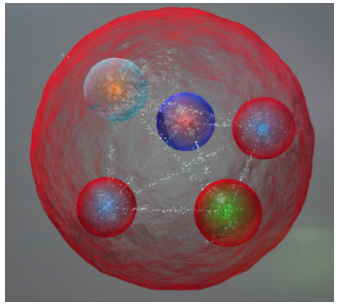
....

Terminology in this talk: tetra(penta) quarks indicate just the number of valence quarks in the state; it is not meant to say anything on how quarks are clustered in them

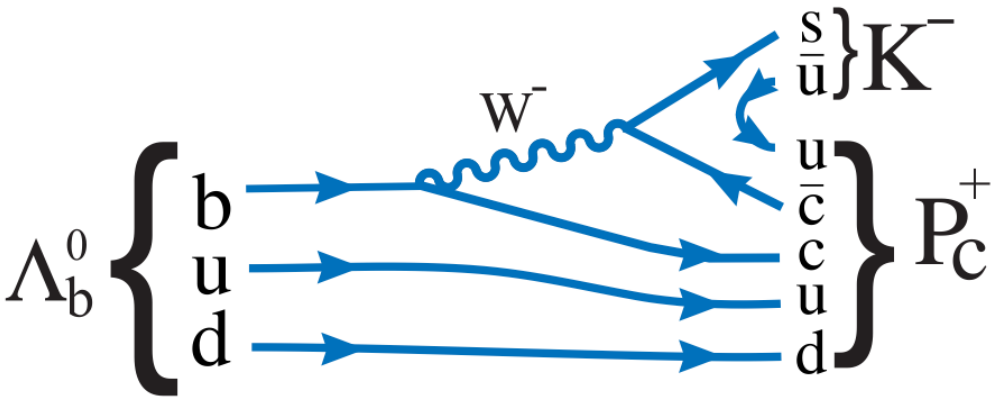
Pentaquarks

LHCb: 1507.03414

14th July 2015

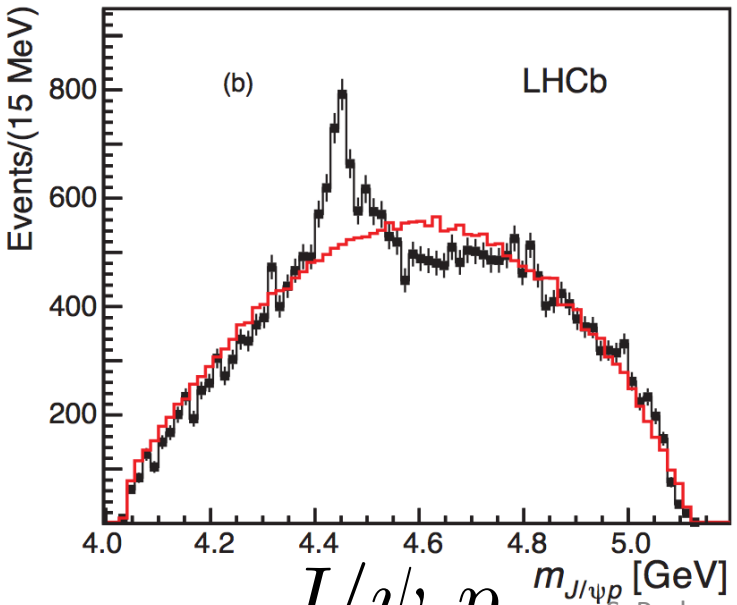


$$\Lambda_b^- \to K^- J/\psi \; p$$

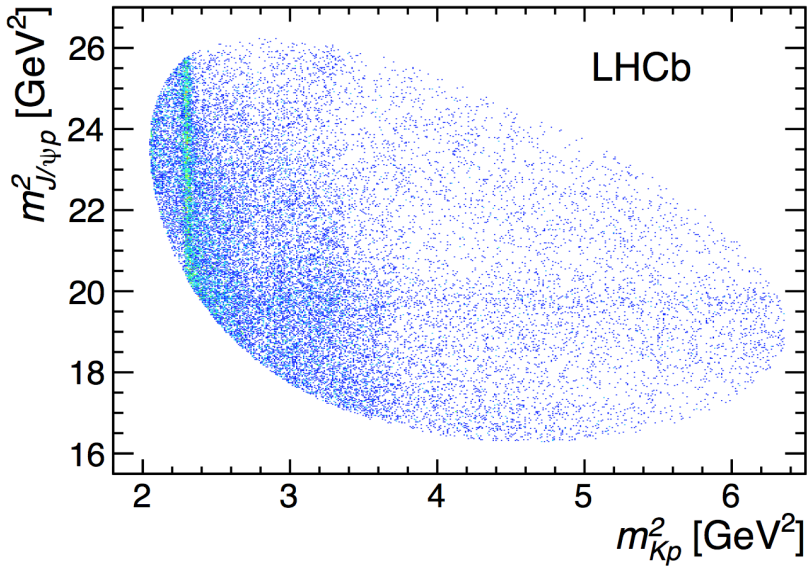


$$P_c^+ \to J/\psi \;\; p$$

$$\bar{c}c \;\; uud$$



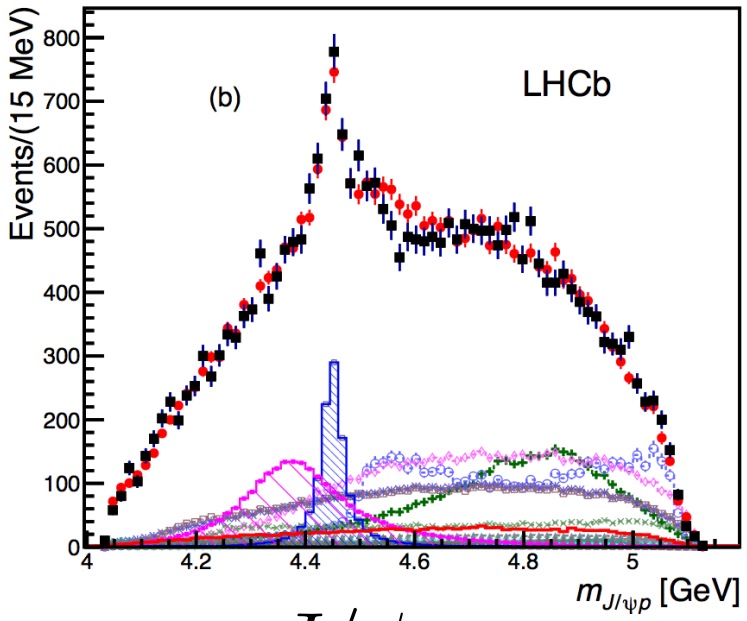
$J/\psi \;\; p$



$J/\psi \;\; p$

Two pentaquark candidates from LHCb

LHCb: 1507.03414

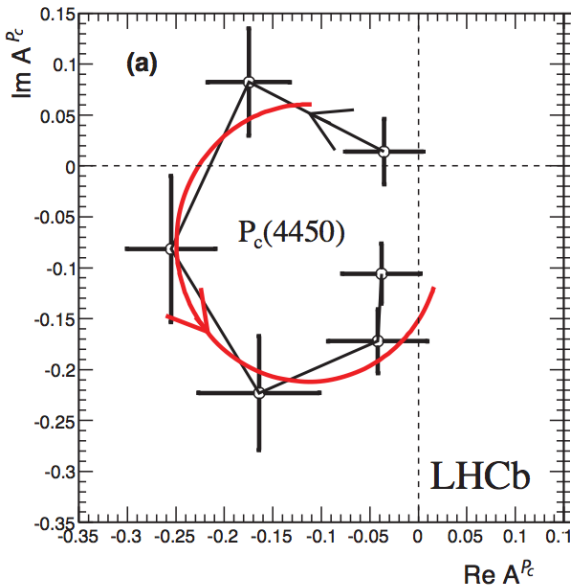


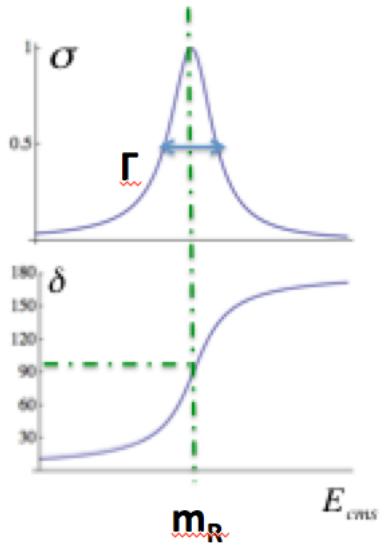
$J/\psi\ p$

scattering matrix T for elastic scattering

$$A \simeq T = \frac{1}{2i}(e^{2i\delta(E)} - 1)$$

$$|T + \frac{1}{2i}| = |\frac{1}{2i}e^{2i\delta(E)}| = \frac{1}{2}$$





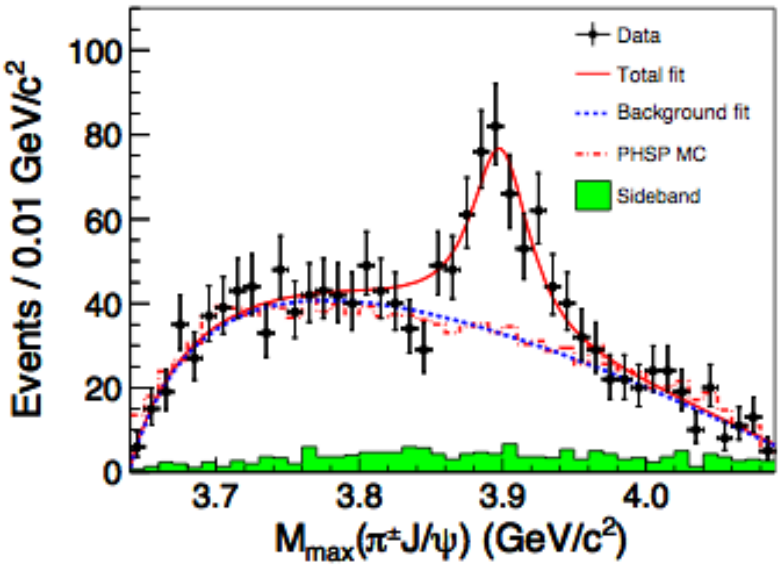
$P_c(4380)$: $M_1 = 4380 \pm 8 \pm 29$ MeV $\Gamma_1 = 205 \pm 18 \pm 86$ MeV $J^P = 3/2^+$ or $5/2^+$

$P_c(4450)$: $M_2 = 4449.8 \pm 1.7 \pm 2.5$ MeV $\Gamma_2 = 39 \pm 5 \pm 19$ MeV $J^P = 5/2^-$ or $3/2^-$

Tetraquarks Z_c

BESIII, Belle, Cleo-c 2013

Example: $Z_c^+(3900)$



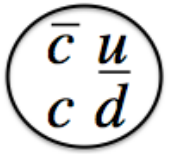
DD* thr.

[BESIII, 2013, 1303.5949, PRL]

state confirmed by Belle, Cleo-c

$$Z_c^+(3900) \to J/\psi\ \pi^+$$

flavor content of the hadron: $\bar{c}c\ \ \bar{d}u$



M = 3888 ± 3 MeV , Γ = 35 ± 7 MeV

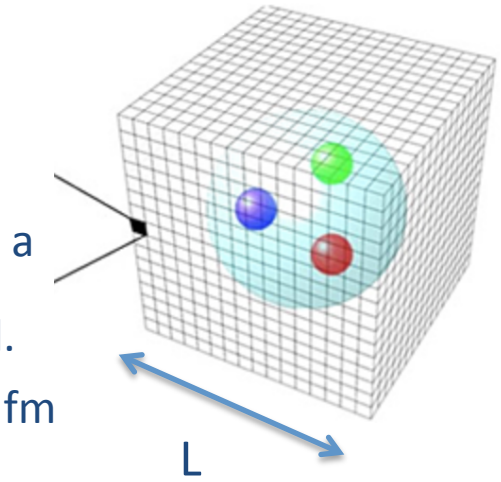
Non-perturbative method: QCD on lattice

$$L_{QCD} = -\tfrac{1}{4} G^a_{\mu\nu} G^{\mu\nu}_a + \sum_{q=u,d,s,c,b,t} \bar{q} i \gamma_\mu (\partial^\mu + i g_s G^\mu_a T^a) q - m_q \bar{q} q$$

input: g_s , m_q

output: hadron properties

hadron interactions



precision cal.

a~0.05-0.15 fm

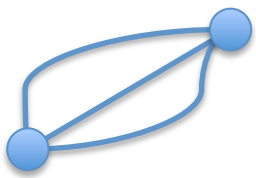
L~ 2-8 fm

Evaluation of Feynman path integrals in discretized space-time

quantum mechanics

$$\int Dx \, e^{i S/\hbar}$$

$$S = \int dt \, L[x(t)]$$



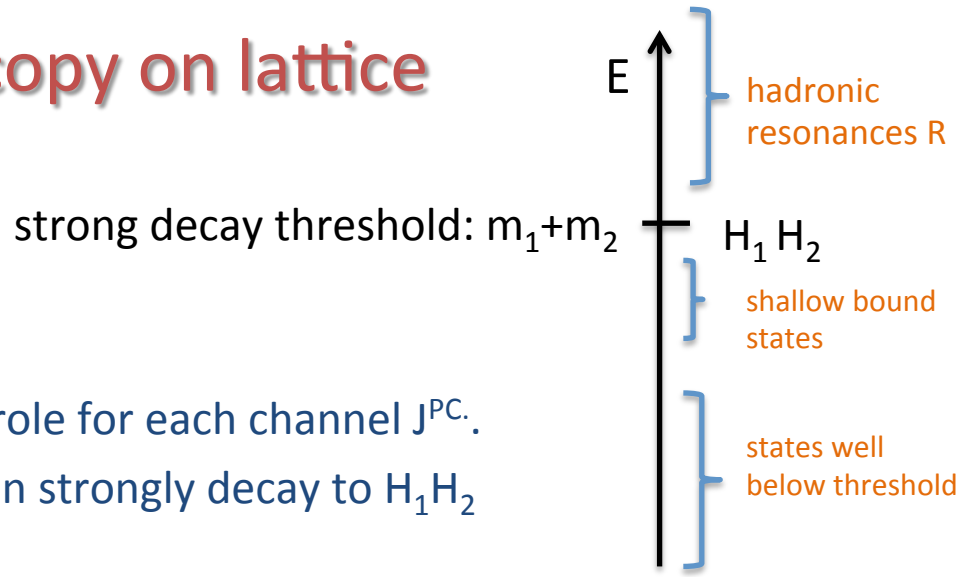
quantum field theory in Euclidian space-time

$$\int DG \, Dq \, D\bar{q} \, e^{-S_{QCD}/\hbar}$$

$$S_{QCD} = \int d^4x \, L_{QCD}[G(x), q(x), \bar{q}(x)]$$

x,t (Minkovsky) → x, i t (Euclidean)

Status of hadron spectroscopy on lattice



The location of the threshold pays a major role for each channel J^{PC}.
If the hadron mass is above threshold, it can strongly decay to H_1H_2 as long as quantum numbers allow decay.

- Hadrons well bellow threshold: *"easy"*
  Unfortunately none of exotic experimental candidates is found well below threshold.

- Hadron resonances above threshold and states slightly below threshold : *challenging*
  - until recently: so-called single-hadron approximation
  - now: rigorous treatment by determining scattering matrix for two hadrons H_1 H_2

- lattice searches for manifestly exotic hadrons (usually above several thresholds): *very challenging*

Lattice setup

PACS-CS

	Ensemble (1)	Ensemble (2)
$N_L^3 \times N_T$	$16^3 \times 32$	$32^3 \times 64$
N_f	2	2+1
a [fm]	0.1239(13)	0.0907(13)
L [fm]	1.98(2)	2.90(4)
m_π [MeV]	266(3)(3)	156(7)(2)

- Wilson-clover quarks
- Fermilab method for c and b : $E_M(p) = M_1 + \frac{\mathbf{p}^2}{2M_2} - \frac{a^3 W_4}{6} \sum_i p_i^4 - \frac{(\mathbf{p}^2)^2}{8M_4^3} + \dots$

  [El Khadra, Kronfeld, Mackenzie, 1997]

  Only splittings with respect to a chosen reference mass are compared to experiment.

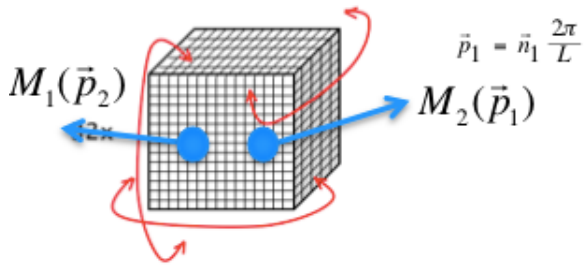
- evaluating Wick contr.: distillation (Ensemble 1) [Peardon et. al., HSC, 2009]

  stochastic distillation (Ensemble 2) [Morningstar et al., 2011]

Discrete energy spectrum from correlators

Example: meson channel with given J^{PC}

$$\mathcal{O} = \bar{q}\Gamma q, \quad (\bar{q}\Gamma_1 q)(\bar{q}\Gamma_2 q), \quad [\overline{qq}][qq]$$



$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t)\, \mathcal{O}_j^+(0) | 0 \rangle = \sum_n Z_i^n\, Z_j^{n*}\, e^{-E_n t} \qquad Z_i^n = \langle 0 | \mathcal{O}_i | n \rangle$$

Energies and overlaps extracted using GEVP $\quad C(t)u^{(n)}(t) = \lambda^{(n)}(t)C(t_0)u^{(n)}(t)$

All physical states with given J^{PC} appear as energy levels E_n in principle : single particle, two-particle,...

channel	:	"eigenstates"
$J^{PC} = 1^{--}$, $\bar{u}d$	:	$\rho,\ \pi\pi$
$J^{PC} = 1^{++}$, $\bar{c}c$	:	$\chi_{c1}, X(3872),\ DD^*$
$J^{PC} = 1^{+-}$, $\bar{c}c\bar{u}d$	:	$J/\psi,\ DD^*$

Two-meson states:

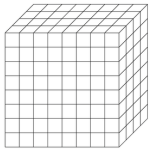
- In experiment: two-meson decay products with continuous E.
- On lattice: discrete E due to finite L and periodic BC: p=n 2π/L

Bound state and narrow resonance:

typically lead to extra energy level (in addition to two-meson levels)

States well below threshold: "easy" precision spectrum

- $m=E_n$ for P=0 $a\to 0$, $L\to\infty$, $m_q\to m_q^{phy}$



- Available from a number of lattice QCD collaborations for a number of years
- Only few examples shown

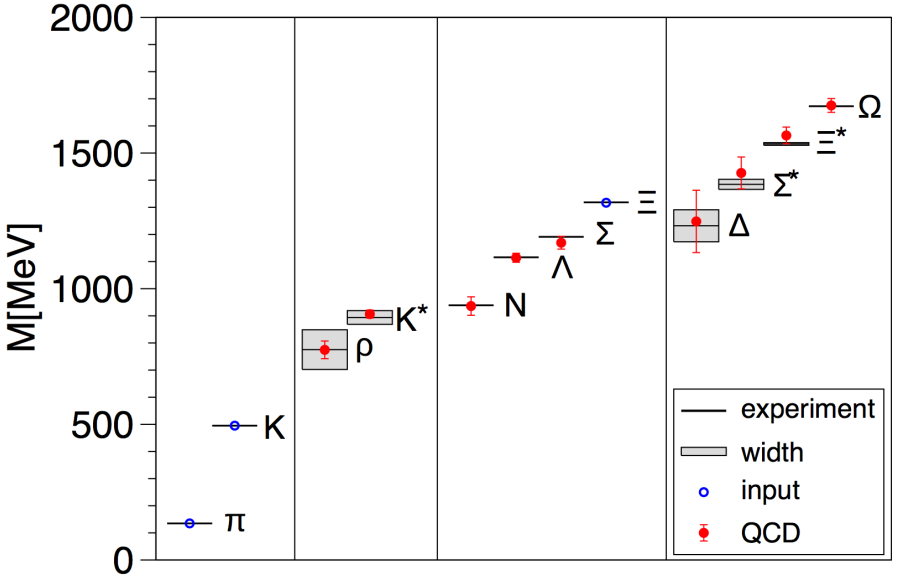
Proton and neutron

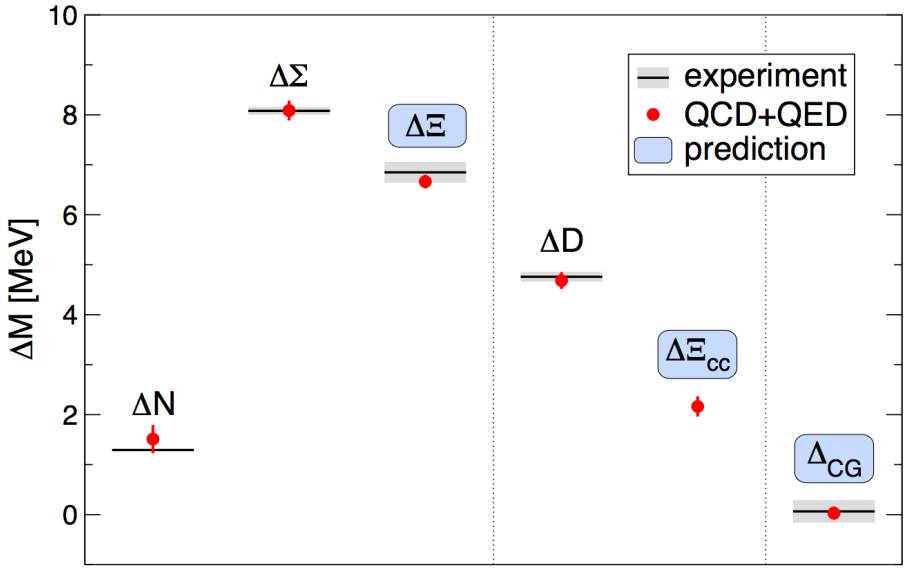
constitute more than 99%
of the bright side of universe

$m_u c^2 \approx 2\ \text{MeV}$ $\quad m_d c^2 \approx 5\ \text{MeV}$

$m_p c^2 \approx 938.3\ \text{MeV}$ $\quad m_n c^2 \approx 939.5\ \text{MeV}$

Higgs mechanism provides tiny contribution, the rest of visible mass in universe is due to the strong interaction in hadrons





BMW collaboration
Science 322,2008

QCD

$m_u = m_d$

m_n , m_p

m_n - m_p

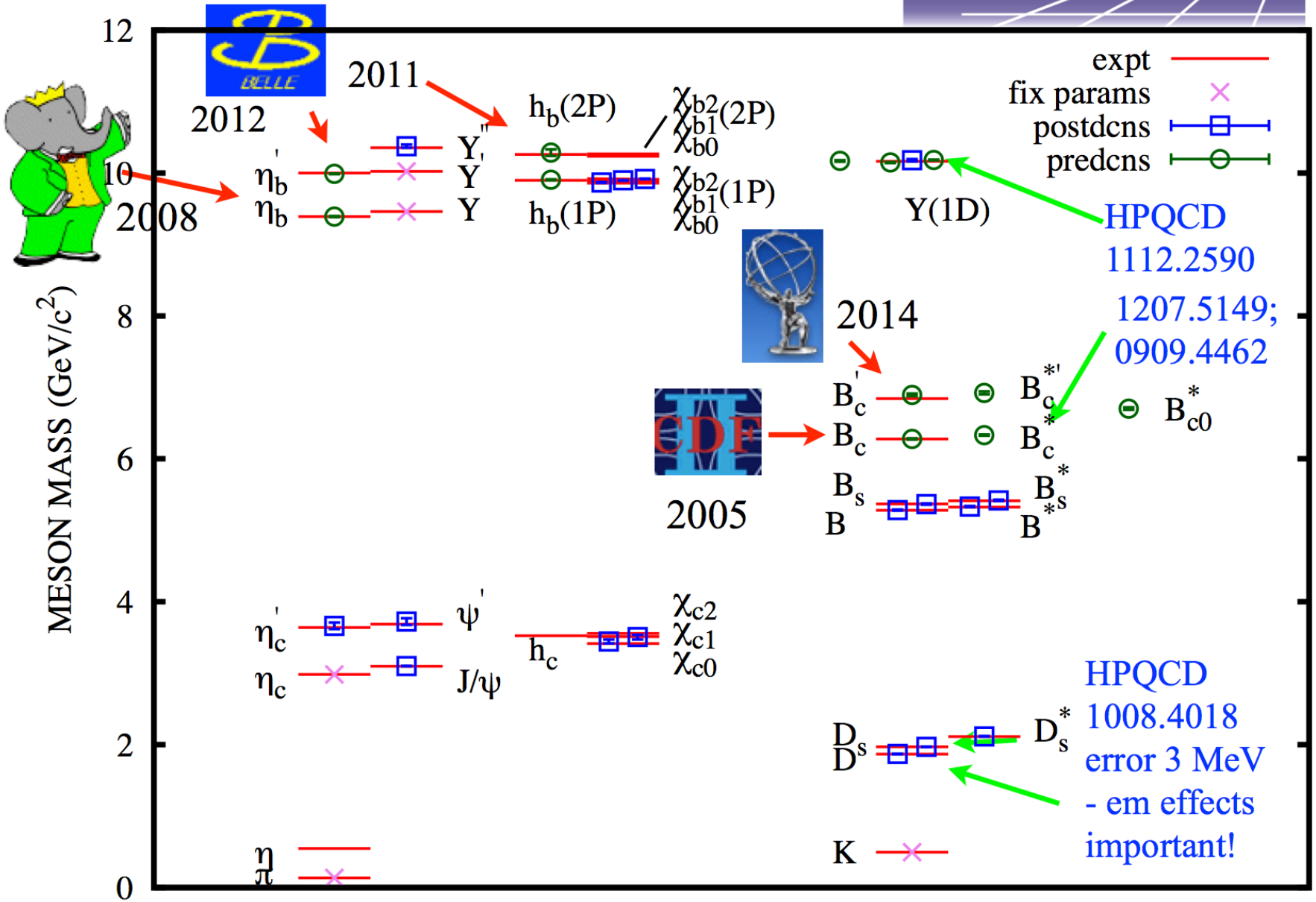
BMW collaboration
Science 347,2015

QCD + QED

$m_u \neq m_d$

High Precision QCD

$m=E_n$ The gold-plated meson spectrum



Hadrons near or above strong decay threshold challenging

Single–hadron approximation used until few years ago:

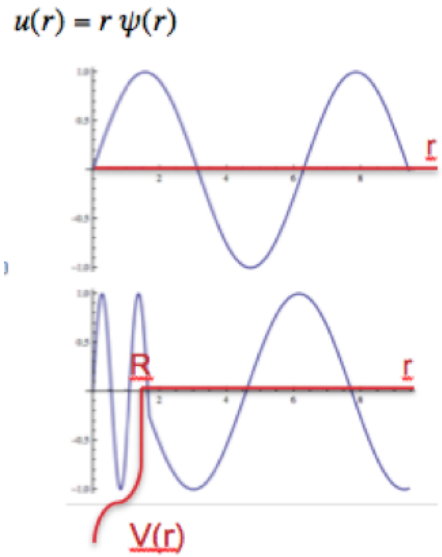
ignores strong decays of resonances and effects of thresholds

I will concentrate on "Rigorous approach" which addresses

scattering of two hadrons ..... in experiment or theory

Two-hadron scattering

Scattering phase shift δ

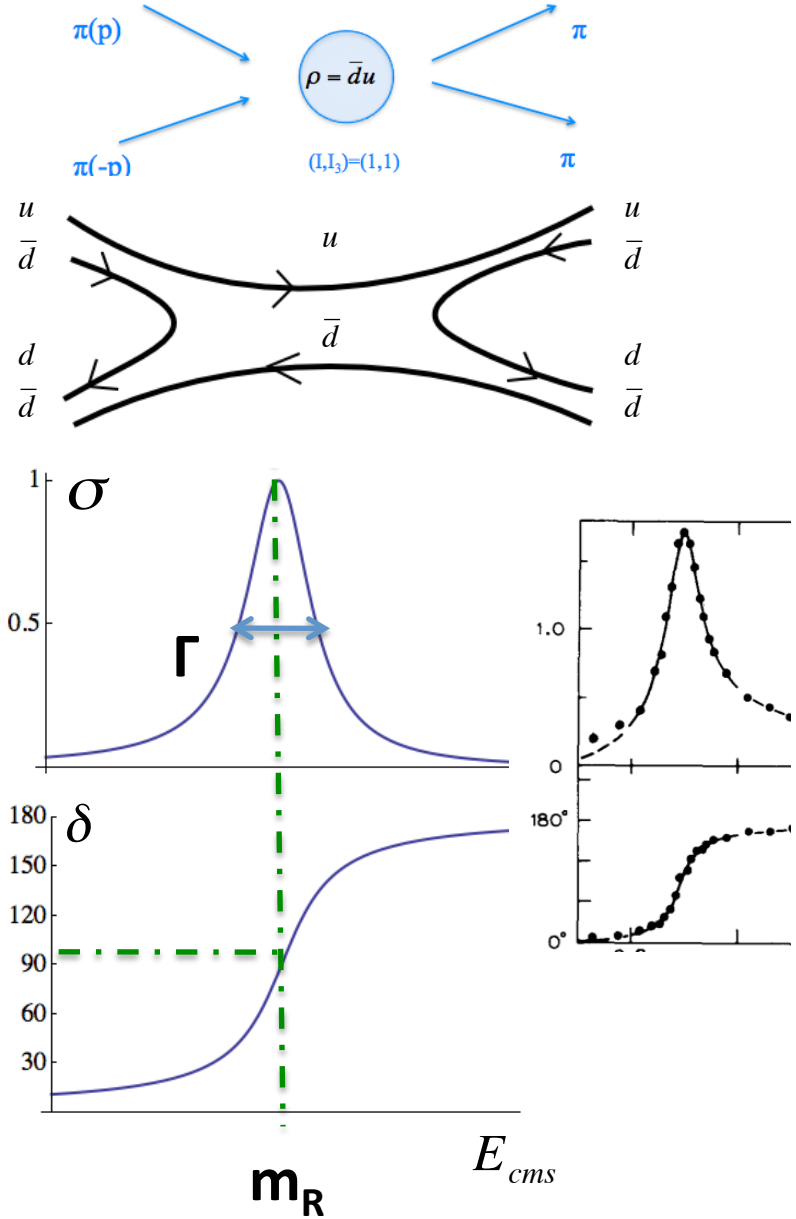


QM interpretation:

$$\psi(r) \propto \frac{\sin(pr)}{r}$$

$$\psi(r) \propto \frac{\sin(pr+\delta)}{r} \qquad r > R$$

$$\text{unknown} \qquad r < R$$



scattering matrix T(E) from δ(E) for elastic scat:

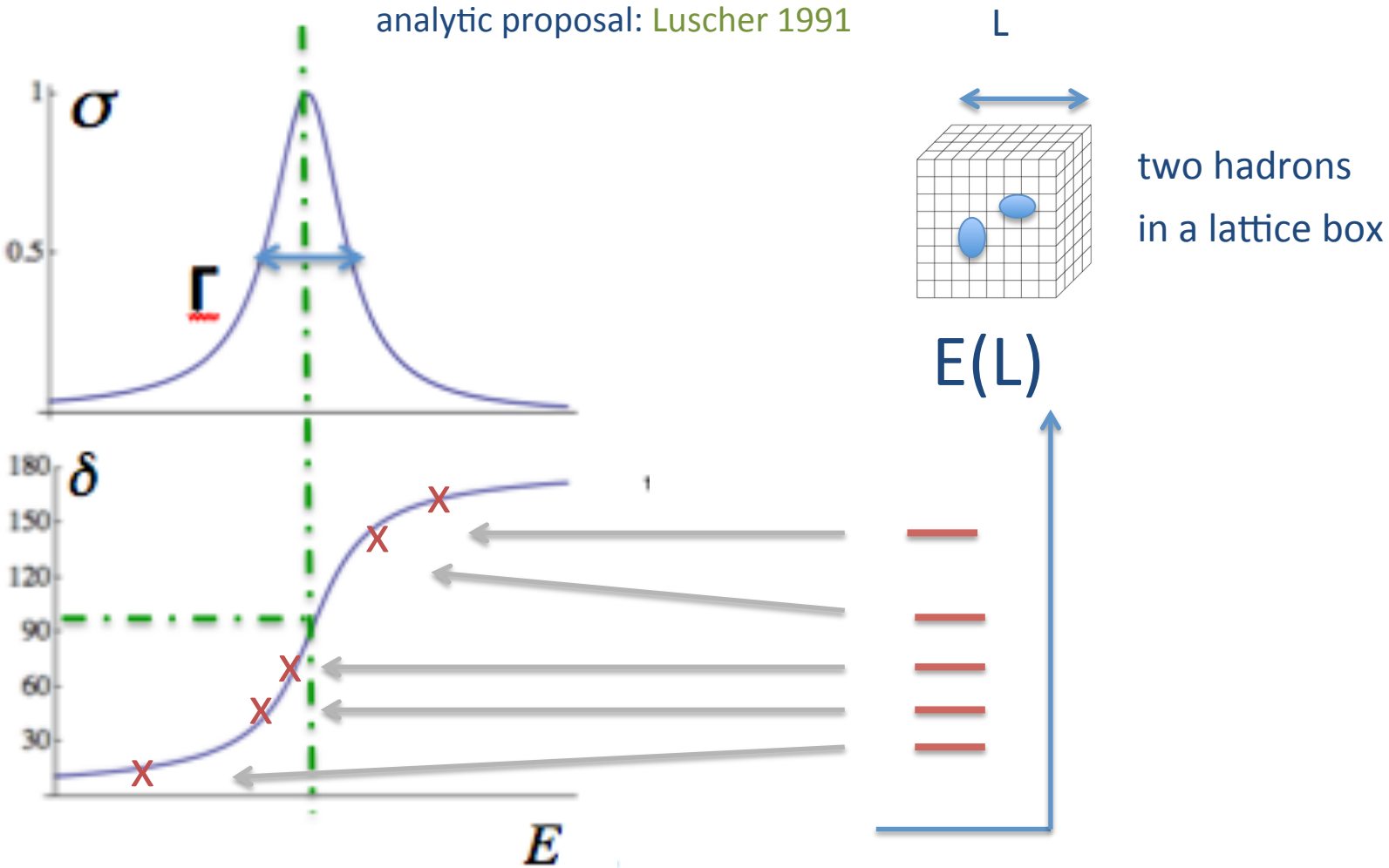
$$S(E) = 1 + 2\,i\,T(E) = e^{2\,i\,\delta(E)}$$

$$T = \frac{e^{2i\delta} - 1}{2i} = \sin\delta\; e^{i\delta} = \frac{1}{\cot\delta - i}$$

$$\sigma \propto |T|^2 = \sin^2\delta$$

Rigorous treatment: scattering info from spectrum

analytic proposal: Luscher 1991



scattering phase shifts at infinite volume $\delta(E)$ ← $E(L)$ energies from lattice with spatial extent L

Relation between E and δ(E)

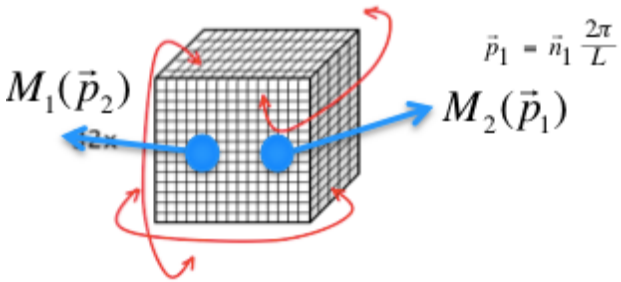
analytic proposal: Luscher 1991

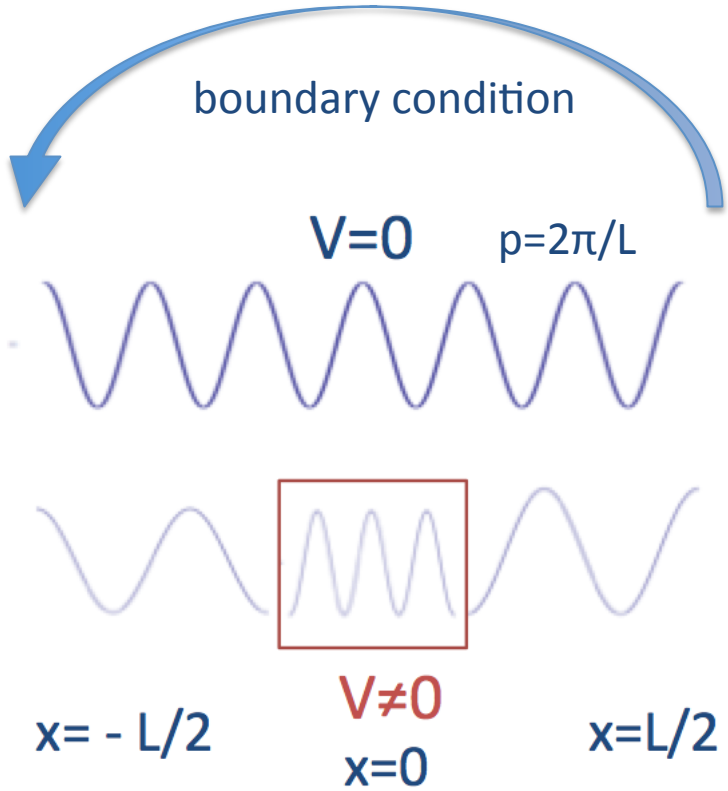
Energies of two hadrons in a box:

$$E(L) = \sqrt{m_1^2 + \vec{p}_1^{\,2}} + \sqrt{m_2^2 + \vec{p}_2^{\,2}} + \Delta E$$

- due to strong interaction
- gives rigorous info on δ

$$\vec{p}_{1,2} = \frac{2\pi}{L}\vec{n}$$

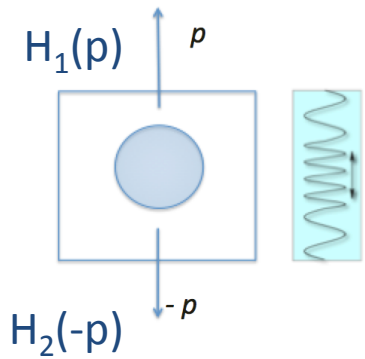




Scattering of two hadrons

elastic scattering with total momentum P=0: $E=E_{cm}$

$E_n(L) \xrightarrow{\text{Luscher's eq.}} \delta(E)$



Scattering matrix for partial wave l:

$$S(E) = e^{2i\delta(E)}, \quad S(E) = 1 + 2iT(E), \quad T(E) = \frac{1}{\cot\delta(E) - i}$$

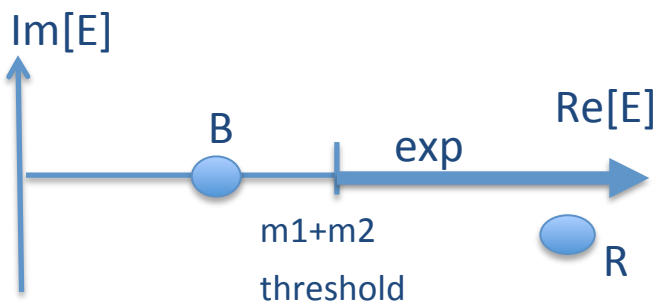
Bound state (B):

$\cot[\delta(E_B)] = i$, $E_B < m_1+m_2$

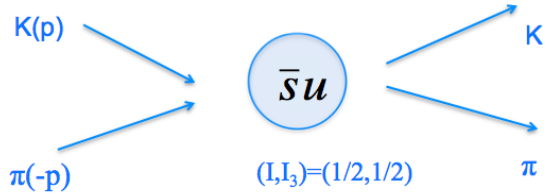
Resonance (R) (of Breit-Wigner type):

$$T(E) = \frac{-E\,\Gamma}{E^2 - m_R^2 + i\,E\,\Gamma}, \quad \Gamma(E) = g^2\,\frac{p^{2l+1}}{E^2}$$

Locations of poles of T(E)
for res. and bound st.



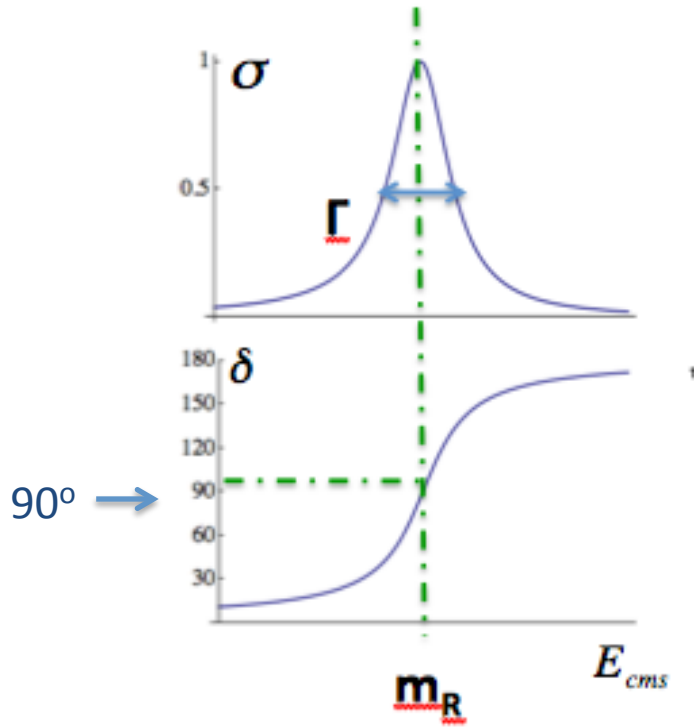
Resonances above threshold



Decay quickly through strong interaction

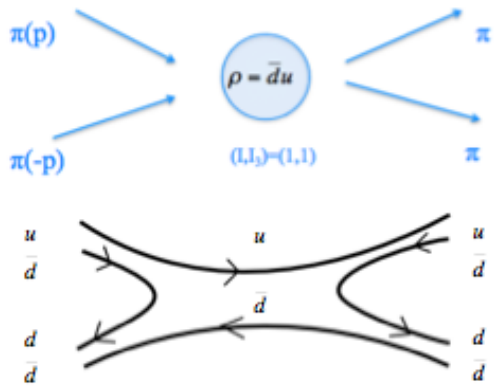
$$\tau = \hbar / \Gamma$$

Γ=10–300 MeV



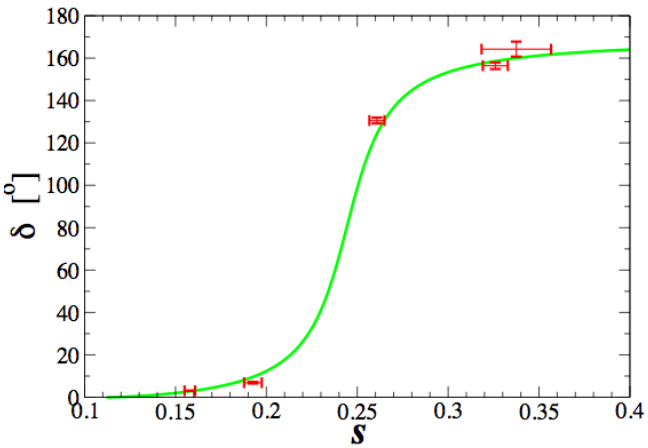
ρ resonance

(the only "mature" hadronic resonance from lattice)



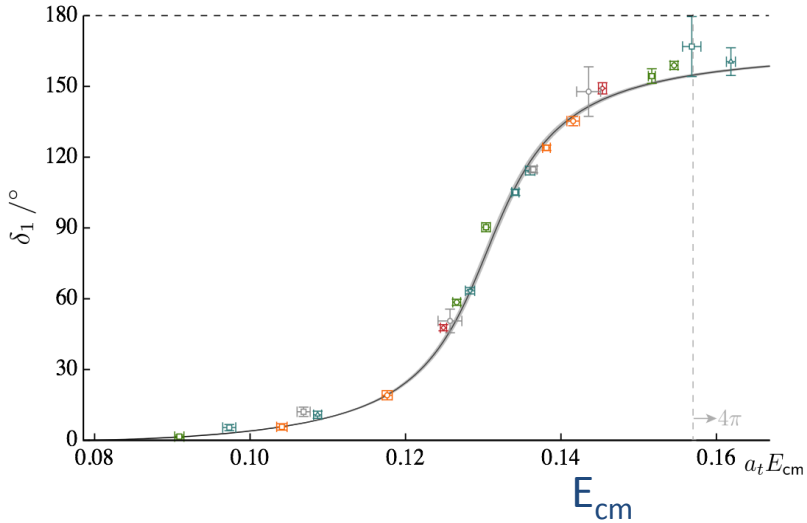
$$\mathcal{O}: \ \bar{u}d, \ \pi\pi$$

$E_{cm} \rightarrow \delta(E_{cm})$

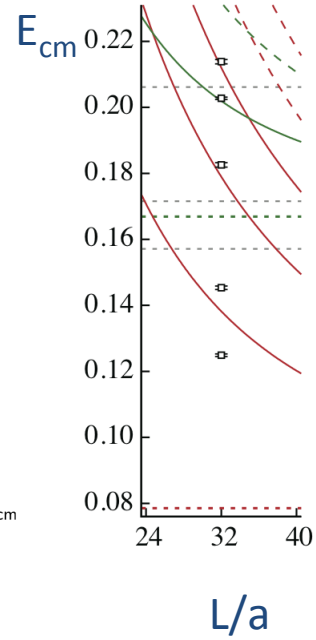


Lang, Mohler, S.P. , Vidmar, PRD 2011,

$m_\pi \approx 266$ MeV



[HSC, PRD 1507.02599] $m_\pi \approx 236$ MeV



$$\text{BW}: \quad \delta = \text{acot}\frac{m_R^2 - E_{cms}^2}{m_R\,\Gamma} \quad \rightarrow \quad m_R, \quad \Gamma \text{ or } g_{\rho\pi\pi}$$

Other resonances simulated on the lattice

$\bar{c}c$	$\bar{u}u$	$\bar{s}u$	$\bar{c}u$
$\eta_c(1S)$	$\pi^\pm$	$K^\pm$	$D^\pm$
$J/\psi(1S)$	π^0	K^0	D^0
$\chi_{c0}(1P)$	η	$K_S{}^0$	$D^*(2007)^0$
$\chi_{c1}(1P)$	$f_0(500)$ or σ was $f_0(60$	$K_L{}^0$	$D^*(2010)^\pm$
$h_c(1P)$	$\rho(770)$	$K_0^*(800)$ or κ	$D_0^*(2400)^0$
$\chi_{c2}(1P)$	$\omega(782)$	$K^*(892)$	$D_0^*(2400)^\pm$
$\eta_c(2S)$	$\eta'(958)$	$K_1(1270)$	$D_1(2420)^0$
$\psi(2S)$	$f_0(980)$	$K_1(1400)$	$D_1(2420)^\pm$
$\psi(3770)$	$a_0(980)$	$K^*(1410)$	$D_1(2430)^0$
$X(3872)$	$\phi(1020)$	$K_0^*(1430)$	$D_2^*(2460)^0$
$\chi_{c0}(2P)$wa	$h_1(1170)$	$K_2^*(1430)$	$D_2^*(2460)^\pm$
$\chi_{c2}(2P)$	$b_1(1235)$	$K(1460)$	$D(2550)^0$
$X(3940)$	$a_1(1260)$	$K_2(1580)$	$D(2600)$
$\psi(4040)$	$f_2(1270)$	$K(1630)$	$D^*(2640)^\pm$
$X(4050)^\pm$	$f_1(1285)$	$K_1(1650)$	$D(2750)$
$X(4140)$	$\eta(1295)$	$K^*(1680)$	
$\psi(4160)$	$\pi(1300)$	$K_2(1770)$	
$X(4160)$	$a_2(1320)$	$K_3^*(1780)$	
$X(4250)^\pm$	$f_0(1370)$	$K_2(1820)$	
	$h_1(1380)$	$K(1830)$	
	$\pi_1(1400)$		
	$\eta(1405)$		
	$f_1(1420)$		
	$\omega(1420)$		
	$f_2(1430)$		
	$a_0(1450)$		
	$\rho(1450)$		

well below open charm decay th.

well below strong decay th.

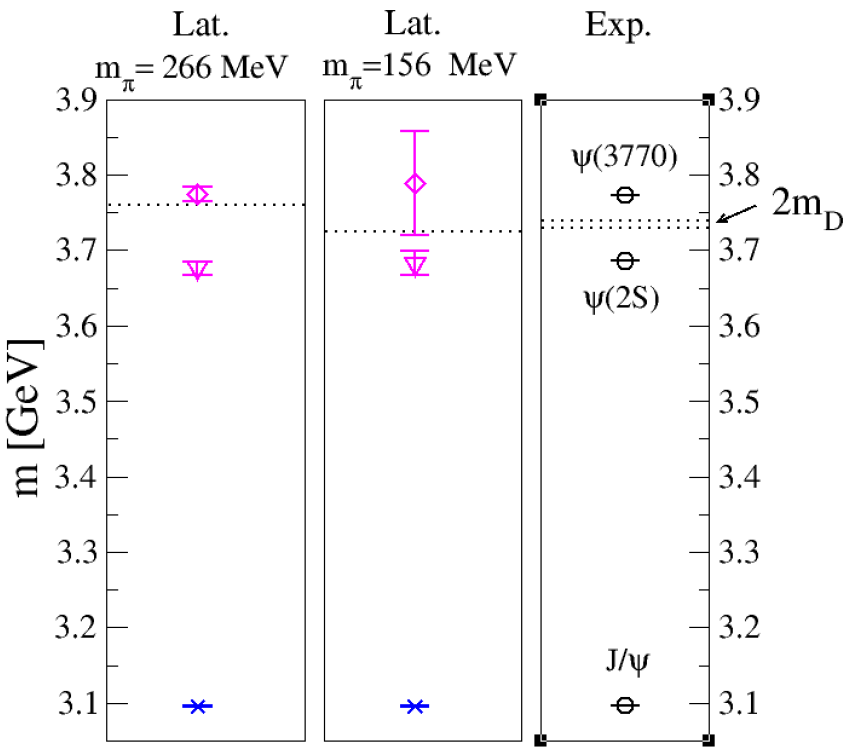
pioneering simulations of resonances with rigorous treatment 2012-2015

cu [Mohler, S.P. , Woloshyn: PRD 2013] su [S.P., Lang, Leskovec, Mohler, PRD 2013; Wilson, Dudek, Edwards, Thomas, PRL 20114, PRD 2015] cc [Lang, Leskovec, Mohler, S.P., 1503.05363] a1,b1 [Lang, Leskovec, Mohler, S.P., JHEP 2014]

General result: m_R and Γ (or coupling) close to experimental values.

Exception: $\chi_{c0}(2P)$ - puzzling state

Resonance ψ(3770) and bound st. ψ(2S) from DD̄ scattering in p-wave



$\mathcal{O}$: $\bar{c}\,c$, $D\bar{D}$, $J^{PC} = 1^{--}$

DD̄ scat. in p-wave is simulated

T-matrix is determined from E_n

Fit of T-matrix gives:

BW resonance ψ(3770):

m_R (magenta diamonds)

Γ (given below)

Bound state ψ(2S) from pole in T:

m_B (magetna triangles)

$\eta_c(1S)$
$J/\psi(1S)$
$\chi_{c0}(1P)$
$\chi_{c1}(1P)$
$h_c(1P)$
$\chi_{c2}(1P)$
$\eta_c(2S)$
$\psi(2S)$
$2m_D$
$\psi(3770)$
$X(3872)$
$\chi_{c0}(2P)$wa
$\chi_{c2}(2P)$
$X(3940)$
$\psi(4040)$
$X(4050)^{\pm}$
$X(4140)$
$\psi(4160)$
$X(4160)$
$X(4250)^{\pm}$

ψ(3770)	Mass [MeV]	g (no unit)
Lat (m_π=266 MeV)	3774 ±6±10	19.7 ±1.4
Lat (m_π=156 MeV)	3789 ±68±10	28 ± 21
Exp.	3773.15± 0.33	18.7 ± 1.4

$$\Gamma = \frac{g^2}{6\pi}\frac{p^3}{s}$$

Lang, Leskovec, Mohler, S.P., 1503.05363

States slightly below threshold

$p\eta_c$ pentaquark

D_{s0}^*, D_{s1}, B_{s0}^*, B_{s1}, X(3872)

deuterium-like systems

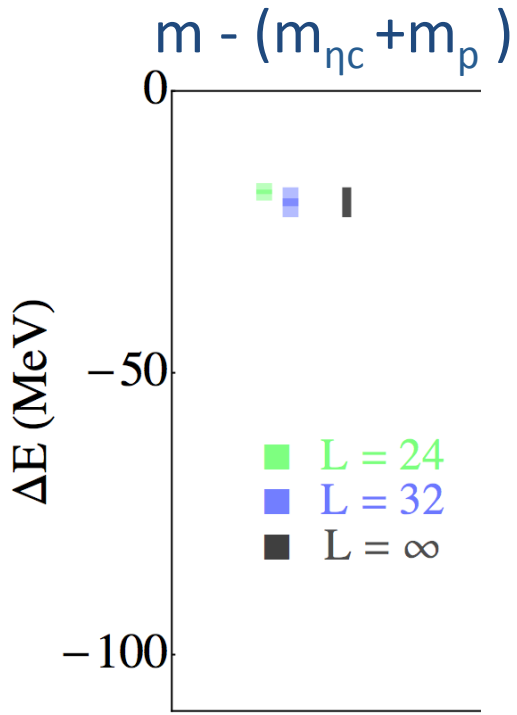
Bound state of a η_c and p from lattice

[NPLQCD, 1410.7069, PRD, m_π~800 MeV]

$\eta_c \; p$

$\bar{c}c \; uud$

~ 20 MeV below th. $m_{\eta c}+m_p$



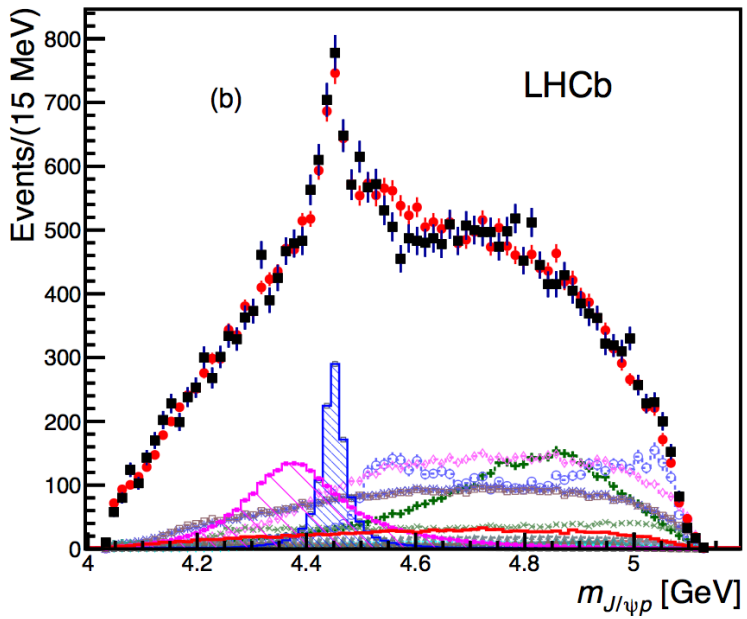
Two pentaquark resonances in J/ψ and p from exp

LHCb: 1507.03414

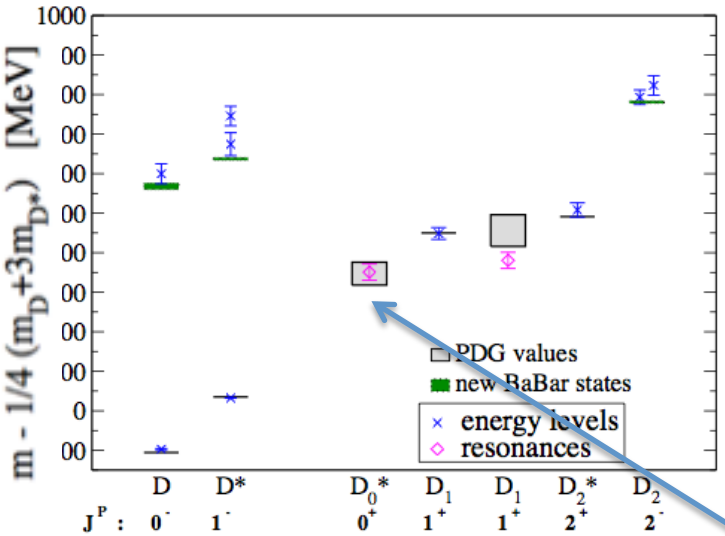
$J/\psi \; p$

$\bar{c}c \; uud$

~ 400 MeV above th. $m_{J/\psi}+m_p$



D_s and D scalar meson puzzle



D mesons (resonances)

scalar D resonance in Dπ

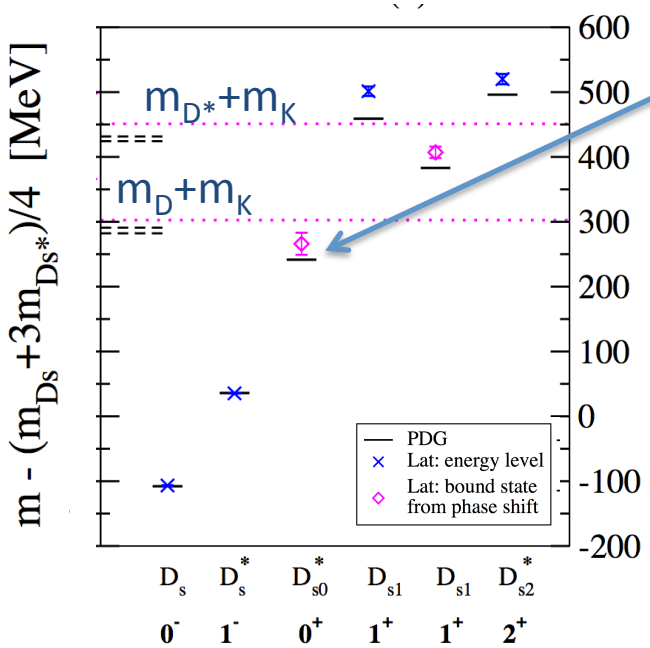
[D. Mohler, S.P. , R. Woloshyn: PRD 2013]

- why do these scalar partners have mass so close ?

$$D_0^*(2400): \quad M \approx 2318 \text{ MeV} \quad \Gamma \approx 267 \text{ MeV} \quad \bar{c}u \text{ or } \bar{c}u\bar{s}s \,?$$

$$D_{s0}(2317): \quad M \approx 2318 \text{ MeV} \quad \Gamma \approx 0 \text{ MeV} \quad \bar{c}s \text{ or } \bar{c}s[\bar{u}u + \bar{d}d] \,?$$

1) is D_0^* mass pushed up : valence ss pair ?? ✖
2) is D_{s0}^* mass pushed down : effect of DK threshold ?? ✔

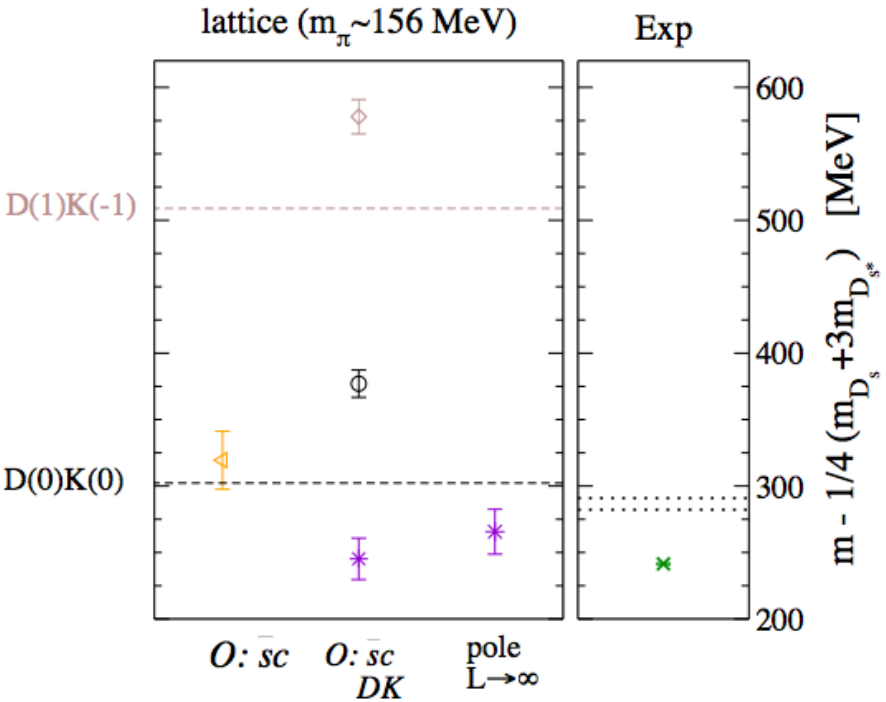


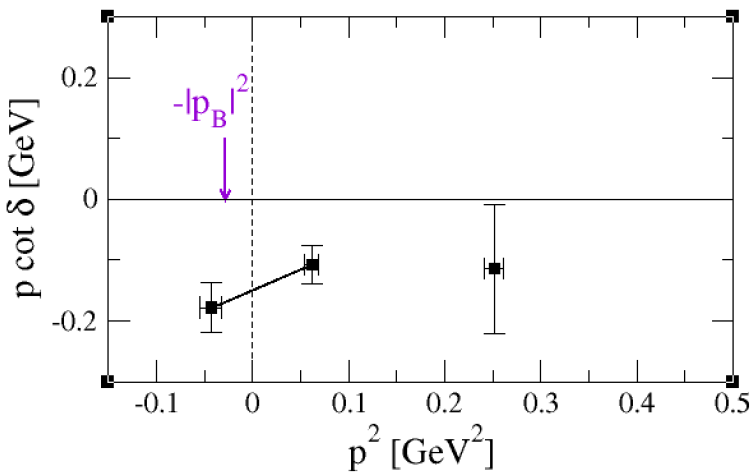
D_s mesons (near-threshold)

[C. Lang, L. Leskovec, D. Mohler, S.P. , R. Woloshyn: PRL 2013, PRD 2014]

$D_{s0}^*(2317)$ and DK scattering

Energies from a lattice





- $E=E_K(p)+E_D(-p)$ renders p
- δ for DK scattering in s-wave

  extracted using Luscher's relation

$$p\cot\delta(p) = \frac{2\mathcal{Z}_{00}(1;(\frac{pL}{2\pi})^2)}{L\sqrt{\pi}}$$

- Interpolation: effective range

$$p\cot\delta(p) = \frac{1}{a_0} + \frac{1}{2}r_0 p^2$$

$$a_0 = -1.33 \pm 0.20 \text{ fm}$$

$$r_0 = 0.27 \pm 0.17 \text{ fm}$$

a_0<0 indicates a state below th.

- pole position of $D_{s0}^*(2317)$

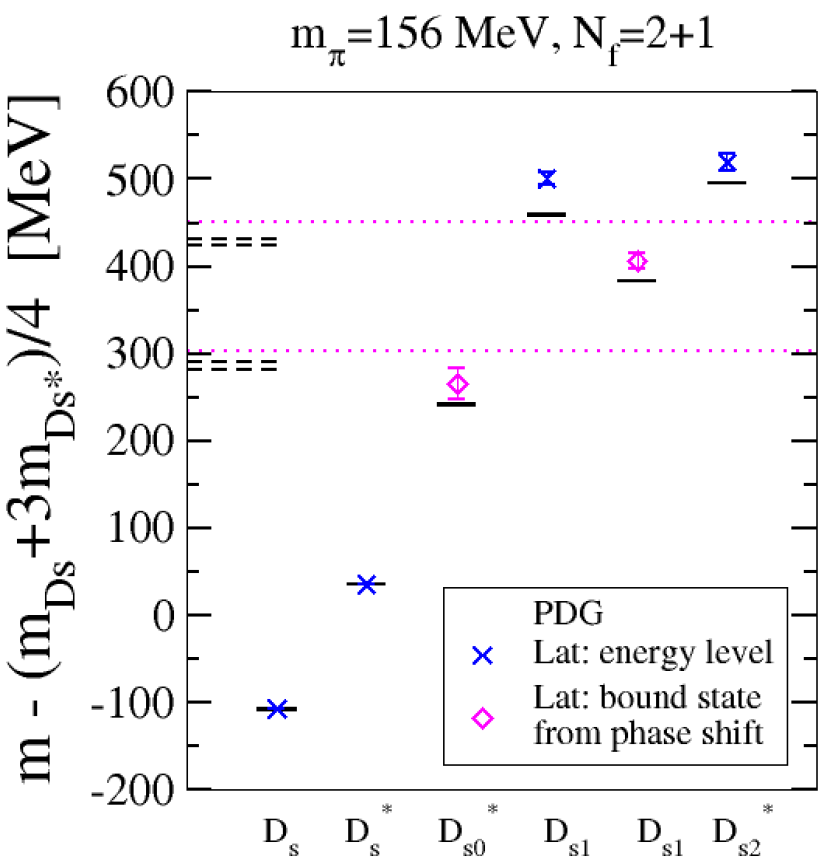
$$S \propto [\cot\delta - i]^{-1} = \infty, \quad \cot\delta(p_B) = i$$

$$K^{-1} = p_B \cot\delta = (i\,|\,p_B\,|)i = -\,|\,p_B\,|$$

$$m_{D_{s0}}^{lat,L\to\infty} = E_D(p_B) + E_K(p_B)$$

D. Mohler, C. Lang, L. Leskovec, S.P. , R. Woloshyn: 1308.3175, PRL : m_π≈156 MeV, L≈2.9 fm, Nf=2+1

D_{s0} and D_{s1} below DK and D^*K thresholds



- D_{s0} and D_{s1} have been observed experimentally
- Quark models expected them above thresholds but they were found below them
- Our post-dictions agree with measured masses
- Threshold effect lowers their masses

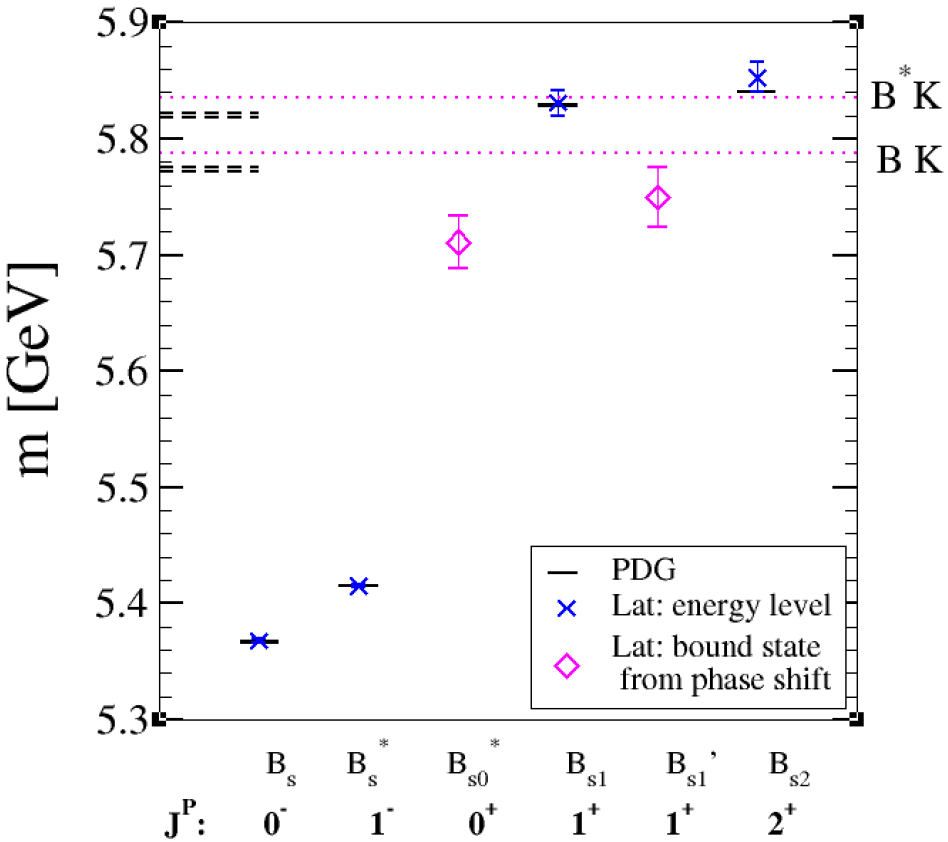
- Composition of resulting D_{s0} and D_{s1} analyzed via Weinberg-type compositness conditions: [Martinez Torres, E. Oset, S.P., A. Ramos: 1412.1706]

[D. Mohler, C. Lang, L. Leskovec, S.P. , R. Woloshyn: 1308.3175, Phys. Rev. Lett 2013 1403.8103, PRD 2014]

Mass prediction for missing B_{s0} and B_{s1}

$$\mathcal{O} = \bar{s}b, \;\; BK$$

m_π=156 MeV



Quantities shown:

for two bound states :

$$m_B = (m_B - E_{th})^{lat} + E_{th}^{exp}$$

for other states :

$$m = (m - \bar{m})^{lat} + \bar{m}^{lat}$$

for dotted lattice thresholds :

$$E_{th} = (E_{th} - \bar{m})^{lat} + \bar{m}^{lat}$$

$$\bar{m} \equiv \frac{1}{4}(m_{Bs} + 3m_{Bs^*})$$

- B_{s1}' and B_{s2} agree well with exp

Predictions:

- B_{s0} bound state below BK th.
- B_{s1} bound state below B*K th.

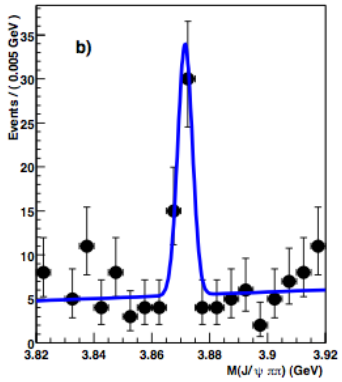
[C. Lang, D. Mohler, S.P. , R. Woloshyn: 1501.0164]

X(3872) , $J^{PC}=1^{++}$, charmonium-like

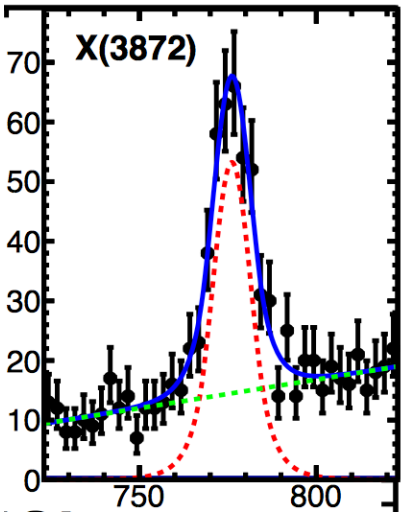
- First charmonium-like state discovered [Belle, PRL, 2003]
- sits within 1 MeV of $D^0\underline{D}^{0*}$ threshold
  8 MeV below D^+D^{*-} threshold
  } isospin breaking effects my be important
- believed to have a large molecular $D^0\underline{D}^{0*}$ Fock component
- $\Gamma < 1.2$ MeV
- decays to I=0, 1 equally important

  X(3872) → J/Ψ ω (I=0)

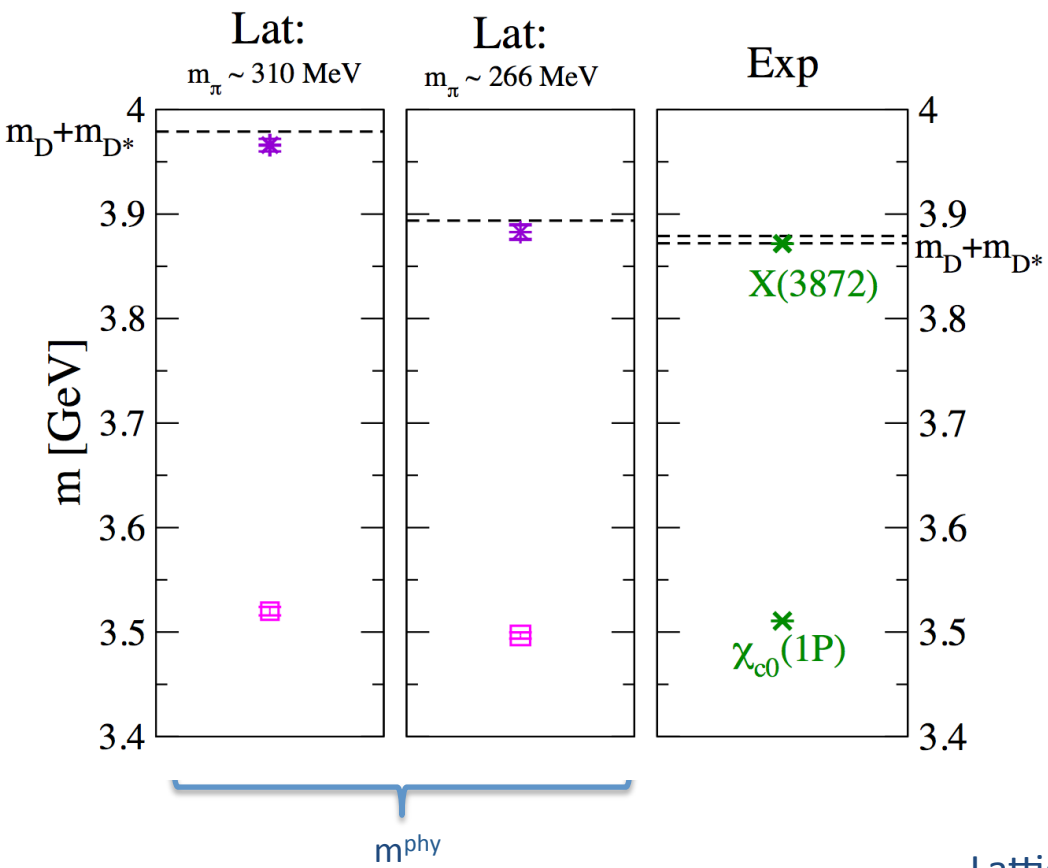
  X(3872) → J/Ψ ρ (I=1)



[LHCb, PRL 2013]



X(3872) as bound state from DD* scattering, $J^{PC}=1^{++}$, I=0



$\mathcal{O}: \bar{c}\, c,\ D\bar{D}^*$

- ground state: $\chi_{c1}(1P)$
- DD* scattering matrix near th. determined
- A pole of $T \propto \frac{1}{\cot\delta - i} = \infty$ found just below th. (violet star)
- The pole attributed to X(3872), which is a shallow bound state in both simulations
- Position of DD* threshold depends on $m_{u/d}$, and may be affected by discretization effects related to charm quark

X(3872)	m - ($m_{D0}+m_{D0*}$)
lat (m_π=310 MeV)	- 13 ± 6 MeV
lat (m_π=266 MeV)	- 11 ± 7 MeV
exp	- 0.14 ± 0.22 MeV

Lattice evidence for X(3872):

- $m_\pi \approx 266$ MeV, a=0.124 fm, L= 2 fm

[S.P. and Leskovec: 1307.5172, PRL 2013]

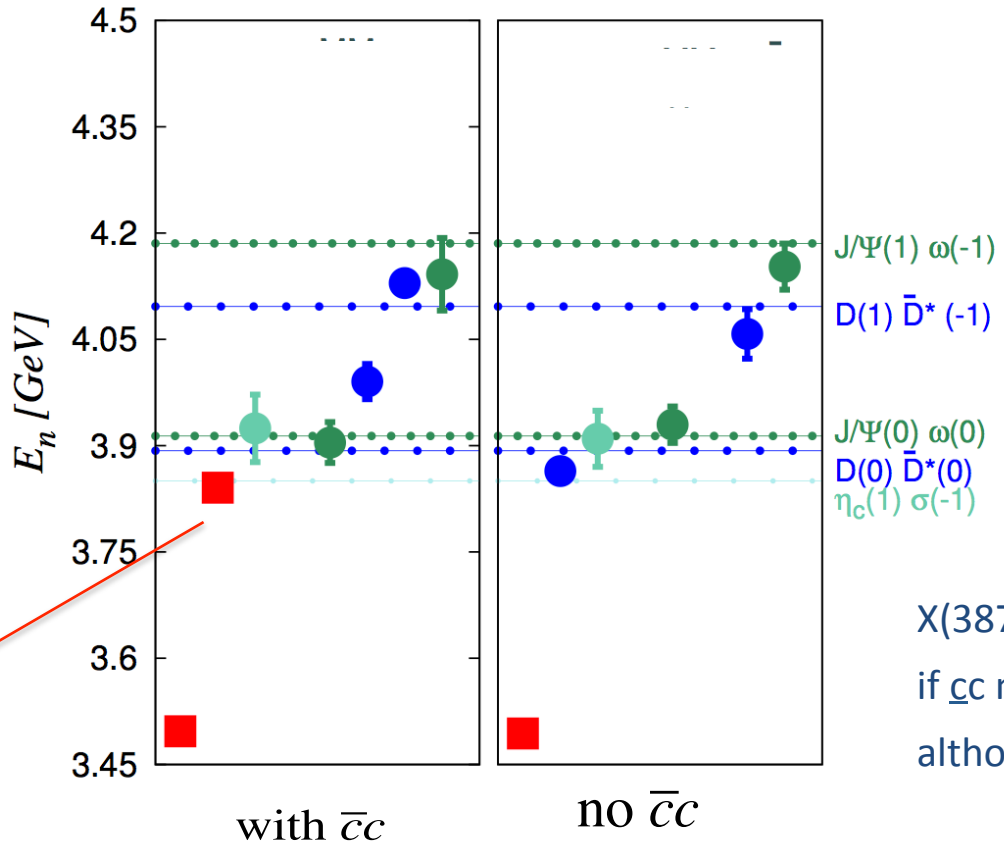
- $m_\pi \approx 310$ MeV, a=0.15 fm, L=2.4 fm , HISQ

[Lee, DeTar, Na, Mohler , update of proc 1411.1389]

Which Fock components are essential for X(3872) with I=0 ?

$J^{PC}=1^{++}$ $\quad \mathcal{O}: \underbrace{\bar{c}\,c, \quad D\bar{D}^*}_{\text{essential}}, \; J/\psi\,\omega, \; \chi_{c1}\eta, \; \eta_c\sigma, \; \underbrace{[\bar{c}u]_{3c}[cu]_{3c}, \; [\bar{c}u]_{6c}[cu]_{6c}}_{\text{do no seem not essential}}$

energies of eigenstates



X(3872) found only if cc in the basis

X(3872) not found if cc not in the basis, although [cu][cu] in the basis

[M. Padmanath, C.B. Lang, S.P., 1503.03257]

Searches for hadrons with exotic flavor

very challenging

$$Z_c = \bar{c}c\bar{d}u, \quad Z_b = \bar{b}b\bar{d}u, \quad P_c = \bar{c}cuud$$

few searches — to challenging — to recent

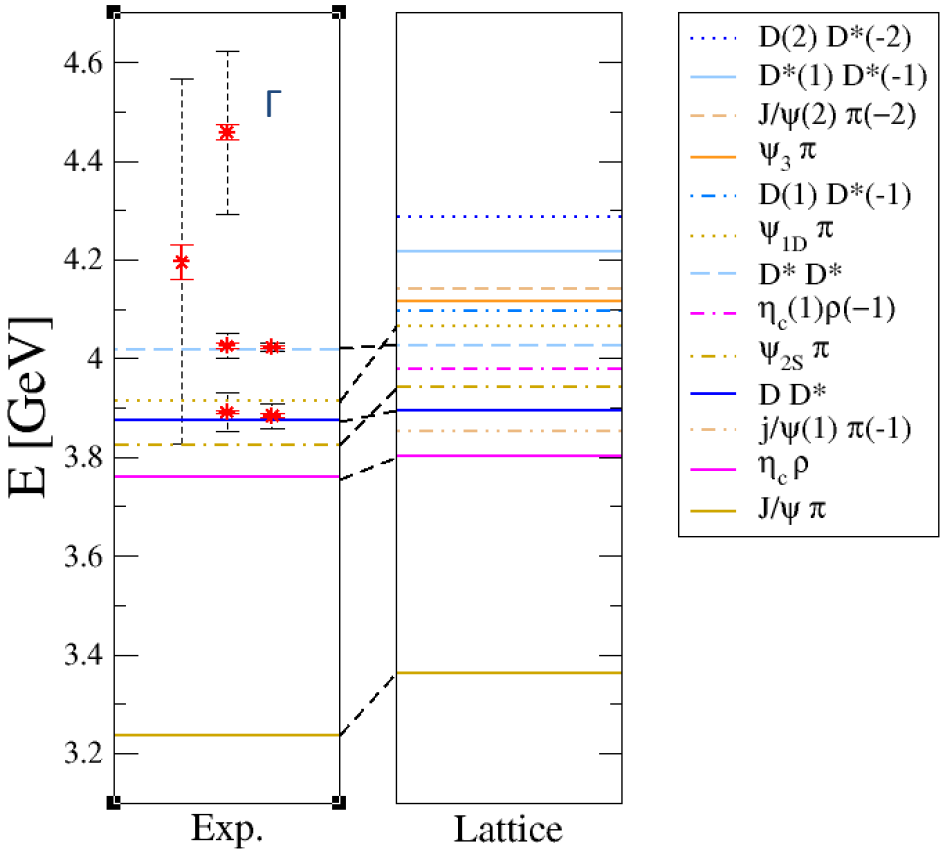
Even more challenging since most of experimental exotic states

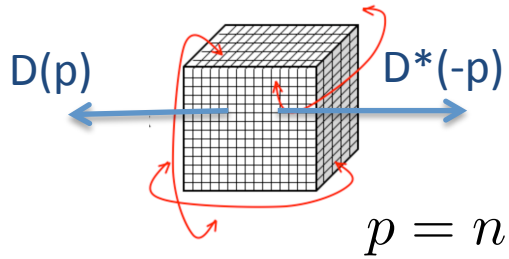
- are above several thresholds and decay to several two-meson final states
- require simulation of coupled-channels

Coupled-channel scattering matrix extracted on lattice so far only in

- a-la Luscher: Kπ, Kη system [Willson, Dudek, Edward, Thomas, HSC, PRL 2014]

Z_c^+ channel : $I^G=1^+$, $J^{PC}=1^{+-}$





$$p = n\frac{2\pi}{L}$$

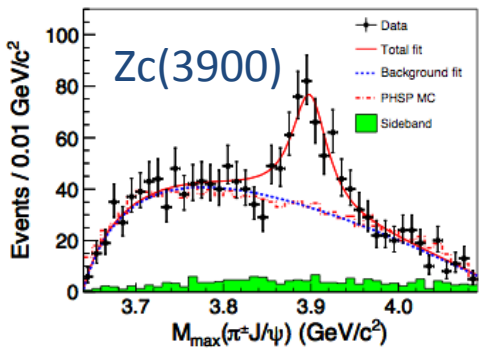
Lattice:

Horizontal lines represent energies of 13 two-meson states in non-interacting case

$E = E[\, M_1(p_1)\,] + E[\, M_2(p_2)\,]$

Extracting 13 two-meson states is a challenge

$$\mathcal{O} : (\bar{c}u)(\bar{d}c),\ \ (\bar{c}c)(\bar{d}u),\ \ [\bar{c}\bar{d}][cu]$$

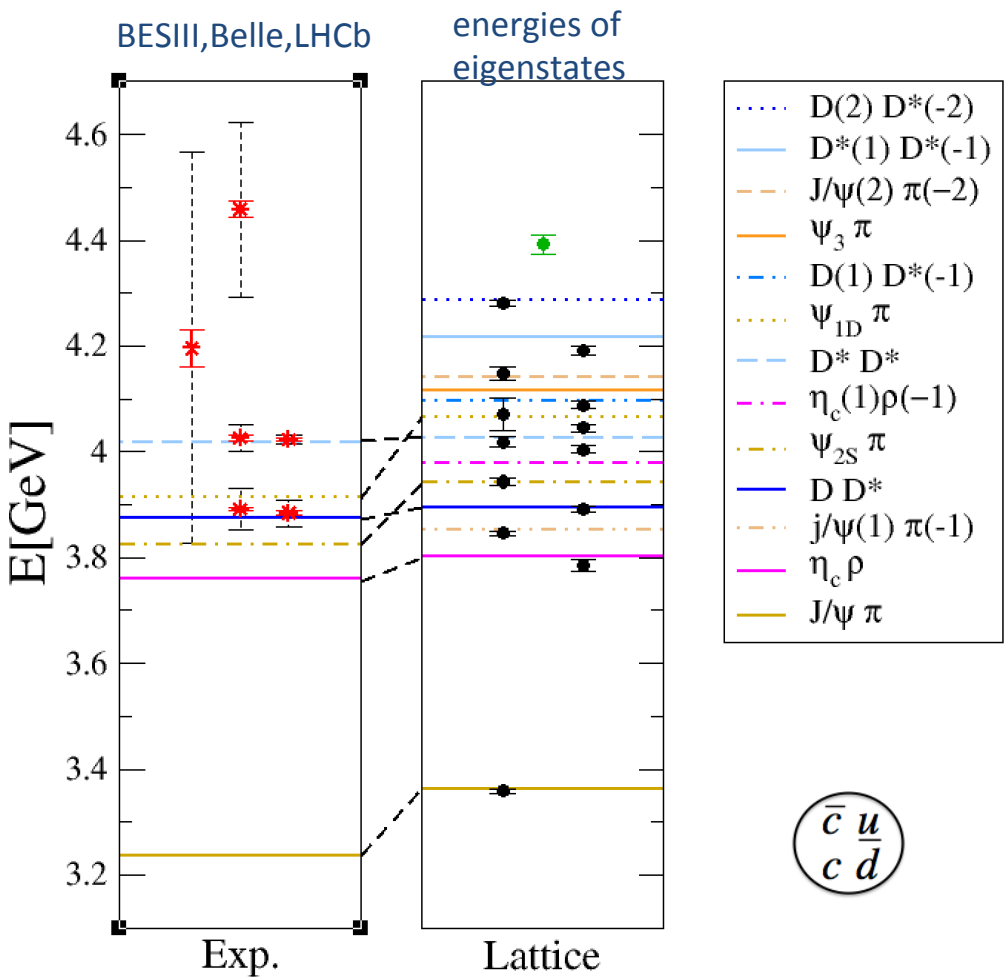


[S.P., Lang, Leskovec, Mohler, 1405.7612, PRD 2015]

Ensemble (2), $m_\pi \approx 266$ MeV, $L \approx 2$ fm, Nf=2

Z_c^+ channel : $I^G=1^+, J^{PC}=1^{+-}$

$$\mathcal{O} : (\bar{c}u)(\bar{d}c),\ (\bar{c}c)(\bar{d}u),\ [\bar{c}\bar{d}][cu]$$



Lattice results:

- 13 expected two-meson eigenstates found as expected (black circles)
- no additional eigenstate below 4.2 GeV
- no candidate for Z_c^+ below 4.2 GeV

Exp:

- $Z_c^+(3900)$ confirmed by three exp

Puzzle:

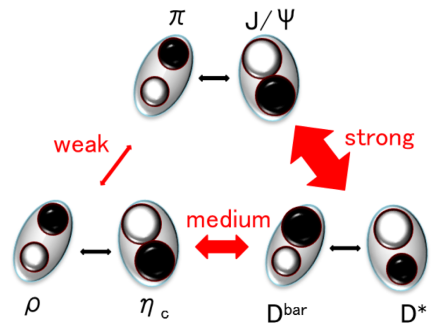
- why no eigenstate for Zc(3900) ?
- it would be naively expected if Zc(3900) related to a resonance pole
- is Zc(3900) of a different origin?
- perhaps coupled channel effect?

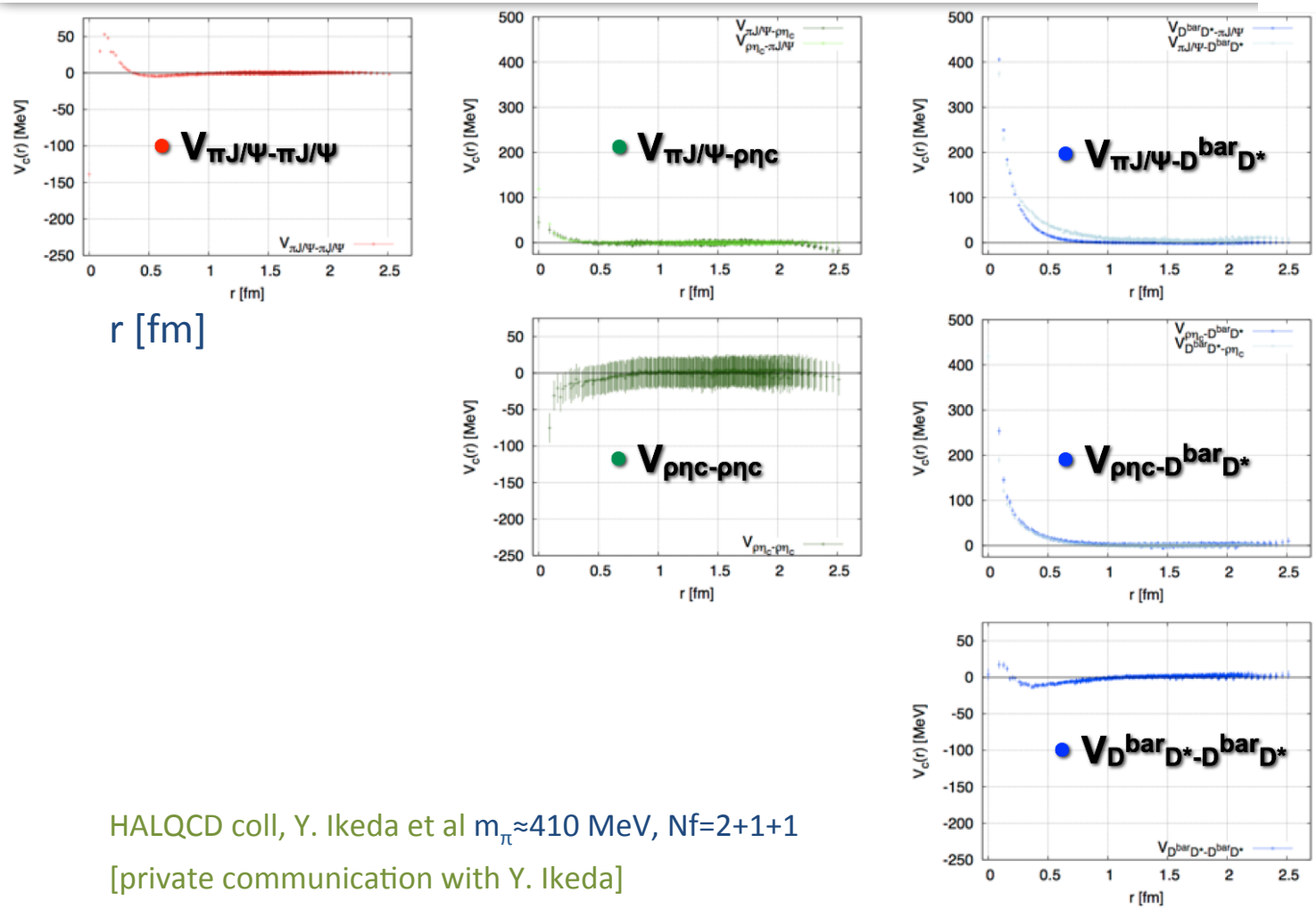
[S.P., Lang, Leskovec, Mohler, 1405.7612, PRD 2015]

similar conclusion [S.-H. Lee, C. DeTar, H. Na, 1411.1389]

Z_c^+ channel : $I^G=1^+$, $J^{PC}=1^{+-}$

- HALQCD method [application on H-dibaryon: HALQCD, 1504.01717]





r [fm]

indication for coupled channel eff.

HALQCD coll, Y. Ikeda et al $m_\pi \approx 410$ MeV, Nf=2+1+1

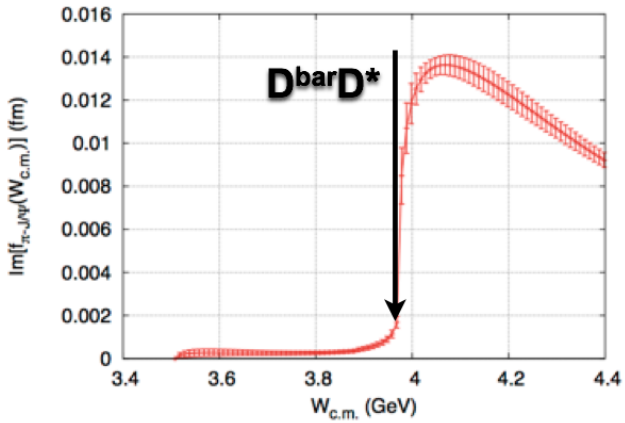
[private communication with Y. Ikeda]

Z_c^+ channel : $I^G=1^+$, $J^{PC}=1^{+-}$

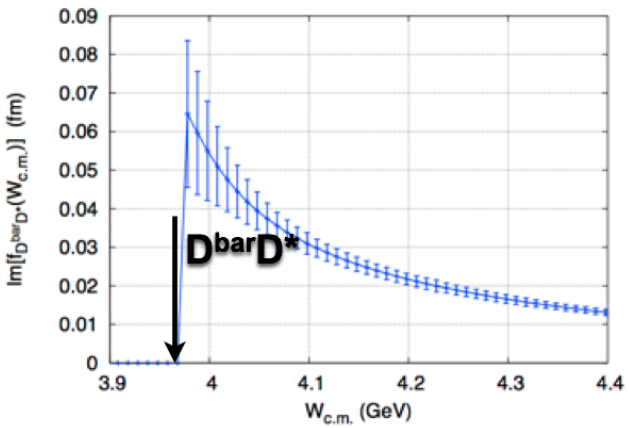
Lattice quantity related to cross-section

Experimental cross-section related to $Z_c^+(3900)$

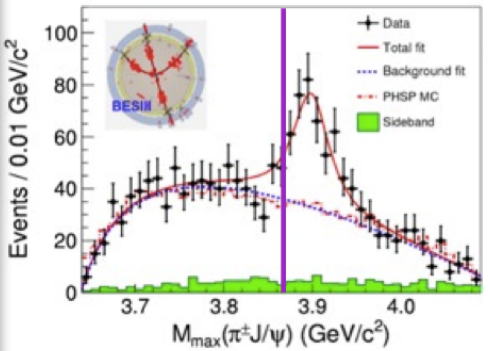
- **πJ/Ψ invariant mass (m_π=410MeV)**



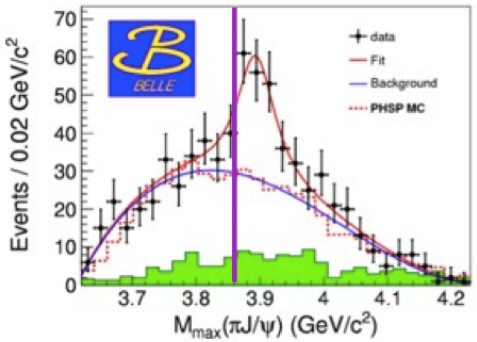
- **$D^{bar}D^*$ invariant mass (m_π=410MeV)**



- **e^+e^- --> π(πJ/Ψ) @ 4.26GeV**

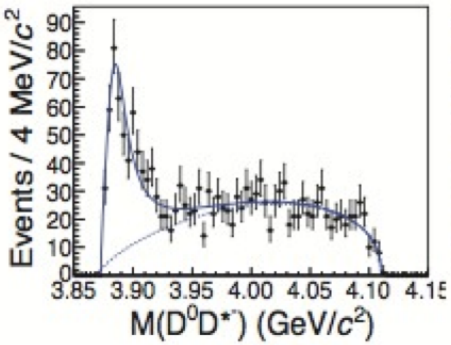


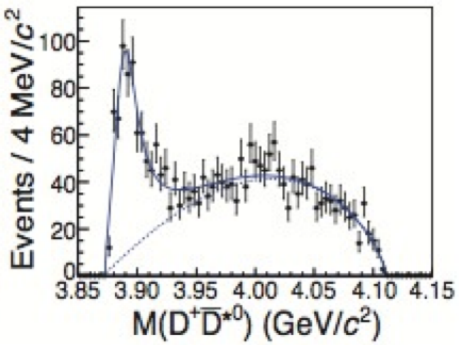
BESIII Coll., PRL110 (2013).



Belle Coll., PRL110 (2013).

- **e^+e^- --> $\pi^{+/-}$ $(D^{bar}D^*)^{-/+}$**





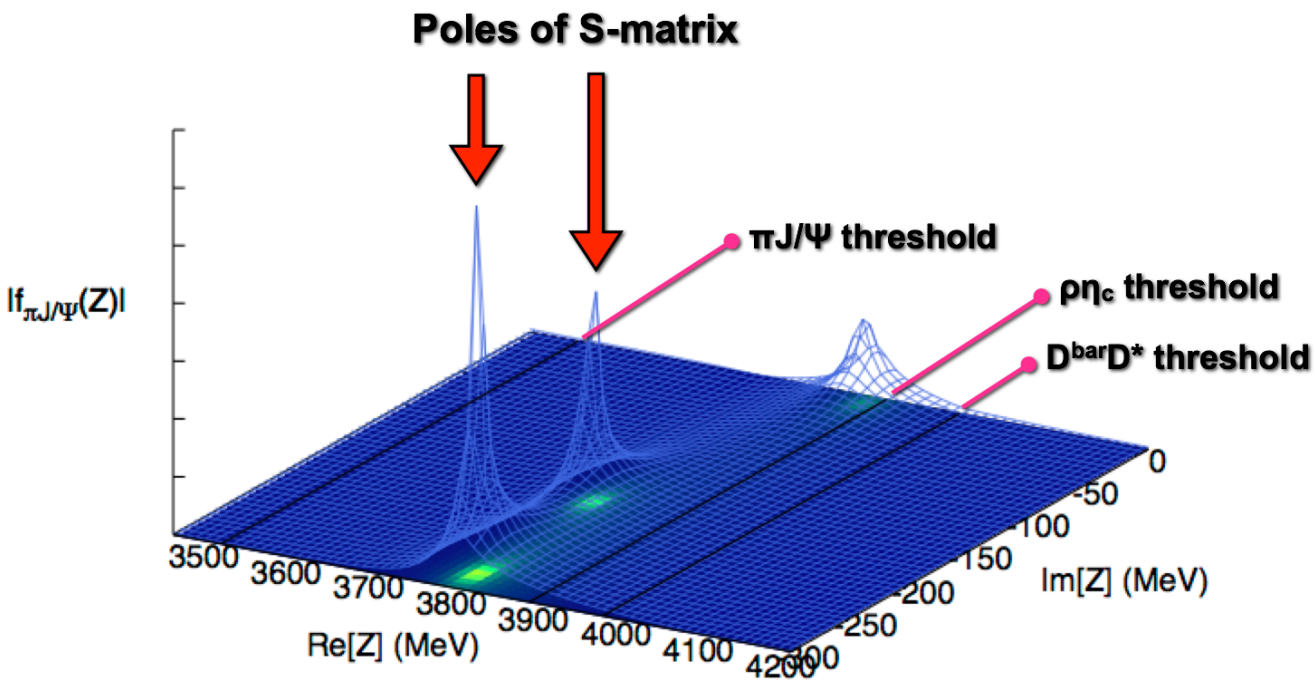
BESIII Coll., PRL112 (2014).

HALQCD coll, Ikeda et al m_π≈410 MeV, Nf=2+1+1

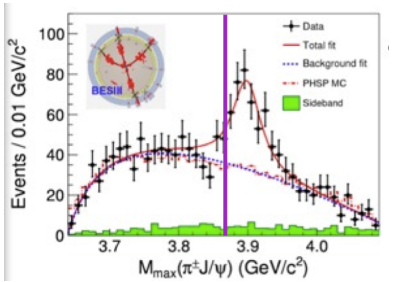
[private communication with Y. Ikeda]

Lineshapes resemble experimental Zc(3900).

Z_c^+ channel : $I^G=1^+$, $J^{PC}=1^{+-}$



DD*



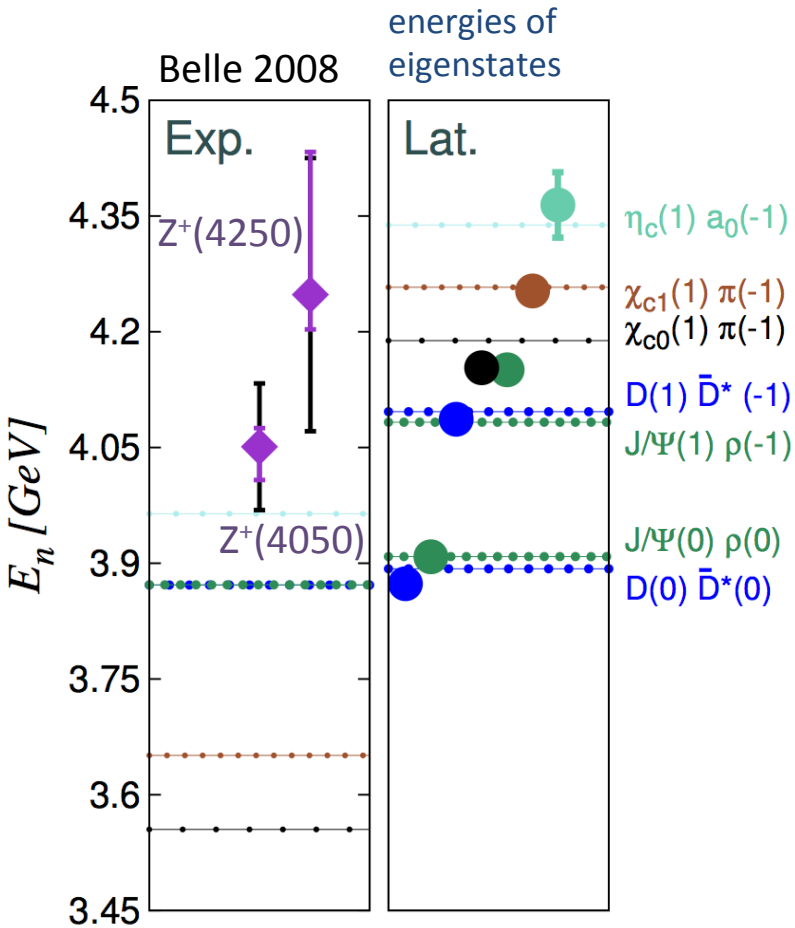
- exp: $Z_c^+(3900)$ peak appears above DD* th.
- HALQCD: poles found BELOW DD* th.

  pole NOT interpreted as a resonance (such a pole would be expected above DD*)
- Remains to be seen if HALQCD result is consistent with absence of Zc eigenstate in S.P. et al, PRD 2015
- Scattering matrix has not yet been extracted from Luscher approach

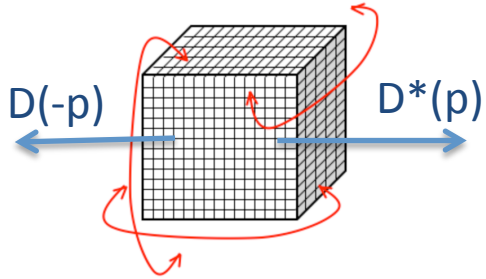
HALQCD coll, Ikeda et al $m_\pi \approx 410$ MeV, Nf=2+1+1

[private communication with Y. Ikeda]

charged partner of X(3872); channel $I^G=1^-$, $J^{PC}=1^{++}$, $\underline{c}\underline{c}\underline{d}u$



[M. Padmanath, C.B. Lang, S.P., 1503.03257]



$$p = n\frac{2\pi}{L}$$

$$\mathcal{O} : (\bar{c}u)(\bar{d}c),\ \ (\bar{c}c)(\bar{d}u),\ \ [\bar{c}\bar{d}][cu]$$

- Horizontal lines: energies of expected two-meson states in limit of no interaction: $E = E[\ M_1(p_1)\ ] + E[\ M_2(p_2)\ ]$
- Circles: energies of eigenstates from latt
- Only expected two-meson states observed.
- No lattice candidate for charged X(3872). In agreement with absence of such state in exp.
- No lattice candidate for other charged state. Two Belle 2008 states are exp. unconfirmed.

Y(4140), $J^{PC}=?^{?+}$, $c\bar{c}s\bar{s}$

$$Y(4140) \to J/\psi\ \phi$$
$$\bar{c}c \quad \bar{s}s$$

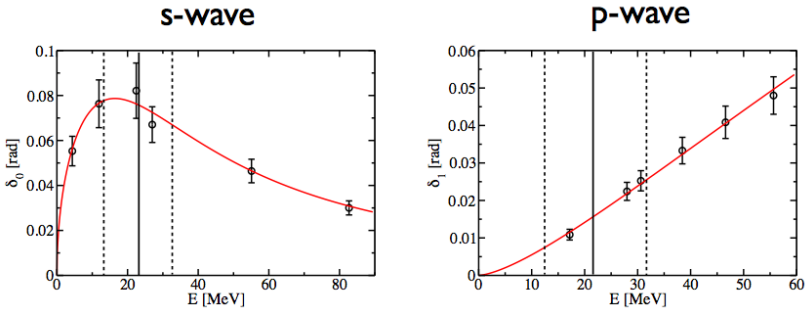
Experiment:

peak in J/ψ Φ just above J/ψ Φ threshold

found: CDF 2009, CMS 2012, D0 2013, Babar 2015

not found: Belle 2010, LHCb 2012

J/Ψ Φ scattering phase shift [rad]



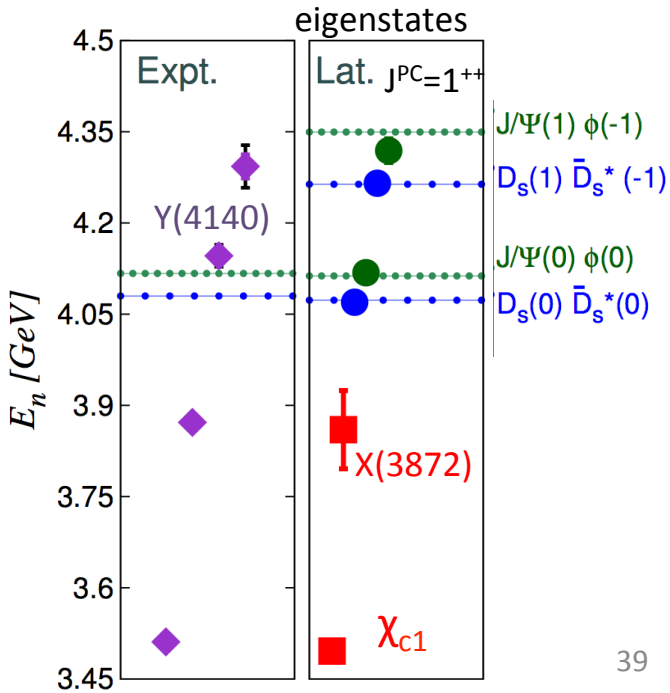
Lattice:

- S. Ozaki and S. Sasaki, 1211.5512, PRD
  strange quark annihilation neglected
  no resonant structure found for

- M. Padmanath, C.B. Lang, S.P., 1503.03257

$$\mathcal{O}: \ \bar{c}c, \ (\bar{c}s)(\bar{s}c), \ (\bar{c}c)(\bar{s}s), \ [\bar{c}\bar{s}][cs]$$

channel $J^P=1^+$ considered only: expected two-particle eigenstates found and χ_{c1}, X(3872) but not Y(4140)



Conclusions

- **Experiment** (in brief): exciting tetraquark and pentaquark candidates

- **Phenomenology:** many many interpretations of exotic experimental "bumps", some of them do not need to introduce exotic hadrons to explain experimental data

- **Lattice QCD** (in brief):

❑ Evidence found for many hadrons with **non-exotic** flavor:

- states well below th. : charmonium , D, π, K ... and all the others
- shallow bound states : D_{s0}, D_{s1}, B_{s0}, B_{s1} , X(3872) with I=0
- resonances via BW : ρ, K*, $K_0^*(1430)$, K_2, D_0^*, D_1, a_1, b_1 , Ψ(3770)

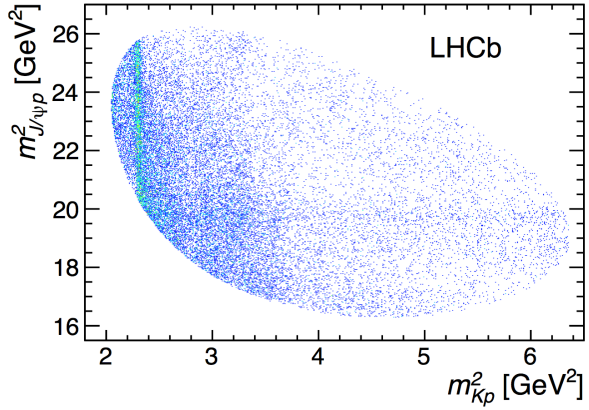
❑ Hadrons with **exotic** flavor

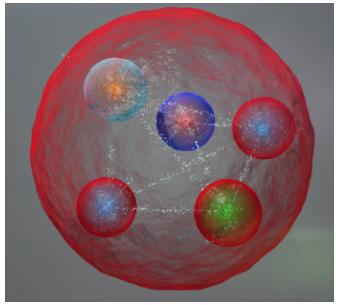
- if exotic hadrons were below strong decay threshold, they would be easy to (dis)prove in LQCD
- unfortunately, most of exotic hadrons can decay into several channels via strong interaction
- therefore lattice has not given yet a final answer which (if any) exotic hadrons arise from QCD
- I have (hopefully) given you some flavor in which direction progress is going
- this is an exciting topic at present and I am looking forward to face further challenges it poses ...

Backup slides

Pentaquarks

$$\Lambda_b^- \to K^- J/\psi \; p$$





LHCb: 1507.03414

14th July 2015

$$P_c^+ \to J/\psi \;\; p$$

$$\bar{c}c \;\; uud$$

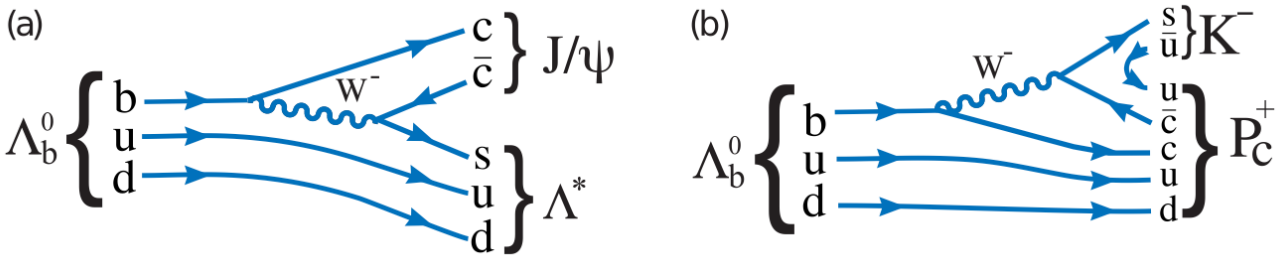
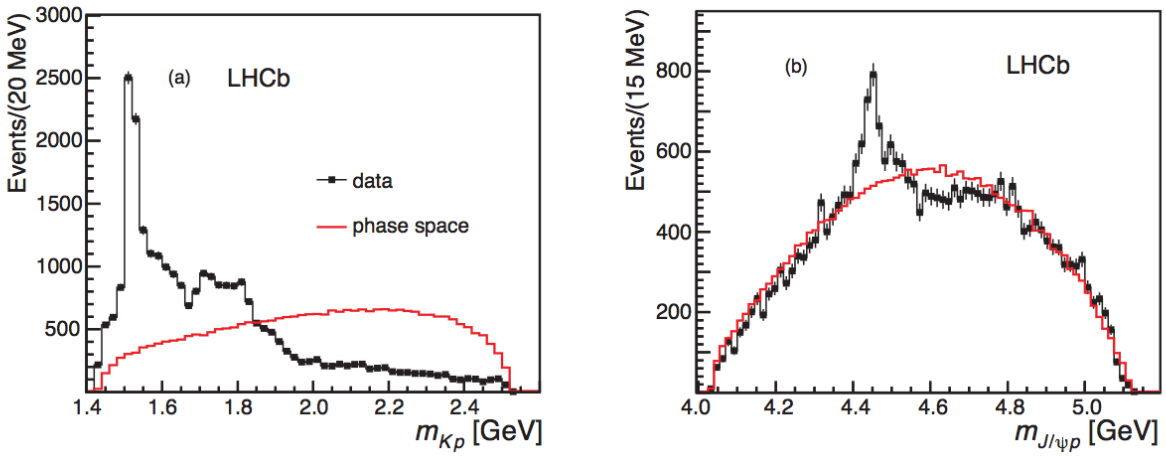


Figure 1: Feynman diagrams for (a) $\Lambda_b^0 \to J/\psi\,\Lambda^*$ and (b) $\Lambda_b^0 \to P_c^+ K^-$ decay.

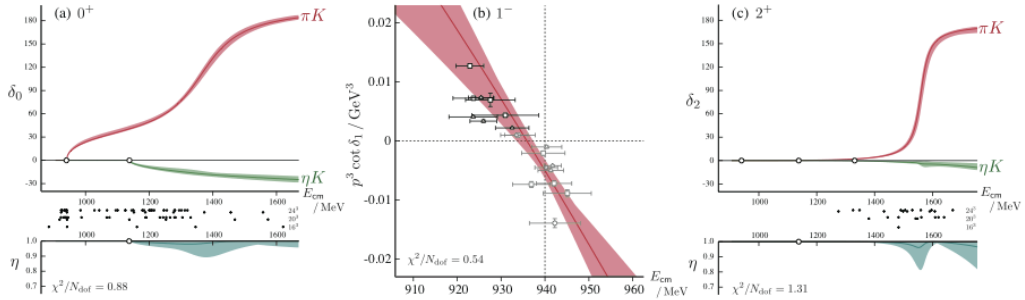


State	J^P	M_0 (MeV)	Γ_0 (MeV)
$\Lambda(1405)$	$1/2^-$	$1405.1^{+1.3}_{-1.0}$	50.5 ± 2.0
$\Lambda(1520)$	$3/2^-$	1519.5 ± 1.0	15.6 ± 1.0
$\Lambda(1600)$	$1/2^+$	1600	150
$\Lambda(1670)$	$1/2^-$	1670	35
$\Lambda(1690)$	$3/2^-$	1690	60
$\Lambda(1800)$	$1/2^-$	1800	300
$\Lambda(1810)$	$1/2^+$	1810	150
$\Lambda(1820)$	$5/2^+$	1820	80
$\Lambda(1830)$	$5/2^-$	1830	95
$\Lambda(1890)$	$3/2^+$	1890	100
$\Lambda(2100)$	$7/2^-$	2100	200
$\Lambda(2110)$	$5/2^+$	2110	200
$\Lambda(2350)$	$9/2^+$	2350	150
$\Lambda(2585)$	?	$\approx$2585	200

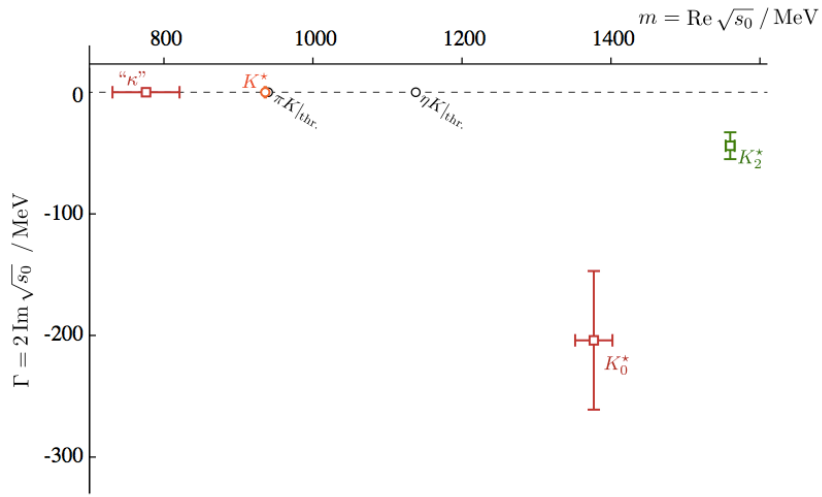
Resonances in Kπ, Kη coupled channels

- $\bar{q}q$, Kπ, Kη interpolators
- a number of different 0<P≤2
- for each E_n: one determinant equation for many unknowns
- T-matrix parametrized to get around this problem
- the location of poles of T-matrix in complex given below
- K*(892) and κ are below threshold for this m
- K_0^* , K_2^* are resonances
- m_π=391 MeV, N_L=16, 20, 24

[Dudek, Edwards, Thomas, Wilson, HSC, 1406.4158, PRL 2014]



$$t_{ii} = \frac{(\eta e^{2i\delta_i} - 1)}{2i\rho_i}, \; t_{ij} = \frac{\sqrt{1-\eta^2}\, e^{i(\delta_i+\delta_j)}}{2\sqrt{\rho_i\,\rho_j}}$$

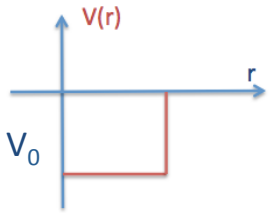


$$\det\left[\delta_{ij}\delta_{JJ'} + i\rho_i\, t_{ij}^{(J)}(E_{\mathrm{cm}})\left(\delta_{JJ'} + i\mathcal{M}_{JJ'}^{\vec{P}\Lambda}(p_i L)\right)\right] = 0,$$

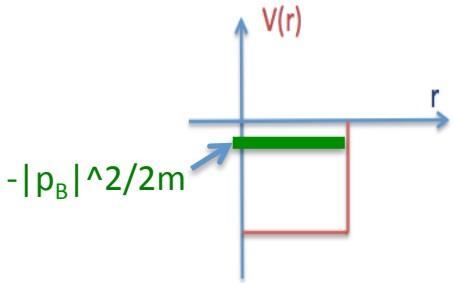
location of poles in T matrix in complex plane

How to identify a shallow bound state? (continued)

V0 that provides no bound st.

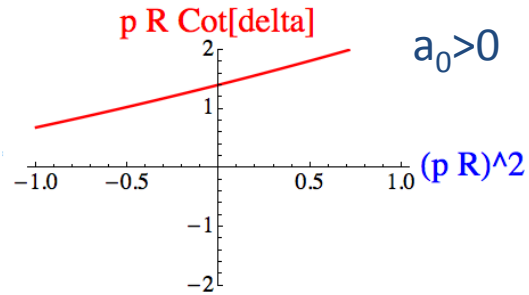


V0 that provides one shallow bound st.

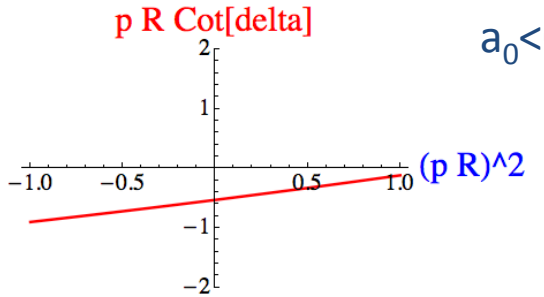


$$\lim_{p\to 0} p \cot\delta(p) = \frac{1}{a_0}$$

a_0=scattering length

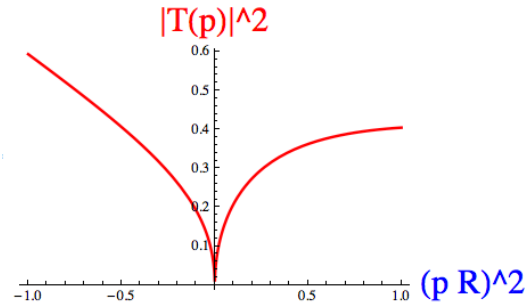


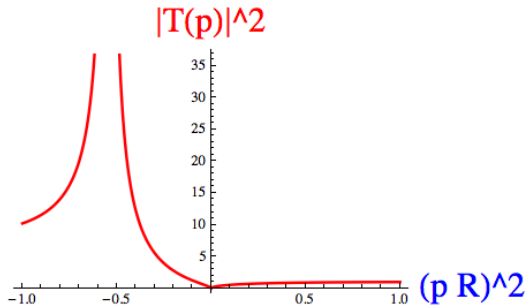
$a_0>0$



$a_0<0$

$$|T(p)|^2 \propto \left|\frac{1}{\cot\delta - i}\right|^2$$





$-(|p_B|\,R)$^2

scattering matrix
has a pole

Hadrons near or above threshold: single-hadron approximation

- only interpolating fields $\mathcal{O} \approx \bar{q}\, q$ or $\mathcal{O} \approx q\, q\, q$
- assumptions: all energy levels correspond to "one-hadron" states

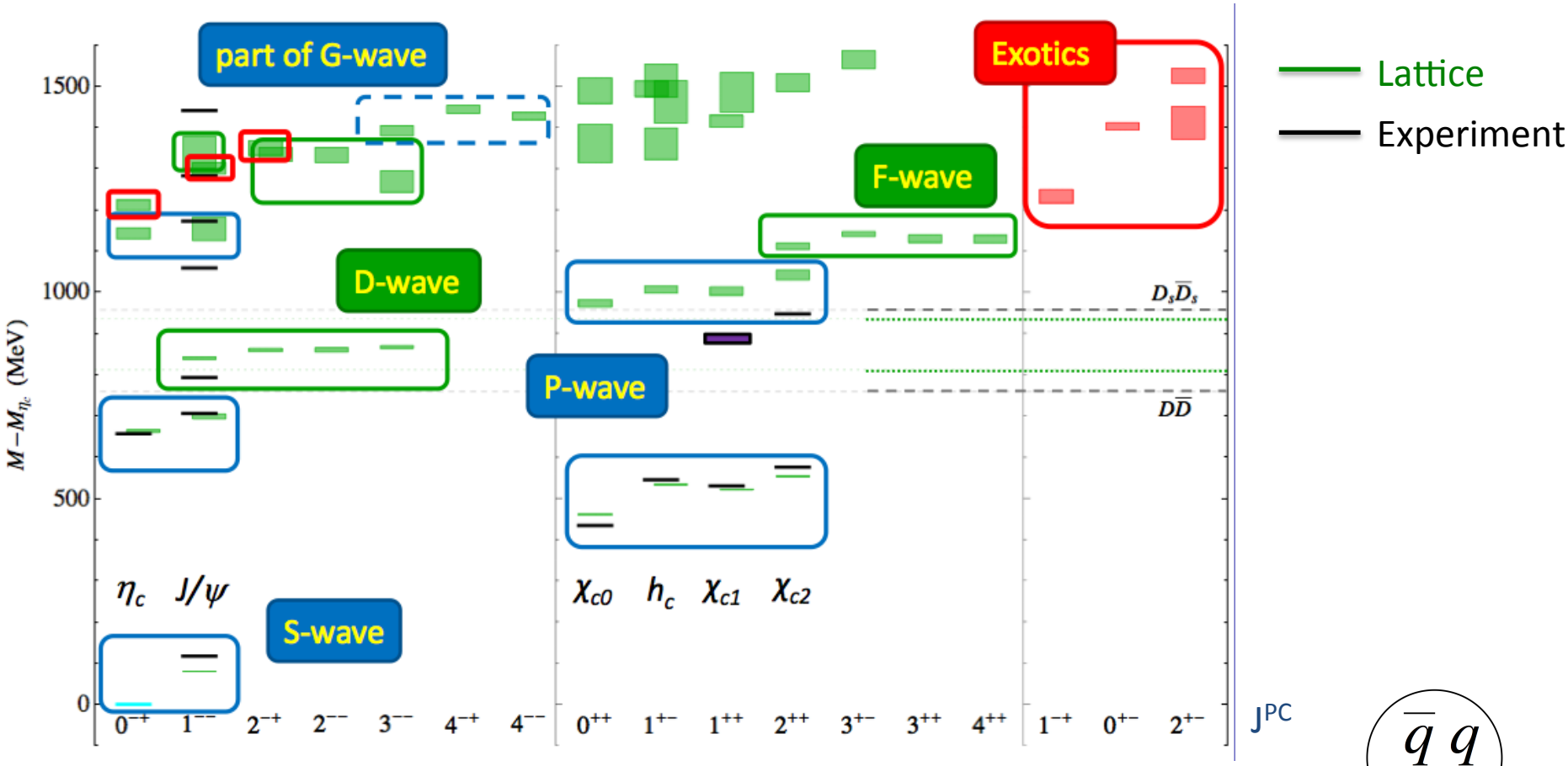
  no two-hadron state is seen

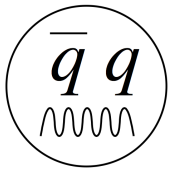
  m=E (for P=0)

  these are strong assumptions ...

  but results still present valuable reference point

$\underline{c}c$ spectrum: single-hadron approximation





[HSC , L. Liu et al: 1204.5425, JHEP]

- $m_\pi \approx 400$ MeV, $L \approx 2.9$ fm, Nf=2+1
- identification with $n\ ^{2S+1}L_J$ multiplets using $\langle O|n\rangle$
- green: lat, black: exp

Hybrids:

some of them have exotic J^{PC}

large overlap with $O = \underline{q}\, F_{ij}\, q$

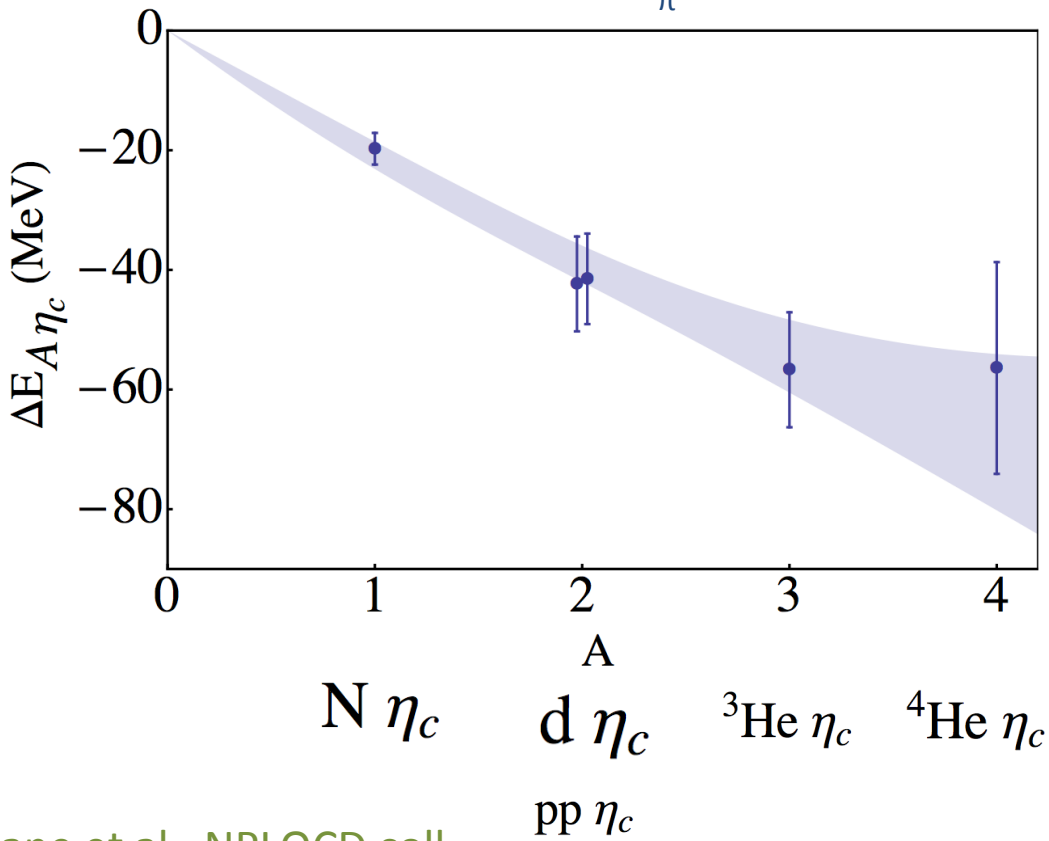
Bound states of a charmonium and a nucleus

Exp: no reliable candidates for such states yet

Lat: SU(3) symmetric point, $m_u=m_d=m_s$, m_π~800 MeV, L=3.4, 4.5, 6.7 fm

Bound states found at this m_π. Not known yet whether these bound states survive at m_π^{phys}

m_π~800 MeV



Binding energies

for a system composed of

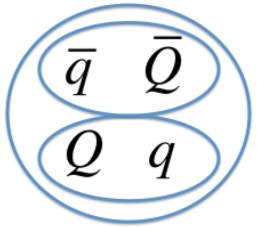
η_c + nucleus (A)

number of nucleons in nucleus

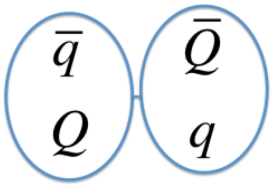
[Beane et al., NPLQCD coll. , 1410.7069]

Phenomenological approaches

Approaches that support new exotic hadrons



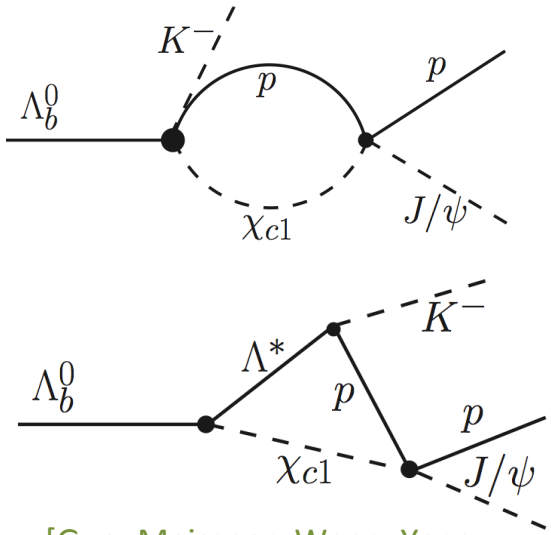
- compact tetraquark: diquark antidiquark



- meson molecules
- ...

Approaches that do not contain new exotic hadrons, and explain bumps in experimental cross sections via

- kinematical effects due to nearby threshold
- coupled channel effects
- kinematics interplay in triangle diagrams
- ...



[Guo, Meissner, Wang, Yang 1507.04950]

I will try to address the question whether observed exotic and conventional hadrons arise directly from QCD.

ρ resonance: lattice results

from T. Yamazaki [lat14 plenary, 1503.0867]

$$\Gamma = \frac{g^2_{\rho\pi\pi}}{6\pi}\frac{p^3}{E^2}$$